#### International Journal of Contemporary Mathematical Sciences Vol. 16, 2021, no. 1, 21 - 34 HIKARI Ltd, www.m-hikari.com https://doi.org/10.12988/ijcms.2021.91467

# On the Calderón-Zygmund Singular Integral Operators and Local Hardy-type Amalgam Spaces (Part 1)

#### Hon Ming Ho

Department of Mathematics
Hong Kong University of Science and Technology, Hong Kong

This article is distributed under the Creative Commons by-nc-nd Attribution License. Copyright © 2021 Hikari Ltd.

#### Abstract

In this paper, we prove that Calderón-Zygmund singular integral operators are bounded from local Hardy-type amalgam space  $l^q(h^p)$  into  $l^p(h^p)$  where  $1 < q < \infty$  and 0 .

Mathematics Subject Classification: 30E25

**Keywords:** Local Hardy-type Amalgam Spaces, Calderón-Zygmund Singular Operators

#### 1 Introduction

Amalgam spaces are defined as follows:

$$l^{q}(L^{p}) \stackrel{\text{def}}{=} \left\{ f \colon \|f\|_{l^{q}(L^{p})} = \left[ \sum_{Q \in \mathcal{Q}} \|f\|_{L^{p}(Q)}^{q} \right]^{\frac{1}{q}} < \infty \right\},$$

where  $0 < p, q < \infty$ , Q is the collection of all cubes with uniform side length 1 and their vertices on integral lattices of  $\mathbb{R}^n$ . One of the main reason to study the amalgam spaces is that they allow us to separate the global behaviour from the local behaviour of a function. In [14] the author proved that pseudo-differential operators of order zero simultaneously preserve local and

global behaviour of functions in the amalgam spaces in the sense that they are bounded from  $l^q(L^p)$  into  $l^q(L^p)$ , where  $0 < q \le 1 < p < \infty$ . In [15], the author proved that pseudo-differential operators of order zero are bounded from local Hardy-type amalgam space  $l^q(h^p)$  into  $l^q(h^p)$  where  $0 < q \le p \le 1$ . In [16], the author proved that Calderón-Zygmund singular integral operators are bounded from  $l^1(L^p)$  into  $weak - l^1(L^p)$ . In this paper, we prove that Calderón-Zygmund singular integral operators are bounded from local Hardy-type amalgam space  $l^q(h^p)$  into  $l^p(h^p)$  where  $1 < q < \infty$  and  $0 . The main results in this paper serve as a kind of preparation for proving the boundedness of Calderón-Zygmund singular integral operators from <math>l^1(h^p)$  into  $weak - l^1(h^p)$ , where 0 , which will appear in a separate paper [17].

## 2 Preliminary

Throughout this paper, we will adopt the following notations.

- (a) If  $1 \leq p \leq \infty$ , then p' is denoted as the exponent conjugate to p. Q(x,r) is the closed cube centered at x with side length r. l(Q) is the side length of the cube Q. B(x,r) is an open ball centered at x with radius r. If t > 0, B is an open ball and Q is a closed cube in  $\mathbb{R}^n$ , then tB, tQ have the same centers as B, Q respectively but whose radius and side length are expanded by the factor t. If E is a subset of  $\mathbb{R}^n$ , then the complement of E in  $\mathbb{R}^n$  will be denoted as  $^cE$ .  $m_n$  is the usual Lebesgue measure on  $\mathbb{R}^n$ .
- (b)  $L^p$  will indicate the usual function spaces on  $\mathbb{R}^n$ .  $L^p_c, C_c, C^\infty_c$  are denoted as the spaces of all compactly supported  $L^p$ -functions, continuous functions and smooth functions respectively.  $\mathcal{S}$  is the Schwartz class of all rapidly decreasing functions. Its dual space consists of all tempered distributions, denoted as  $\mathcal{S}'$ .
- (c) Let  $A_r^p f(x)$  be the average of f over the closed cube Q(x,r) in  $p^{th}$ -mean, i.e.  $A_r^p f(x) = \left[\frac{1}{r^n} \int_{Q(x,r)} |f(\omega)|^p d\omega\right]^{\frac{1}{p}}$  provided that  $0 . When <math>p = \infty$ , we define  $A_r^\infty f(x) = \|f\|_{L^\infty(Q(x,r))}$ . When r = 1, we simply write  $A^p f(x)$ .
- (d) The Hardy-Littlewood maximal function and its uncentered version are defined as

$$Mf(x) = \sup_{r>0} \frac{1}{c_n r^n} \int_{B(x,r)} |f| dm_n , \ \widetilde{M}f(x) = \sup_{x \in B(y,r)} \frac{1}{c_n r^n} \int_{B(y,r)} |f| dm_n,$$

respectively, where  $c_n$  is the volume of the unit ball.

(e)  $l^q(L^p) \stackrel{\text{def}}{=} \left\{ f \colon \|f\|_{l^q(L^p)} \stackrel{\text{def}}{=} \left[ \sum_{Q \in \mathcal{Q}} \|f\|_{L^p(Q)}^q \right]^{\frac{1}{q}} < +\infty \right\}$ , where  $0 < p, q < \infty$ ,  $\mathcal{Q}$  is the collection of closed cubes with uniform side length 1 and their vertices on integral lattices of  $\mathbb{R}^n$ . If  $q = \infty$ , then  $\|f\|_{l^\infty(L^p)} = \sup_{Q \in \mathcal{Q}} \|f\|_{L^p(Q)}$ .

The following four lemmas will be frequently used throughout the thesis.

**Lemma 2.1** Let  $f : \mathbb{R}^n \to \mathbb{C}$  be a Borel measurable function,  $\lambda, r > 0, k \in \mathbb{Z}^+, 0 . Then$ 

$$|\{\ A_r^p f > \lambda\ \}| \leqslant c\ 2^{kn}\ |\{\ A_{r\cdot 2^{-k}}^p f > 2^{-n/p}\ \lambda\ \}|,$$

where c is a constant independent of f, r and  $\lambda$ .

