

A Solution of the Two-dimensional Modified Poisson-Boltzmann Equation

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Abstract

In this work, we find solutions to the modified Poisson-Boltzmann (PB) equation using the Tanh and Riccati solitary wave techniques.

Keywords: Modified Poisson-Boltzmann equation, Tanh method

1 Introduction

When in an electrolyte solution is incorporated the finite size of particles, an entropic effect, called excluded volume, we get the modified Poisson-Boltzmann equation [1]-[4]. In the search to find new solutions to non-linear partial differential equations, in recent years, we have witnessed the creation of a huge number of methods called solitary waves. Among them, we use the tanh [5] and Riccati [6] in order to find solutions to the modified Poisson-Boltzmann equation.

2 The Modified Poisson-Boltzmann equation

We start from the two-dimensional modified Poisson-Boltzmann equation [1]-[4]:

$$\nabla^2 \Psi = \frac{\partial^2 \Psi}{dx^2} + \frac{\partial^2 \Psi}{dy^2} = \frac{8\pi en_0}{\epsilon} \frac{\sinh\left(\frac{e\Psi}{k_B T}\right)}{1 - \phi_0 + \phi_0 \cosh\left(\frac{e\Psi}{k_B T}\right)} \quad (1)$$

$$\psi = \frac{e\Psi}{k_B T} \quad (2)$$

$$\frac{\partial^2 \psi}{dx^2} + \frac{\partial^2 \psi}{dy^2} = \frac{8\pi e^2 n_0}{\epsilon k_B T} \frac{\sinh(\psi)}{1 - \phi_0 + \phi_0 \cosh(\psi)} \quad (3)$$

We define $\kappa^2 = \frac{8\pi e^2 n_0}{\epsilon k_B T}$. Then, we apply the following coordinate transformation:

$$u = x + y \quad (4)$$

The derivatives change like:

$$\frac{\partial}{\partial x} = \frac{d}{du}; \quad \frac{\partial}{\partial y} = \frac{d}{du}; \quad \frac{\partial^2}{\partial x^2} = \frac{d^2}{du^2}; \quad \frac{\partial^2}{\partial y^2} = \frac{d^2}{du^2} \quad (5)$$

Replacing eqs. (5) in eq. (3), we get:

$$2 \frac{d^2 \psi}{du^2} = \kappa^2 \frac{\sinh(\psi)}{1 - \phi_0 + \phi_0 \cosh(\psi)} \quad (6)$$

Also, we do the next transformation

$$\psi = 2 \tanh^{-1}(v) \quad (7)$$

So, the derivatives are:

$$\frac{d\psi}{du} = \frac{2}{1-v^2} \frac{dv}{du}; \quad \frac{d^2 \psi}{du^2} = \frac{4v}{(1-v^2)^2} \left(\frac{dv}{du}\right)^2 + \frac{2}{1-v^2} \frac{d^2 v}{du^2} \quad (8)$$

And replacing in eq. (6)

$$\left(\frac{4}{1-v^2} \frac{d^2 v}{du^2} + \frac{8v}{(1-v^2)^2} \left(\frac{dv}{du}\right)^2 \right) (1 - \phi_0 + \phi_0 \cosh(2 \tanh^{-1}(v))) = \kappa^2 \sinh(2 \tanh^{-1}(v)) \quad (9)$$

Taking into account that

$$\sinh(2 \tanh^{-1}(v)) = \frac{2v}{1-v^2}; \quad \cosh(2 \tanh^{-1}(v)) = \frac{1+v^2}{1-v^2} \quad (10)$$

Then, replacing eq. (10) in eq. (9)

$$4(1-v^2)(1-v^2+2\phi_0 v^2) \frac{d^2 v}{du^2} + 8v(1-v^2+2\phi_0 v^2) \left(\frac{dv}{du}\right)^2 = \kappa^2 2v(1-v^2)^2 \quad (11)$$

Now, we introduce a new independent variable [5]:

$$Y = \tanh(\mu u) \quad (12)$$

Then, the derivatives of Y , are:

$$\frac{d}{du} = \mu(1 - Y^2) \frac{d}{dY}; \quad \frac{d^2}{du^2} = -2\mu^2 Y(1 - Y^2) \frac{d}{dY} + \mu^2(1 - Y^2)^2 \frac{d^2}{dY^2} \quad (13)$$

The solutions are postulated as:

$$v = \sum_{i=1}^m a_i Y^i \quad (14)$$

Then, replacing

$$\begin{aligned} & 4\mu^2(1 - v^2)(1 - v^2 + 2\phi_0 v^2)(-2Y(1 - Y^2) \frac{dv}{dY} \\ & + (1 - Y^2)^2 \frac{d^2 v}{dY^2}) + 8\mu^2 v(1 - v^2 + 2\phi_0 v^2)((1 - Y^2) \frac{d}{dY})^2 \\ & - 2\kappa^2 v(1 - v^2)^2 = 0 \end{aligned} \quad (15)$$

Now, we balance the highest-order linear derivative with the highest order nonlinear terms in eq. (15). Then

$$v^2 Y^4 \frac{d^2 v}{du^2} \rightarrow v^5 \rightarrow 2m + 4 + m - 2 = 5m \rightarrow 3m + 2 = 5m \rightarrow m = 1 \quad (16)$$

