

About Bertrand Curves and Type-1 Bishop Frame in Three-Dimensional Weyl Space W_3

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Abstract

In this paper, we have defined Bertrand curves in three-dimensional Weyl space. Then we have given the relations between the Bishop vector fields of Bertrand curve pair. Finally, while $(C, \overline{C})$ is Bertrand curve pair, we have obtained the equalities depending on $\overline{K}_1$ and $\overline{K}_2$.

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1 Introduction

A manifold with a conformal metric g_{ij} and a symmetric connection ∇_k satisfying the compatibility condition

$$\nabla_k g_{ij} - 2T_k g_{ij} = 0 \quad (1)$$

is called a Weyl space which will be denoted by $W(g_{ij}, T_k)$. The vector field T_k is named the complementary vector field. Under a normalization of the metric tensor g_{ij} in the form

$$\tilde{g}_{ij} = \lambda^2 g_{ij} \quad (2)$$

the complementary vector field T_k is transformed by the law

$$\tilde{T}_k = T_k + \partial_k \ln \lambda \quad (3)$$

where λ is a scalar function [9]. If under the transformation (2), the quantity A is called a satellite of g_{ij} with weight $\{p\}$.

The prolonged derivative and prolonged covariant derivative of A are defined as

$$\dot{\partial}_k = \partial_k A - p T_k A \quad (4)$$

and

$$\dot{\nabla}_k A = \nabla_k A - p T_k A \quad (5)$$

respectively [4], [10]. The v^i_r ($i, r = 1, 2, 3$) be the contravariant components of the vector field v in $W_3(g_{ij}, T_k)$. Suppose that the vector field v are normalized by the conditions $g_{ij} v^i_r v^j_r = 1$ ($j, r = 1, 2, 3$).

The prolonged covariant derivative of the vector field v is given by [14]

$$\dot{\nabla}_k v^i_r = \overset{s}{T}_k \overset{s}{v}^i_r (s = 1, 2, 3). \quad (6)$$

The quantities

$$\overset{q}{\tau}_{rs} = \overset{q}{T}_k \overset{q}{v}^k_s (q = 1, 2, 3; r \neq s) \quad (7)$$

and

$$\overset{r}{\mathcal{C}}_s = \overset{r}{T}_k \overset{r}{v}^k_s \quad (8)$$

are called the Chebyshev curvature of the first kind and geodesic curvature of the net (v_1, v_2, v_3) , respectively [14].

The vector fields

$$\overset{a}{a}^i_{rs} = \overset{q}{\tau}_{rs} \overset{q}{v}^i_r, \quad \overset{c}{c}^i_s = \overset{r}{\mathcal{C}}_s \overset{r}{v}^i_s (i, q, r, s = 1, 2, 3) \quad (9)$$

are called the Chebyshev vector fields of the first kind and geodesic vector fields of the net (v_1, v_2, v_3) , respectively [14].

Since the net (v_1, v_2, v_3) is an orthogonal net, we have [14]

$$\overset{r}{T}_k = 0, \quad \overset{p}{T}_k + \overset{r}{T}_k = 0 \quad (r \neq p). \quad (10)$$

Bertrand curves have been determined by J. Bertrand [1] in 1850. Later, Bertrand curves have been studied by L.R. Pears [13] in 1935, by J.K. Wittermore [16] in 1940 and J.F. Burke [2] in 1960, respectively. Besides the properties of these curves provided in various studies [5], they have been handled in different spaces such as Riemann-Otsuki space [17], Galilean space [11], three-dimensional sphere [7], three-dimensional space forms [3], Euclidean 3-space [8] and Minkowski space-time[15].