**Proof** The reader may refer to [14].

**Lemma 2.2 (Bertrandias, Datry, Dupuis [2])** Suppose that  $0 < p, q < \infty$ . Then there exist two numbers  $c_1, c_2$  dependent only on p, q, n such that for each  $f: \mathbb{R}^n \to \mathbb{C}$  be Borel measurable,

$$c_1 \|A^p f\|_q \leqslant \left[ \sum_{Q \in \mathcal{Q}} (\|f\|_{L^p(Q)})^q \right]^{\frac{1}{q}} \leqslant c_2 \|A^p f\|_q.$$

**Lemma 2.3** Let  $f: \mathbb{R}^n \to \mathbb{C}$  be a Borel measurable function,  $r > 0, k \in \mathbb{Z}^+, 0 < p, q < \infty$ . Then

$$||A_r^p f||_q \leqslant c(2^{kn})^{\frac{1}{q}} ||A_{r\cdot 2^{-k}}^p f||_q$$

where c is a constant independent of f and r but dependent on p, q.

**Proof** It is a consequence of Lemma 2.1 or 2.2. The proof is complete.

**Lemma 2.4 (Maximal Theorem)** The uncentered maximal operator  $\widetilde{M}$  is bounded from  $L^1$  into weak- $L^1$  and from  $L^p$  into  $L^p$  for each 1 .

**Proof** For the proof, the reader can refer to [21].

The class of operators which we will consider in this thesis is defined as follows. We call operators in this class to be the Calderón-Zygmund singular integral operators throughout the thesis.

Definition 2.5 (The Calderón-Zygmund Singular Integral Operators) Let T be a bounded linear operator on  $L^2$  that are translation-invariant. It is a well-known fact (see Stein [21]) that T is representable as

$$\widehat{Tf}(\xi) = m(\xi)f(\xi)$$
 and  $|m(\xi)| \leq A$ .

for all  $f \in L^2$ . Now let W be the tempered distribution with  $\hat{W} = m$ . Then we have, for all  $f \in \mathcal{S}$ ,

$$Tf = W * f.$$

We make the a priori assumption that the distribution W coincides, away from the origin, with a locally integrable function K. Then we have

$$Tf(x) = \int K(x-y)f(y)dy,$$

for all  $f \in L^2_c$  and almost all  $x \notin supp(f)$ . Moreover we assume that

$$|K(x-y) - K(x-y')| \le c \frac{|y-y'|^{\delta}}{|x-y|^{n+\delta}} \text{ for all } |x-y| \ge 2|y-y'|,$$

where  $\delta \in (0,1]$  and c > 0 are fixed constants. Then we call the operator T to be a Calderón-Zygmund singular integral operator.

Such operator was known to be bounded on  $L^p(1 and from <math>L^1$  into weak- $L^1$  (see Stein [21]). The typical examples are the Hilbert transform on  $\mathbb{R}^1$ , whose kernel is  $\frac{1}{x}$ , and its n-dimensional analogue, the Riesz transforms on  $\mathbb{R}^n$ . The kernel of the  $j^{th}$  Riesz transform is  $\frac{y_j}{|y|^{n+1}}$ , where  $1 \le j \le n$ .

# 3 Local Hardy-Type Amalgam Spaces $l^q(h^p)$ , $1 \le q < \infty$ , 0

### **3.1** Definition of $l^q(h^p)$ , $1 \leqslant q < \infty$ , 0

Definition 3.1.1 Suppose  $f \in \mathcal{S}'$ ,  $x \in \mathbb{R}^n$ .

(a) Let 
$$F_N = \{ \psi \in C_0^{\infty}(B(0,1)) : ||\partial^{\beta}\psi||_{\infty} \leq 1, |\beta| \leq N \},$$

(i) 
$$M_1^{F_N} f(x) = \sup_{0 < t < 1} \sup_{\psi \in F_N} |f * \psi_t(x)|,$$

(ii) 
$$M^{F_N} f(x) = \sup_{t>0} \sup_{\psi \in F_N} |f * \psi_t(x)|.$$

(b) Let  $\phi \in \mathcal{S}$  with  $\int \phi \neq 0$ .

(i) 
$$M_1^{\phi} f(x) = \sup_{0 < t < 1} |f * \phi_t(x)|,$$

(ii) 
$$M^{\phi} f(x) = \sup_{t>0} |f * \phi_t(x)|.$$

**Definition 3.1.2** Let  $1 \leq q < \infty$ ,  $0 . We denote <math>l^q(h^p) = \{f \text{ is a distribution}: A^p(M_1^{F_N}f) \in L^q\}$ , where  $F_N$  is the collection of smooth functions defined in Definition 3.1.1 and  $N = \max\{0, [n(\frac{1}{p}-1)]+1\}, [n(\frac{1}{p}-1)] \text{ is the integral part of } n(\frac{1}{p}-1).$ 

Throughout this section, N will be used to denote the number  $\max\{0, [n(\frac{1}{p}-1)]+1\}$ . We remark that with this choice of N, it is not difficult to verify that for each local (2,p)-atom a (see definition in Goldberg [12]),  $\|M_1^{F_N}a\|_{L^p} \leq c$  where c is a constant independent of the choice of a. As a consequence, for  $f \in h^p$ ,  $\|M_1^{F_N}f\|_{L^p} \simeq \|f\|_{h^p}$ . Besides, on this range of norm indices p,q, the basic assumption on f is that f is a distribution which, by definition, is a continuous linear functional on  $C_c^{\infty}$ . We recall a characterization of distributions. A linear functional f on  $C_c^{\infty}$  is a distribution if and only if for each compact set  $K \subset \mathbb{R}^n$ , there exist C > 0 and  $m \in \mathbb{Z}^+$  such that

$$|f(\phi)| \leq C \sum_{|\alpha| \leq m} \|\partial^{\alpha} \phi\|_{\infty}$$
 whenever  $f \in C^{\infty}$  supported inside  $K$ .

**Lemma 3.1.3** Suppose that (i)  $\phi \in C_0^{\infty}(B(0,1))$ , for each  $|\alpha| \leq N$ ,  $\|\partial^{\alpha}\phi\|_{\infty} \leq 1$ , (ii)  $\eta \in C_0^{\infty}(Q_{\eta})$  where  $Q_{\eta}$  is a cube with  $l(Q_{\eta}) \geq 1$ , for each  $|\alpha| \leq N$ ,  $\|\partial^{\alpha}\eta\|_{\infty} \leq 1$ . Then there exists a constant c (dependent only on N) such that

$$M_1^{\phi}(f\eta)(x) \leqslant c \cdot \chi_{3Q_{\eta}}(x) M_1^{F_N} f(x), \quad \text{whenever } f \in \mathcal{S}', x \in \mathbf{IR}^{\mathbf{n}},$$

where by definition,  $(f\eta)(\varphi) := f(\eta\varphi)$  for all  $\varphi \in \mathcal{S}$ .

**Proof** Noticing that for all  $x, \omega \in \mathbb{R}^n$ ,  $t \in (0, 1)$ ,

$$\eta(\omega)\phi_t(x-\omega) = \eta_{x,t}(\frac{x-\omega}{t})\frac{1}{t^n}\phi(\frac{x-\omega}{t}) \quad \text{where } \eta_{x,t}(\xi) = \eta(x-t\xi),$$

$$= \frac{1}{t^n}(\eta_{x,t} \cdot \phi)(\frac{x-\omega}{t}).$$

For all multi-index  $|\beta| \leq N, \ \omega \in {\rm I\!R}^{\bf n}, \ t \in (0,1),$ 

$$\partial_{\omega}^{\beta}(\eta_{x,t} \cdot \phi)(\omega) = \sum_{\gamma \leqslant \beta} {\beta \choose \gamma} (\partial_{\omega}^{\gamma} \eta_{x,t})(\omega) (\partial_{\omega}^{\beta-\gamma} \phi)(\omega),$$
$$= \sum_{\gamma \leqslant \beta} {\beta \choose \gamma} (-t)^{|\gamma|} (\partial_{\omega}^{\gamma} \eta)(x - t\omega) (\partial_{\omega}^{\beta-\gamma} \phi)(\omega).$$

It follows that  $\|\partial_{\omega}^{\beta}(\eta_{x,t}\cdot\phi)\|_{\infty} \leqslant \sum_{\gamma\leqslant\beta} {\beta\choose\gamma} \leqslant c_N$  because  $\operatorname{supp}(\phi)\subseteq B(0,1), \ \forall |\gamma|\leqslant N, \ \|\partial^{\gamma}\eta\|_{\infty}, \ \|\partial^{\gamma}\phi\|_{\infty}\leqslant 1 \ \text{and} \ 0< t<1.$  So  $\frac{1}{c_N}(\eta_{x,t}\cdot\phi)\in F_N$ . Moreover  $\operatorname{supp}(M_1^{\phi}(f\eta))\subseteq 3Q_{\eta}$ . Consequently, for all  $x\in\mathbf{R}^{\mathbf{n}},\ M_1^{\phi}(f\eta)(x)\leqslant c_N\cdot\chi_{3Q_{\eta}}(x)(M_1^{F_N}f)(x)$ . The proof of the lemma is complete.

We first start with a decomposition theorem for distributions in  $l^q(h^p)$ , in which the building blocks are required to be localized in some dyadic cubes.

As a consequence, we show that  $C_c^{\infty}$  is a dense subspace of  $l^q(h^p)$ , and each distribution in  $l^q(h^p)$  is locally a  $h^p$ -distribution.

To state the lemma below, we introduce a smooth partition of unity as below: Let  $\mathcal{Q}$  be a prescribed collection of closed cubes with uniform side length 1 whose interiors do not intersect with each other and union is equal to  $\mathbb{R}^n$ . Denote  $x_k = \text{centre of } Q_k, \ Q_k \in \mathcal{Q}, \ k = 1, 2, 3, \dots$  Take a  $C^{\infty}$ -function  $\psi$  such that  $0 \leq \psi \leq 1$  on  $\mathbb{R}^n$ ,  $\psi \equiv 1$  on  $\overline{Q(0,1)}$  and  $\sup(\psi) \subseteq \overline{Q(0,2)}$ . Define  $\eta_k(x) = \frac{\psi(x-x_k)}{\sum\limits_{j=1}^{\infty} \psi(x-x_j)}$  for each  $x \in \mathbb{R}^n$ ,  $k = 1, 2, \dots$  Then for each

 $x \in \mathbf{IR}^{\mathbf{n}}, \ \sum_{j=1}^{\infty} \psi(x - x_j) \geqslant 1.$  Moreover, for each  $\alpha, k$ ,  $\|\partial^{\alpha} \eta_k\|_{\infty} \leqslant c_{\alpha} \|\partial^{\alpha} \psi\|_{\infty}$  where  $c_{\alpha}$  is dependent on  $\alpha$  and the dimension n only.

**Lemma 3.1.4** Let  $1 \le q < \infty$ ,  $0 . Suppose that <math>\{\eta_k\}_1^{\infty}$  is the smooth partition of unity stated in the previous paragraph. If f is a distribution, then the following three quantities are mutually comparable with bounds independent of f:

- (i)  $||A^p(M_1^{F_N}f)||_{L^q}$ ,
- (ii)  $\left[\sum_{k=1}^{\infty} \|f\eta_k\|_{h^p}^q\right]^{\frac{1}{q}}$ ,
- (iii)  $\inf\{\left[\sum |\lambda_Q|^q\right]^{\frac{1}{q}}: f = \sum_{Q \in \mathcal{Q}} \lambda_Q a_Q, \|a_Q\|_{h^p} \leq 1, \sup\{a_Q\} \subseteq 4Q\}, \text{ where the convergence is taken in the sense of distribution.}$