2 Preliminaries

Let $C : x^i = x^i(s)$ be a curve in three-dimensional Weyl space W_3 (s is the arc length parameter of C). Let $\{v_1, v_2, v_3\}$ be Frenet frame and $\{v_1, n_1, n_2\}$ be Bishop frame of the curve C such that K_1, K_2 are the first and second curvatures and k_1, k_2 are the Bishop curvatures of C . The Frenet and Bishop formulas [6] of C are

$$\begin{aligned} v_1^k \dot{\nabla}_k v_1^i &= K_1 v_2^i, \\ v_1^k \dot{\nabla}_k v_2^i &= -K_1 v_1^i + K_2 v_3^i \\ v_1^k \dot{\nabla}_k v_3^i &= -K_2 v_2^i, \end{aligned} \quad (11)$$

and

$$\begin{aligned} v_1^k \dot{\nabla}_k v_1^i &= k_1 v_1^i + k_2 v_2^i, \\ v_1^k \dot{\nabla}_k n_1^i &= -k_1 v_1^i, \\ v_1^k \dot{\nabla}_k n_2^i &= -k_2 v_1^i. \end{aligned} \quad (12)$$

Since v_2^i is orthogonal to v_1 , v_2^i can be written as $v_2^i = a n_1^i + b n_2^i$ [6] where $a = g_{ij} v_2^i n_1^j = \cos\theta$, $b = g_{ij} v_2^i n_2^j = \cos(\frac{\pi}{2} - \theta) = \sin\theta$ and $\theta = \angle(v_2, n_1)$. In addition, since $v_3^i = \varepsilon_{ijk} v_1^j v_2^k$ ($k = 1, 2, 3$), the following equality is satisfied:

$$\begin{aligned} v_3^i &= \varepsilon_{ijk} v_1^j (\cos\theta n_1^k + \sin\theta n_2^k), \\ v_3^i &= n_2^i \cos\theta - n_1^i \sin\theta. \end{aligned} \quad (13)$$

Therefore

$$\begin{pmatrix} v_1^i \\ v_2^i \\ v_3^i \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\theta & \sin\theta \\ 0 & -\sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} v_1^i \\ n_1^i \\ n_2^i \end{pmatrix} \quad (14)$$

is valid.

From (11), $v_1^k \dot{\nabla}_k v_1^i = K_1 v_2^i$. Multiplying this equality by $g_{ij} v_2^j$ and taking summation on i and j , we get [6]

$$\begin{aligned} K_1 &= v_1^k (\dot{\nabla}_k v_1^i) g_{ij} v_2^j, \\ K_1 &= T_k^p v_1^k v_1^i g_{ij} v_2^j, \\ K_1 &= g_{ij} c_1^i v_2^j, \end{aligned} \quad (15)$$

or

$$K_1 = \frac{2}{1} \quad (16)$$

where $g_{ij} v_2^i v_2^j = 1$, $g_{ij} v_3^i v_2^j = 0$ and $T_k^1 = 0$.

Theorem 2.1 *If $K_1 = 0$, then the geodesic vector field c_1^i of the net (v_1, v_2, v_3) is orthogonal to v_2 .*

From (11), $v_1^k \dot{\nabla}_k v_3^i = -K_2 v_2^i$. Multiplying this relation by $g_{ij} v_2^j$ and summing on i and j , we obtain [6]

$$\begin{aligned} -K_2 &= v_1^k (\dot{\nabla}_k v_3^i) g_{ij} v_2^j, \\ -K_2 &= T_k^p v_3^k v_1^i g_{ij} v_2^j, \\ K_2 &= -g_{ij} a_{31}^i v_2^j, \end{aligned} \quad (17)$$

or

$$K_2 = -\frac{2}{31} \quad (18)$$

where $g_{ij} v_1^i v_2^j = 0$, $g_{ij} v_2^i v_2^j = 1$ and $T_k^3 = 0$.

Theorem 2.2 *If $K_2 = 0$, then the Chebyshev vector field of the first kind a_{31}^i of the net (v_1, v_2, v_3) is orthogonal to v_2 .*