**Proof** Denote  $N_1(f), N_2(f), N_3(f)$  to be the expressions in (i), (ii), (iii) respectively. For the proof of  $N_2(f) \leqslant cN_1(f)$  where c being independent of the choice of f, it is already contained in the proof of Theorem 3.3.2 in [15]. We will not repeat it here. Besides, it is trivial to see that  $N_3(f) \leqslant N_2(f)$ . What we are going to show is that  $N_1(f) \leqslant cN_3(f)$ . Now let  $f = \sum_{Q \in \mathcal{Q}} \lambda_Q a_Q$  where  $a_Q$  's are satisfied the above mentioned properties. Then case 1:  $1 \leqslant p \leqslant q$ . We have

$$\int_{\mathbb{R}^{n}} |A^{p}(M_{1}^{F_{N}}f)(y)|^{q} dy = \sum_{V \in \mathcal{Q}} \int_{V} |A^{p}(M_{1}^{F_{N}}f)(y)|^{q} dy,$$

$$\leq \sum_{V \in \mathcal{Q}} \int_{V} |\sum_{Q \in \mathcal{Q}} A^{p}(M_{1}^{F_{N}}(\lambda_{Q}a_{Q}))(y)|^{q} dy,$$

$$= \sum_{V \in \mathcal{Q}} \int_{V} |\sum_{4Q \cap 4V \neq \emptyset} A^{p}(M_{1}^{F_{N}}(\lambda_{Q}a_{Q}))(y)|^{q} dy,$$

$$\leq \sum_{V \in \mathcal{Q}} (\sum_{4Q \cap 4V \neq \emptyset} ||\lambda_{Q}a_{Q}||_{h^{p}})^{q},$$

$$\leq c_{n,q} \sum_{V \in \mathcal{Q}} |\lambda_{V}|^{q},$$

where  $c_{n,q}$  only depends on the dimension n and norm q. Case 2: 0 . We have

$$\begin{split} \int_{\mathbf{R}^{\mathbf{n}}} |A^{p}(M_{1}^{F_{N}}f)(y)|^{q} dy &= \sum_{V \in \mathcal{Q}} \int_{V} |A^{p}(M_{1}^{F_{N}}f)(y)|^{q} dy, \\ &\leqslant \sum_{V \in \mathcal{Q}} \int_{V} |\sum_{Q \in \mathcal{Q}} |\lambda_{Q}|^{p} \int_{Q(y,1)} |M_{1}^{F_{N}}a_{Q}|^{p} dm_{n} |^{\frac{q}{p}} dy, \\ &= \sum_{V \in \mathcal{Q}} \int_{V} |\sum_{4Q \cap 4V \neq \emptyset} |\lambda_{Q}|^{p} \int_{Q(y,1)} |M_{1}^{F_{N}}a_{Q}|^{p} dm_{n} |^{\frac{q}{p}} dy, \\ &\leqslant \sum_{V \in \mathcal{Q}} (\sum_{4Q \cap 4V \neq \emptyset} |\lambda_{Q}|^{p})^{\frac{q}{p}}, \\ &\leqslant \sum_{V \in \mathcal{Q}} (m^{\frac{q}{p}-1} \sum_{4V \cap 4Q \neq \emptyset} |\lambda_{Q}|^{q}), \\ &\leqslant m^{\frac{q}{p}} \sum_{V \in \mathcal{Q}} |\lambda_{V}|^{q}, \end{split}$$

where m is the number of all distinct elements in  $\{Q \in \mathcal{Q} : 4Q \cap 4V \neq \emptyset\}$ . Note that m is independent of the chioce of V. The second last inequality follows from the Hölder's inequality and the fact that  $p \leqslant q$ . This completes the proof of the lemma.

As a consequence of the above lemma, each distribution in  $l^q(h^p)$   $(1 \le p \le q < \infty)$  is locally an  $L^p$ -function because each  $f\eta_k \in L^p$ . As a result, it is easy to see that  $l^q(h^p)$  is a Banach space. Moreover since  $||f\eta_k||_{h^p} \simeq ||f\eta_k||_{L^p}$  where  $1 , <math>||A^p(M_1^{F_N}f)||_q \simeq ||A^pf||_q$ . For  $0 , <math>(l^q(h^p), ||\cdot||_{l^q(h^p)})$  defines a topological vector space with a complete translation invariant metric,  $d(f,g) = ||f-g||_{l^q(h^p)}^p$ . Besides, we can show that  $C_c^\infty$  is a dense subspace of  $l^q(h^p)$ . Indeed, it is obvious that  $||f-\sum_{k=1}^N f\eta_k||_{l^q(h^p)} \to 0$  as  $N \to \infty$  by the previous lemma. Together with the fact that  $h^p$  space is stable under multiplication by S (see Goldberg [12]), each  $f\eta_k$  can be approximated in  $h^p$ -norm by smooth functions which are compactly supported inside  $4Q_k$  (supp $(f\eta_k) \subset 2Q_k$ ). Finally by the previous lemma, f can be approximated in  $l^q(h^p)$ -norm by compactly supported smooth functions.

In the rest of this section, we consider  $A_r^p f(x)$  as the average of f over the open ball B(x,r).

## 3.2 The Calderón-Zygmund Singular Integral Operators on $l^q(h^p)$ , $1 < q < \infty$ , 0

**Proposition 3.2.1** Let 0 . Then there exists a constant <math>c > 0 (dependent only on p) such that

$$\sup_{\lambda>0} \lambda |\{A^{p}(M^{F_{N}}f) > \lambda\}| \leq c \|A^{p}(M_{1}^{F_{N}}f)\|_{L^{1}}, \quad whenever \quad f \in l^{1}(h^{p}).$$

**Proof** Let  $f \in l^1(h^p)$ . Denote  $Gf = \sup_{\psi \in F_N} \sup_{t \geqslant 1} |f * \psi_t|$ . Then  $A^p(M^{F_N}f) \leqslant cA^p(M_1^{F_N}f) + cA^p(Gf)$ . So

$$|\{A^p(M^{F_N}f) > \lambda\}| \le |\{A^p(M_1^{F_N}f) > \frac{\lambda}{2c}\}| + |\{A^p(Gf) > \frac{\lambda}{2c}\}|.$$

Therefore, it is sufficient for us to show that the third quantity is less than  $c' \lambda^{-1} ||f||_{l^1(h^p)}$ . To finish the proof, it suffices to show that

$$||Gf||_{C(B(y,1))} \leqslant c \widetilde{M}(A^p(M_1^{F_N}f))(y), \text{ for each } y \in \mathbb{R}^n,$$
 (3.1)

where c is independent of y, f. Indeed, we fix  $y_0 \in \mathbb{R}^n$ ,  $x_0 \in B(y_0, 1)$ . Then take a function  $\phi \in C_0^{\infty}(B(0, \frac{1}{4}))$  with  $\int \phi = 1$ . (The choice of  $\phi$  is independent of f) Fix  $t \ge 1$  and  $\psi \in F_N$ . Then

$$f * \psi_t(x_0) = \frac{1}{t^n} f(\psi(\frac{x_0 - \cdot}{t})) = \frac{1}{t^n} \int_{B(x_0, t + \frac{1}{2})} f(\psi(\frac{x_0 - \cdot}{t}) \phi(w - \cdot)) dw.$$

Now for each  $w \in B(x_0, t + \frac{1}{2}), \beta \in B(w, \frac{1}{4})$ , we define

$$g_{w,\beta}(y) = 2^{-n}\psi(\frac{x_0 - \beta + 2^{-1}y}{t})\phi(w - \beta + 2^{-1}y).$$

So  $f(\psi(\frac{x_0-\cdot}{t})\phi(w-\cdot)) = f(2^n \ g_{w,\beta}(2\ (\beta-\cdot)))$ . We claim that for such  $g_{w,\beta}$  function,  $\operatorname{supp}(g_{w,\beta}) \subseteq B(0,1)$ . It follows from the fact that  $\operatorname{supp}(\phi) \subseteq B(0,\frac{1}{4})$  and  $|w-\beta|<\frac{1}{4}$ . Moreover by the Leibnitz' formula, there exists a positive number L (independent of  $x_0,\beta,w$  and t, but dependent on  $\phi$  and N) such that for each  $|\alpha| \leq N$ ,  $\|\partial_y^{\alpha}(g_{w,\beta})\|_{\infty} \leq L$ . So

$$|f(\psi(\frac{x_0-\cdot}{t})\phi(w-\cdot))| \leq L\cdot (M_1^{F_N}f)(\beta), \text{ for each } \beta\in B(w,\frac{1}{4}).$$

Integrating both sides in the  $p^{th}$ -mean over the ball  $B(w, \frac{1}{4})$ , we have

$$|f(\psi(\frac{x_0-\cdot}{t})\phi(w-\cdot))| \leqslant cL \cdot A_{\frac{1}{4}}^p(M_1^{F_N}f)(w).$$

Since  $t \ge 1$  and  $y_0 \in B(x_0, t+1)$ , we finally have

$$|f * \psi_t(x_0)| \leq c \ \tilde{M}(A^p(M_1^{F_N}f))(y_0).$$

So (3.1) is established. The proof is therefore complete.

**Theorem 3.2.2 (Goldberg [12])** Let  $0 with <math>\int \psi = 1$  and  $\int x^{\alpha} \psi(x) dx = 0$  for each  $\alpha \ne 0$ . Then for each  $f \in h^p, f - f * \psi \in H^p$  and  $\|f - f * \psi\|_{H^p} \le c \|f\|_{h^p}$  where c is independent of f. Moreover,  $f * \psi = \sum \lambda_i b_i$  a.e. on  $\mathbb{R}^n$  where  $b_i$ 's are local (2, p)-atoms with their supporting cubes,  $l(Q) \ge 1$ , and  $\sum |\lambda_i|^p \le c \|f\|_{h^p}^p$ .

**Theorem 3.2.3 (Stein [22])** Let 0 , <math>T be the Calderón-Zygmund singular integral operator with its kernel satisfying the following assumption

$$|\partial^{\beta}K(x)| \leqslant c |x|^{-n-|\beta|}$$
 for  $|\beta| \leqslant [n(\frac{1}{p}-1)], x \neq 0.$ 

Then T is bounded from  $H^p$  into  $H^p$ .

Now we come to the main results in this paper.

**Theorem 3.2.4** Let  $1 < q < \infty$ , 0 , <math>T be the Calderón-Zygmund singular integral operator. (i) If 1 , then there exists a constant <math>c (dependent only on p, q, T) such that

$$||Tf||_{l^q(h^p)} \leqslant c ||f||_{l^q(h^p)}$$
 whenever  $f \in C_c^{\infty}$ .

So T admits a unique bounded linear extension on  $l^q(h^p)$ . (ii) If 0 and K satisfies the following stronger assumption

$$|\partial^{\beta}K(x)| \leqslant c |x|^{-n-|\beta|}$$
 for  $|\beta| \leqslant [n(\frac{1}{p}-1)] + 2, \ x \neq 0,$ 

then there exists a constant c' (dependent only on p, q, T) such that

$$||Tf||_{l^q(h^p)} \leqslant c' ||f||_{l^q(h^p)}$$
 whenever  $f \in C_c^{\infty}$ .

So T admits a unique bounded linear extension on  $l^q(h^p)$ .

**Proof** Take a smooth function  $\psi$  vanishing outside B(0,1) and  $\int \psi dx = 1$ . Let  $f \in C_c^{\infty}$ . Then  $T(f * \psi)$  is a well-defined function in  $L^q$ . We now write  $Tf = T(f * \psi) + T(f - f * \psi)$ . We claim that  $\|T(f * \psi)\|_{l^q(h^p)} \leqslant c_q \|f\|_{l^q(h^p)}$ . Indeed, from (3.1), we know  $|f * \psi| \leqslant c \widetilde{M}(A^p(M_1^{F_N}f))$ . So by the maximal theorem, the  $L^q$ -boundedness property of T and the Jensen inequality  $(1 \leqslant \frac{q}{p})$ , we have  $\|T(f * \psi)\|_{l^q(h^p)} \leqslant c_q \|T(f * \psi)\|_{L^q} \leqslant c \|f * \psi\|_{L^q} \leqslant c \|\widetilde{M}(A^p(M_1^{F_N}f))\|_{L^q} \leqslant c \|f\|_{l^q(h^p)}$ .