Using $v_1^k \dot{\nabla}_k v_3^i = -K_2 v_2^i$ and [13] and then multiplying obtained equation by $g_{ij} v_2^j$, we have $K_2 = v_1^k \dot{\nabla}_k \theta$ ($\theta = \theta(s)$) where $g_{ij} n_2^i v_2^j = \sin \theta$, $g_{ij} n_1^i v_2^j = \cos \theta$ and $g_{ij} v_1^i v_2^j = 0$. Multiplying $v_1^k \dot{\nabla}_k v_1^i = k_1 n_1^i + k_2 n_2^i = K_1 v_2^i$ by $g_{ij} n_1^j$ we get

$$(v_1^k \dot{\nabla}_k v_1^i) g_{ij} n_1^j = k_1 = K_1 \cos \theta \quad (19)$$

and multiplying $v_1^k \dot{\nabla}_k v_1^i = k_1 n_1^i + k_2 n_2^i = K_1 v_2^i$ by $g_{ij} n_2^j$ we have

$$(v_1^k \dot{\nabla}_k v_1^i) g_{ij} n_2^j = k_2 = K_1 \sin \theta. \quad (20)$$

From (19) and (20), we obtain $k_1^2 + k_2^2 = K_1^2$.

Theorem 2.3 *If $k_1 = 0$, then the geodesic vector field c_1^i of the net (v_1, v_2, v_3) is orthogonal to n_1 .*

Theorem 2.4 *If $k_2 = 0$, then the geodesic vector field c_1^i of the net (v_1, v_2, v_3) is orthogonal to n_2 .*

3 Bertrand Curves in W_3

Let $\bar{C} : \bar{x}^i = \bar{x}^i(\bar{s})$ be other curve in W_3 ($\bar{s}$ is the arc length parameter of $\bar{C}$). Let us denote Frenet and Bishop components of $\bar{C}$ by $\{\bar{v}_1, \bar{v}_2, \bar{v}_3, \bar{K}_1, \bar{K}_2\}$ and $\{\bar{v}_1, \bar{n}_1, \bar{n}_2, \bar{k}_1, \bar{k}_2\}$, respectively.

Definition 3.1 *If the principal normal vector fields of the curves C and $\bar{C}$ are linear dependent, the curve pair $(C, \bar{C})$ is called Bertrand curve pair.*

If the curve pair $(C, \bar{C})$ is Bertrand curve pair the following equality is satisfied:

$$C(s) = \bar{C}(\bar{s}) + \lambda(\bar{s}) \bar{v}_2^i(\bar{s}). \quad (21)$$

Taking prolonged covariant derivative of (21) in the direction of $\bar{v}_1$ we have

$$\begin{aligned} \bar{v}_1^k \dot{\nabla}_k C &= (v_1^k \dot{\nabla}_k C) f(s) = v_1^i f(s) = \\ &= \bar{v}_1^i + (\bar{v}_1^k \dot{\nabla}_k \lambda) \bar{v}_2^i + \lambda (-\bar{K}_1 \bar{v}_1^i + \bar{K}_2 \bar{v}_3^i) \end{aligned} \quad (22)$$

and multiplying (22) by $g_{ij} \bar{v}_2^j$, we find

$$f(s) g_{ij} v_1^i \bar{v}_2^j = f(s) g_{ij} v_1^i v_2^j = \bar{v}_1^k \dot{\nabla}_k \lambda \quad (23)$$

or

$$0 = \bar{v}_1^k \dot{\nabla}_k \lambda \quad (24)$$

where $g_{ij} v_1^i v_2^j = 0$, $g_{ij} \bar{v}_2^i \bar{v}_2^j = 1$, $g_{ij} \bar{v}_1^i \bar{v}_2^j = 0$ and $g_{ij} \bar{v}_3^i \bar{v}_2^j = 0$.

From (24), we get λ is prolonged covariant constant [12]. On the other hand

$$f(s) = \pm \sqrt{(1 - \lambda \bar{K}_1)^2 + \lambda^2 \bar{K}_2^2}.$$

Let the angle α be between the tangent vector fields v_1 and $\bar{v}_1$ of Bertrand curve pair $(C, \bar{C})$. Since $(C, \bar{C})$ is Bertrand curve pair and $v_1 \perp v_2$, $v_1 \perp \bar{v}_2$ is obtained. Then v_1 can be written $v_1^i = a \bar{v}_1^i + b \bar{v}_3^i$ where $a = g_{ij} v_1^i \bar{v}_1^j = \cos \alpha$ and $b = g_{ij} v_1^i \bar{v}_3^j = \sin \alpha$.

$$v_3^i = \varepsilon_{ijk} v_1^j v_2^k = \varepsilon_{ijk} (\bar{v}_1^j \cos \alpha + \bar{v}_3^j \sin \alpha) \bar{v}_2^k \quad (25)$$

$$v_3^i = \bar{v}_3^i \cos \alpha - \bar{v}_1^i \sin \alpha. \quad (26)$$

From here, the following equality can be written as:

$$\begin{pmatrix} v_1^i \\ v_2^i \\ v_3^i \end{pmatrix} = \begin{pmatrix} \cos \alpha & 0 & \sin \alpha \\ 0 & 1 & 0 \\ -\sin \alpha & 0 & \cos \alpha \end{pmatrix} \begin{pmatrix} \bar{v}_1^i \\ \bar{v}_2^i \\ \bar{v}_3^i \end{pmatrix}. \quad (27)$$

Theorem 3.2 *If $(C, \bar{C})$ is a Bertrand curve pair, then there are the following relations between Bishop vector fields of C and $\bar{C}$:*

$$\bar{v}_1^i = v_1^i \cos \alpha + n_1^i \sin \alpha \sin \theta - n_2^i \sin \alpha \cos \theta \quad (28)$$

$$\begin{aligned} \bar{n}_1^i = & (-\sin \bar{\theta} \sin \alpha) v_1^i + (\cos \bar{\theta} \cos \theta + \sin \bar{\theta} \sin \theta \cos \alpha) n_1^i + \\ & + (\sin \theta \cos \bar{\theta} - \sin \bar{\theta} \cos \theta \cos \alpha) n_2^i \end{aligned} \quad (29)$$

$$\begin{aligned} \bar{n}_2^i = & (\cos \bar{\theta} \sin \alpha) v_1^i + (\sin \bar{\theta} \cos \theta - \cos \bar{\theta} \sin \theta \cos \alpha) n_1^i + \\ & + (\sin \bar{\theta} \sin \theta + \cos \bar{\theta} \cos \theta \cos \alpha) n_2^i. \end{aligned} \quad (30)$$

Proof. If $(C, \bar{C})$ is Bertrand curve pair, from (14) and (27), we have

$$\begin{pmatrix} \bar{v}_1^i \\ \bar{n}_1^i \\ \bar{n}_2^i \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \bar{\theta} & \sin \bar{\theta} \\ 0 & \sin \bar{\theta} & \cos \bar{\theta} \end{pmatrix} \begin{pmatrix} \cos \alpha & 0 & -\sin \alpha \\ 0 & 1 & 0 \\ \sin \alpha & 0 & \cos \alpha \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \theta & \sin \theta \\ 0 & -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} v_1^i \\ v_2^i \\ v_3^i \end{pmatrix} \quad (31)$$

where $\bar{\theta} = \angle(\bar{v}_2^i, \bar{n}_1^i)$.

Theorem 3.3 *If $(C, \bar{C})$ is Bertrand curve pair, the equality $\cos \alpha f(s) = 1 - \lambda \bar{K}_1$ is satisfied.*

Proof. If $(C, \bar{C})$ is Bertrand curve pair, we have $C(s) = \bar{C}(\bar{s}) + \lambda \bar{v}_2^i(\bar{s})$ where λ is prolonged covariant constant. Taking prolonged covariant derivative of this equality in the direction of $\bar{v}_1^k$, we get

$$\begin{aligned} \bar{v}_1^k \bar{\nabla}_k C &= (\bar{v}_1^k \bar{\nabla}_k C) f(s) = \bar{v}_1^k \bar{\nabla}_k \bar{C} + \lambda \bar{v}_1^k \bar{\nabla}_k \bar{v}_2^i \\ \bar{v}_1^i f(s) &= \bar{v}_1^i + \lambda (-\bar{K}_1 \bar{v}_1^i + \bar{K}_2 \bar{v}_3^i) \\ &= (1 - \lambda \bar{K}_1) \bar{v}_1^i + \lambda \bar{K}_2 (-\sin \bar{\theta} \bar{n}_1^i + \cos \bar{\theta} \bar{n}_2^i). \end{aligned} \quad (32)$$