What remains to show is  $||T(f-f*\psi)||_{l^q(h^p)} \leq c ||f||_{l^q(h^p)}$ . For convenience, we write  $\widetilde{T}f = T(f-f*\psi)$ . Take another smooth function  $\eta \in C_0^{\infty}(B(0,5))$  such that  $\eta \equiv 1$  on  $B(0,4), 0 \leq \eta \leq 1$ . Denote  $\eta_x(\cdot) = \eta(x-\cdot)$  for each

 $x \in \mathbb{R}^n$ . We are going to establish the inequality by dividing into three steps to show that for each  $x \in \mathbb{R}^n$ ,

$$A^{p}(M_{1}^{F_{N}}(\widetilde{T}(f\eta_{x})))(x) \leqslant c A_{10}^{p}(M_{1}^{F_{N}}f)(x), \tag{3.2}$$

$$A^{p}(M_{1}^{F_{N}}(\widetilde{T}(f(1-\eta_{x}))))(x) \leqslant c M_{1}'f(x) \qquad \text{if } 1 
$$A^{p}(M_{1}^{F_{N}}(\widetilde{T}(f(1-\eta_{x}))))(x) \leqslant c \|Gf\|_{C(B(x,1))} \qquad \text{if } 0$$$$

$$A^{p}(M_{1}^{F_{N}}(\widetilde{T}(f(1-\eta_{x}))))(x) \leqslant c \|Gf\|_{C(B(x,1))} \quad \text{if } 0 (3.4)$$

where  $M'_1f = \sup_{1 < t} \frac{1}{c_n t^n} \int_{B(x,t)} |f| dm_n$  and Gf is defined as in Proposition 3.2.1. We fix a point  $x_0 \in \mathbb{R}^n$ 

**Step (i).** We are going to establish (3.2). We argue that  $f\eta_{x_0} \in h^p$ . In fact, by Lemma 3.1.3,  $||f\eta_{x_0}||_{h^p} \leq c \ A_{10}^p(M_1^{F_N}f)(x_0) \leq c \ A_{20}^p(M_1^{F_N}f)(\beta)$  for each  $\beta \in B(x_0, 10)$ . Then  $f\eta_{x_0} \in h^p$  follows from Lemma 2.3. Case 1: 1q. The L<sup>p</sup>-boundedness property of T implies that  $A^p(M_1^{F_N}(\widetilde{T}(f\eta_{x_0})))(x_0) \leq$  $\|\widetilde{T}(f\eta_{x_0})\|_{L^p} \leqslant c \|f\eta_{x_0}\|_{L^p} \simeq c \|f\eta_{x_0}\|_{h^p} \leqslant c A_{10}^p(M_1^{F_N}f)(x_0).$  Case 2: 0 <  $p \leqslant 1$ . By invoking Theorem 3.2.2,  $f\eta_{x_0} = g_1 + g_2$  where  $g_1 \in H^p$  with  $||g_1||_{H^p} \leqslant c||f\eta_{x_0}||_{h^p}$  and  $g_2 = \sum \lambda_i b_i$  a.e. on  $\mathbb{R}^n$  for some local (2,p)-atoms  $b_i$ 's with their supporting cubes  $l(Q) \geqslant 1$  and  $\sum |\lambda_i|^p \leqslant c \|f\eta_{x_0}\|_{h^p}^p$ . First we claim that for each  $i, A^p(M_1^{F_N}(\widetilde{T}b_i))(x_0) \leq c$  where c is independent of i and f. Let  $Q_i$  be the smallest supporting cube of  $b_i$ ,  $l(Q_i) \ge 1$ . Then by the Hölder's inequality,

$$|A^{p}(M_{1}^{F_{N}}(\widetilde{T}b_{i}))(x_{0})|^{p} \leq c \left[ \int_{B(x_{0},1)} |M_{1}^{F_{N}}(\widetilde{T}b_{i})(y)|^{p\cdot\frac{2}{p}} dy \right]^{\frac{p}{2}},$$

$$\leq c ||M_{1}^{F_{N}}(\widetilde{T}b_{i})||_{L^{2}}^{p},$$

$$\leq c ||b_{i}||_{L^{2}}^{p},$$

$$\leq c |Q_{i}|^{\frac{p}{2}-1},$$

$$\leq c.$$

The claim is justisfied. Now we come back to establish (3.2). By Theorem 3.2.3,

$$|A^{p}(M_{1}^{F_{N}}(\widetilde{T}(f\eta_{x_{0}})))(x_{0})|^{p} \leq |A^{p}(M_{1}^{F_{N}}(\widetilde{T}g_{1}))(x_{0})|^{p} + |A^{p}(M_{1}^{F_{N}}(\widetilde{T}(\sum \lambda_{i}b_{i})))(x_{0})|^{p},$$

$$\leq ||\widetilde{T}g_{1}||_{H^{p}}^{p} + \sum |\lambda_{i}|^{p} |A^{p}(M_{1}^{F_{N}}(\widetilde{T}b_{i}))(x_{0})|^{p},$$

$$\leq c ||g_{1}||_{H^{p}}^{p} + c ||f\eta_{x_{0}}||_{h^{p}}^{p},$$

$$\leq c ||f\eta_{x_{0}}||_{h^{p}}^{p}.$$

So (3.2) follows. (It is this point where we make use of Theorem 3.2.3.)

**Step (ii).** We are going to establish (3.3). Indeed, for each  $\beta \in B(x_0, 2)$ ,

$$\begin{split} |\widetilde{T}(f(1-\eta_{x_0}))(\beta)| &= |\int_{\mathbf{R}^{\mathbf{n}}} (K(\beta-y) - K * \psi(\beta-y)) \cdot f(y) \cdot (1-\eta_{x_0}(y)) dy|, \\ &\leqslant \int_{c_{B(x_0,4)}} |K(\beta-y) - K * \psi(\beta-y)| |f(y)| dy, \\ &\leqslant \int_{c_{B(x_0,4)}} (\int_{\mathbf{R}^{\mathbf{n}}} |K(\beta-y) - K(\beta-y-t)| |\psi(t)| dt) \cdot |f(y)| dy, \\ &\leqslant c \|\psi\|_1 \cdot \int_{c_{B(x_0,4)}} \frac{|f(y)|}{|\beta-y|^{n+\delta}} dy \quad \text{(since supp } (\psi) \subseteq B(0,1)), \\ &\leqslant c \|\psi\|_1 \cdot \int_{c_{B(x_0,4)}} \frac{|f(y)|}{|x_0-y|^{n+\delta}} dy \quad \text{(since } \beta \in B(x_0,2)), \\ &\leqslant c \|\psi\|_1 \cdot \sum_{j=0}^{\infty} (\frac{1}{2})^{(n+\delta)(j+2)} \int_{2^{j+2} \leqslant |x_0-y| < 2^{j+3}} |f(y)| dy, \\ &\leqslant c \|\psi\|_1 \cdot (M_1' f)(x_0). \end{split}$$

Step (iii). We are going to establish 3.4. For convenience, we write  $\Theta(\cdot) \stackrel{def.}{=} \eta(\frac{\cdot}{2}) - \eta(\cdot)$ . Noticing that for each  $\beta \in B(x_0, 2)$ ,

$$\widetilde{T}(f(1-\eta_{x_0}))(\beta) = \int_{\mathbb{R}^n} \left[ K(\beta-y) - K * \psi(\beta-y) \right] f(y) (1-\eta_{x_0}(y)) dy, 
= \sum_{i=0}^{\infty} \int_{\mathbb{R}^n} \left[ K(\beta-y) - K * \psi(\beta-y) \right] f(y) \Theta(\frac{x_0-y}{2^i}) dy, 
= \sum_{i=0}^{\infty} \int_{\mathbb{R}^n} \left[ K_i(\beta-y) - K_i * \psi(\beta-y) \right] f(y) \Theta(\frac{x_0-y}{2^i}) dy,$$

where  $K_i(\omega) = K(\omega)\zeta(\frac{\omega}{2^i})$  and  $\zeta$  is smooth function such that  $\zeta(y) \equiv 1$  on  $1 \leqslant |y| \leqslant 2^4$  and  $\zeta(y) \equiv 0$  outside  $2^{-1} \leqslant |y| \leqslant 2^5$ . The last equality follows from the fact that  $\Theta(\frac{x_0-\cdot}{2^i})$  vanishes outside  $2^{i+2} \leqslant |x_0-\cdot| \leqslant 5 \cdot 2^{i+1}$ , on which we must have  $2^i + 1 \leqslant |\beta - \cdot| \leqslant 2^{i+4} - 1$ . Now let  $\phi \in F_N, \omega \in B(x_0, 1), 0 < t \leqslant 1$ . Define for each  $y \in \mathbb{R}^n$ ,

$$g_{i,\omega}(y) = 2^{n(i+4)} \left( (K_i - K_i * \psi) * \phi_t \right) \left( 2^{i+4} y \right) \left[ \eta \left( \frac{x_0 - \omega}{2^{i+1}} + 8y \right) - \eta \left( \frac{x_0 - \omega}{2^i} + 16y \right) \right].$$

It follows that

$$\phi_t * \widetilde{T}(f(1 - \eta_{x_0}))(\omega) = \sum_{i=0}^{\infty} \int_{\mathbf{R}^n} \left(\frac{1}{2^{i+4}}\right)^n g_{i,\omega}\left(\frac{\omega - y}{2^{i+4}}\right) f(y) dy.$$

What we need to show is that there exists a positive number L (independent of  $i, \omega, x_0, f, \phi$  and t) such that for each  $i = 0, 1, 2, \ldots, 2^i L$   $g_{i,\omega} \in F_N$ . Indeed,

it is easy to see that  $\operatorname{supp}(g_{i,\omega}) \subseteq B(0,1)$ . Moreover the smoothness of  $\phi$  implies the function  $g_{i,\omega}$  to be also smooth. Observe that for  $|\gamma| \leqslant [n(\frac{1}{p}-1)] + 2$ ,  $\|\partial^{\gamma} K_i\|_{\infty} \leqslant C_{\gamma} 2^{-i(n+|\gamma|)}$  where  $C_{\gamma}$  depends only on  $\zeta, K, N$ , but does not depend on  $i, \omega, x_0, f, \phi, t$ . Then  $g_{i,\omega}$  satisfies the following differential inequality: for each  $|\alpha| \leqslant N$ , (let  $c_{N,\eta} = \max\{2 \cdot 16^N \|\partial^{\beta} \eta\|_{\infty} : |\beta| \leqslant N\}$ )

$$|\partial_{y}^{\alpha}(g_{i,\omega})(y)| \leq 2^{n(i+4)} \cdot c_{N,\eta} \sum_{\gamma \leq \alpha} {\alpha \choose \gamma} 2^{(i+4)|\gamma|} | [\partial_{y}^{\gamma} K_{i} - (\partial_{y}^{\gamma} K_{i}) * \psi] * \phi_{t}(2^{i+4}y) |,$$

$$\leq c 2^{n(i+4)} \cdot c_{N,\eta} \sum_{\gamma \leq \alpha} {\alpha \choose \gamma} 2^{(i+4)|\gamma|} | ||\partial_{y}^{\gamma} K_{i} - (\partial_{y}^{\gamma} K_{i}) * \psi||_{\infty},$$

$$\leq c 2^{n(i+4)} \cdot c_{N,\eta} \sum_{\gamma \leq \alpha} {\alpha \choose \gamma} 2^{(i+4)|\gamma|} \sum_{|\beta|=1} ||\partial_{y}^{\gamma+\beta} K_{i}||_{\infty},$$

$$\leq c 2^{n(i+4)} \cdot c_{N,\eta} \sum_{\gamma \leq \alpha} {\alpha \choose \gamma} 2^{(i+4)|\gamma|} \sum_{|\beta|=1} C_{\gamma+\beta} \cdot 2^{-i(n+|\gamma|+1)},$$

$$= c 2^{-i} \cdot 2^{4(n+N)} \cdot c_{N,\eta} \cdot \sum_{\gamma \leq \alpha} {\alpha \choose \gamma} \sum_{|\beta|=1} C_{\gamma+\beta},$$

$$\leq 2^{-i} L^{-1},$$

where L depends only on  $N, \eta, K$  but does not depend on  $i, \omega, x_0, f, \phi, t$ . The claim is justified. So we have

$$M_1^{F_N}(\widetilde{T}(f(1-\eta_{x_0})))(\omega) \leqslant \sum_{i=0}^{\infty} \frac{1}{2^i L} Gf(\omega) \leqslant c \|Gf\|_{C(B(x_0,1))}.$$

Then (3.4) follows.