Multiplying (32) by $g_{ij} \bar{v}_1^j$, we have $\cos \alpha f(s) = 1 - \lambda \bar{K}_1$ where $g_{ij} \bar{v}_1^i \bar{v}_1^j = \cos \alpha (\alpha = \angle(\bar{v}_1^i, \bar{v}_1^i))$, $g_{ij} \bar{n}_1^i \bar{v}_1^j = g_{ij} \bar{n}_2^i \bar{v}_1^j = 0$.

Proof 2.

$$\begin{aligned} \bar{v}_1^i f(s) &= \bar{v}_1^i + \lambda \bar{T}_k \bar{v}_1^k \bar{v}_p^i \\ &= \bar{v}_1^i + \lambda (\bar{T}_k \bar{v}_1^k \bar{v}_1^i + \bar{T}_k \bar{v}_1^k \bar{v}_3^i) \\ &= \bar{v}_1^i + \lambda (-\bar{T}_k \bar{v}_1^k \bar{v}_1^i + \bar{T}_k \bar{v}_1^k \bar{v}_3^i) \\ &= \bar{v}_1^i + \lambda (-\bar{C} \bar{v}_1^i - \bar{C} \bar{v}_3^i). \end{aligned} \quad (33)$$

$$\begin{aligned} g_{ij} \bar{v}_1^i \bar{v}_1^j f(s) &= g_{ij} \bar{v}_1^i \bar{v}_1^j - \lambda \bar{C} g_{ij} \bar{v}_1^i \bar{v}_3^j \\ \cos \alpha f(s) &= 1 - \lambda \bar{C} = 1 - \bar{K}_1. \end{aligned} \quad (34)$$

is obtained where $g_{ij} \bar{v}_1^i \bar{v}_1^j = 1$ and $g_{ij} \bar{v}_3^i \bar{v}_1^j = 0$

Theorem 3.4 *If $(C, \bar{C})$ is Bertrand curve pair, $\sin \alpha f(s) = \lambda \bar{K}_2$ is valid.*

Proof. Multiplying (32) by $g_{ij} \bar{n}_1^j$ and using Theorem 3.2, we obtain

$$\begin{aligned} g_{ij} \bar{v}_1^i \bar{n}_1^j f(s) &= -\lambda \bar{K}_2 \sin \bar{\theta} \\ -\sin \bar{\theta} \sin \alpha f(s) &= -\lambda \bar{K}_2 \sin \bar{\theta} \\ \sin \alpha f(s) &= \lambda \bar{K}_2. \end{aligned} \quad (35)$$

Proof 2. Multiplying (33) by $g_{ij} \bar{v}_3^j$ and using (27), we get

$$g_{ij} v_1^i \bar{v}_3^j f(s) = \lambda \left(-\frac{2}{31} \right) = \lambda \bar{K}_2 \quad (36)$$

and

$$\begin{pmatrix} \cos\alpha & 0 & -\sin\alpha \\ 0 & 1 & 0 \\ \sin\alpha & 0 & \cos\alpha \end{pmatrix} \begin{pmatrix} v_1^i \\ v_2^i \\ v_3^i \end{pmatrix} = \begin{pmatrix} \bar{v}_1^i \\ \bar{v}_2^i \\ \bar{v}_3^i \end{pmatrix} \quad (37)$$

and

$$\begin{aligned} g_{ij} v_1^i (\sin\alpha v_1^j + \cos\alpha v_3^j) f(s) &= \lambda \bar{K}_2 \\ \sin\alpha f(s) &= \lambda \bar{K}_2 \end{aligned} \quad (38)$$

where $g_{ij} \bar{v}_1^i \bar{v}_3^j = 0$, $g_{ij} \bar{v}_3^i \bar{v}_3^j = 1$, $g_{ij} v_1^i v_1^j = 1$ and $g_{ij} v_1^i v_3^j = 0$.

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