By invoking (3.1), (3.3) and (3.4) imply that  $A^p(M_1^{F_N}(\widetilde{T}(f(1-\eta_{x_0}))))(x_0) \leq c \widetilde{M}(A^p(M_1^{F_N}f))(x_0)$ . Finally we have

$$\|\widetilde{T}f\|_{l^q(h^p)} \leqslant c \|A_{10}^p(M_1^{F_N}f)\|_{L^q} + c \|\widetilde{M}(A^p(M_1^{F_N}f))\|_{L^q} \leqslant c \|f\|_{l^q(h^p)}.$$

The proof is therefore complete.

#### References

[1] J. Alvarez, Continuity of Calderon-Zygmund Type Operators on the Predual of a Morrey Space, Clifford algebras in analysis and related topics (Fayetteville, AR, 1993), Stud. Adv. Math., CRC, Boca Raton, Fl, 1996, 309-319. https://doi.org/10.4324/9781315139548-15

- [2] J. P. Bertrandias, C. Datry, C. Dupuis, Unions et intersections d'espaces  $L^p$  invariants par translation ou convolution, Ann. Inst. Fourier (Grenoble), **28** (1978), 53-84. https://doi.org/10.5802/aif.689
- [3] C. Carton-Lebrun, H. P. Heinig, S. C. Hofmann, Integral Operators on Weighted Amalgams, Studia Mathematica, 109 (2) (1994), 133-157. https://doi.org/10.4064/sm-109-2-133-157
- [4] S. Chanillo, Weighted Norm Inequalities For Strongly Singular Convolution Operators, Tran. Amer. Math. Soc., 281 (1) (1984), 77. https://doi.org/10.1090/s0002-9947-1984-0719660-6
- [5] M. Cowling, S. Meda, R. Pasquale, Riesz Potential and Amalgams, Ann. Inst. Fourier, Grenoble, 49 (4) (1999), 1345-1367. https://doi.org/10.5802/aif.1720
- [6] R. R. Coifman, A real variable characterization of  $H^p$ , Studia Math., 51 (1974), 269-274. https://doi.org/10.4064/sm-51-3-269-274
- [7] R. R. Coifman, Characterizations of Fourier transforms of Hardy spaces, *Proc. Nat. Acad. Sci. U.S.A.*, 71 (1974), 4133-4134. https://doi.org/10.1073/pnas.71.10.4133
- [8] C. Fefferman, E. M. Stein,  $H^p$  spaces of several variables,  $Acta\ Math.$ , 129 (1972), 137-193. https://doi.org/10.1007/bf02392215
- [9] J. J. F. Fournier, J. Stewart, Amalgams of  $L^p$  And  $l^q$ , Bull. Amer. Math. Soc. (N.S.), **13** (1985), 1-21. https://doi.org/10.1090/s0273-0979-1985-15350-9
- [10] M. Frazier, B. Jawerth, Decomposition of Besov spaces, *Indiana Univ. Math. J.*, **34** (1985), 777-799. https://doi.org/10.1512/iumj.1985.34.34041
- [11] M. Frazier, B. Jawerth, A discrete transform and decompositions of distribution spaces, J. Funct. Anal., 93 (1990), 34-170. https://doi.org/10.1016/0022-1236(90)90137-a
- [12] D. Goldberg, A Local Version of Real Hardy Spaces, Duke Math. J., 46
   (1) (1979), 27-42. https://doi.org/10.1215/s0012-7094-79-04603-9
- [13] F. Holland, Harmonic Analysis on Amalgams of  $L^p$  And  $l^q$ , J. London Math. Soc., 10 (1975), 295-305. https://doi.org/10.1112/jlms/s2-10.3.295
- [14] H. M. Ho, Block Decomposition and Boundedness of Pseudodifferential Operators Acting on Amalgams, *Pure Mathematical Sciences*, **10** (1) (2021), 11-26. https://doi.org/10.12988/pms.2021.91251

[15] H. M. Ho, On Block Decomposition of Local Hardy-type Amalgam Spaces, *International Journal of Mathematical Analysis*, **15** (3) (2021), 121 - 138. https://doi.org/10.12988/ijma.2021.912195

- [16] H. M. Ho, Boundedness of Calderón-Zygmund Singular Integral Operators from  $l^1(L^p)$  into  $weak-l^1(L^p)$ , International Journal of Mathematical Analysis, **15** (3) (2021), 107 120. https://doi.org/10.12988/ijma.2021.912191
- [17] H. M. Ho, On the Calderón-Zygmund Singular Integral Operators and Local Hardy-type Amalgam Spaces (Part 2), *submitted*.
- [18] R. H. Latter, A decomposition of  $H^p$  in terms of atoms, *Studia Math.*, **62** (1977), 93-101. https://doi.org/10.4064/sm-62-1-93-101
- [19] Y. Rakotondratsimba, Fractional Maximal and Integral Operators on Weighted Amalagam Spaces, J. Korean Math. Soc., 36 (5) (1999), 855-890.
- [20] Walter Rudin, Functional Analysis, 2nd ed., McGraw-Hill, New York, 1991.
- [21] E. M. Stein, Singular Integrals and Differentialibility Properties of Functions, Princeton University Press, Princeton, New Jersey, 1970.
- [22] E. M. Stein, Harmonic Analysis-Real Variable Methods, Orthogonality and Oscillatory Integrals, Princeton University Press, Princeton, New Jersey, 1993.
- [23] N. Wiener, On the representation of functions by trigonometrical integrals, *Math. Z.*, **24** (1926), 575-616. https://doi.org/10.1007/bf01216799
- [24] N. Wiener, Tauberian theorems, Ann. of Math., 33 (1932), 1-100. https://doi.org/10.2307/1968102
- [25] M. W. Wong, An Introduction to Pseudo-Differential Operators, 2<sup>nd</sup> Ed., World Scientific, 1999. https://doi.org/10.1142/4047

Received: March 29, 2021; Published: April 22, 2021