So, eq. (14) is:

$$v = a_0 + a_1 Y \quad (17)$$

Replacing in eq. (15), we get

$$\begin{aligned} & 4\mu^2(1 - (a_0 + a_1 Y)^2)(1 - (a_0 + a_1 Y)^2 + 2\phi_0(a_0 + a_1 Y)^2) \times \\ & (-2Y(1 - Y^2)a_1) + 8\mu^2(a_0 + a_1 Y)(1 - (a_0 + a_1 Y)^2 + 2\phi_0(a_0 + a_1 Y)^2) \times \\ & ((1 - Y^2)a_1)^2 - 2\kappa^2(a_0 + a_1 Y)(1 - (a_0 + a_1 Y)^2)^2 = 0 \end{aligned} \quad (18)$$

Defining $l_1 = (9 + \sqrt{69})$, $l_2 = (9 - \sqrt{69})$ and $l_3 = (\frac{27}{2} - \frac{3\sqrt{69}}{2})$. Doing some algebra

$$\begin{aligned}
f_1 \rightarrow (a_0 = 0, a_1 = -\sqrt{1 + \frac{1}{3}l_3^{1/3} + \frac{\left(\frac{1}{2}l_1\right)^{1/3}}{3^{2/3}}, \phi_0 = \frac{1}{36}(-6} & \quad (19) \\
-3 \times 2^{2/3} (3l_2)^{1/3} + 2^{1/3} (3l_2)^{2/3} - 3 \times 2^{2/3} (3l_1)^{1/3} + 2^{1/3} (3l_1)^{2/3}), \\
\kappa^2 = \frac{4}{3}l_3^{1/3} \mu^2 + 2 \left(\frac{2}{3}\right)^{2/3} l_1^{1/3} \mu^2)
\end{aligned}$$

$$\begin{aligned}
f_2 \rightarrow (a_0 = 0, a_1 = \sqrt{1 + \frac{1}{3}l_3^{1/3} + \frac{\left(\frac{1}{2}l_1\right)^{1/3}}{3^{2/3}}, \phi_0 = \frac{1}{36}(-6} & \quad (20) \\
-3 \times 2^{2/3} (3l_2)^{1/3} + 2^{1/3} (3l_2)^{2/3} - 3 \times 2^{2/3} (3l_1)^{1/3} + 2^{1/3} (3l_1)^{2/3}), \\
\kappa^2 = \frac{4}{3}l_3^{1/3} \mu^2 + 2 \left(\frac{2}{3}\right)^{2/3} l_1^{1/3} \mu^2)
\end{aligned}$$

$$\begin{aligned}
f_3 \rightarrow (a_0 = 0, a_1 = -(1 - \frac{1}{6}l_3^{1/3} + \frac{il_3^{1/3}}{2\sqrt{3}} - \frac{\left(\frac{1}{2}l_1\right)^{1/3}}{2 \times 3^{2/3}} - \frac{il_1^{1/3}}{2 \times 2^{1/3}3^{1/6}})^{1/2}, & \quad (21) \\
\phi_0 = -\frac{1}{72}i(-12i + 3 \times 2^{2/3}3^{5/6}l_2^{1/3} + 3 \times 2^{1/3}3^{1/6}l_2^{2/3} + 3i2^{2/3} (3l_2)^{1/3} \\
-i2^{1/3} (3l_2)^{2/3} - 3 \times 2^{2/3}3^{5/6}l_1^{1/3} - 3 \times 2^{1/3}3^{1/6}l_1^{2/3} + 3i2^{2/3} (3l_1)^{1/3} \\
-i2^{1/3} (3l_1)^{2/3}), \kappa^2 = -\frac{2}{3}l_3^{1/3} \mu^2 + \frac{2il_3^{1/3} \mu^2}{\sqrt{3}} - \left(\frac{2}{3}\right)^{2/3} l_1^{1/3} \mu^2 - \frac{i2^{2/3}l_1^{1/3} \mu^2}{3^{1/6}})
\end{aligned}$$

$$\begin{aligned}
f_4 \rightarrow (a_0 = 0, a_1 = (1 - \frac{1}{6}l_3^{1/3} + \frac{il_3^{1/3}}{2\sqrt{3}} - \frac{\left(\frac{1}{2}l_1\right)^{1/3}}{2 \times 3^{2/3}} - \frac{il_1^{1/3}}{2 \times 2^{1/3}3^{1/6}})^{1/2}, & \quad (22) \\
\phi_0 = -\frac{1}{72}i(-12i + 3 \times 2^{2/3}3^{5/6}l_2^{1/3} + 3 \times 2^{1/3}3^{1/6}l_2^{2/3} + 3i2^{2/3} (3l_2)^{1/3} \\
-i2^{1/3} (3l_2)^{2/3} - 3 \times 2^{2/3}3^{5/6}l_1^{1/3} - 3 \times 2^{1/3}3^{1/6}l_1^{2/3} + 3i2^{2/3} (3l_1)^{1/3} \\
-i2^{1/3} (3l_1)^{2/3}), \kappa^2 = -\frac{2}{3}l_3^{1/3} \mu^2 + \frac{2il_3^{1/3} \mu^2}{\sqrt{3}} - \left(\frac{2}{3}\right)^{2/3} l_1^{1/3} \mu^2 - \frac{i2^{2/3}l_1^{1/3} \mu^2}{3^{1/6}})
\end{aligned}$$

$$\begin{aligned}
f_5 \rightarrow (a_0 = 0, a_1 = -(1 - \frac{1}{6}l_3^{1/3} - \frac{il_3^{1/3}}{2\sqrt{3}} - \frac{\left(\frac{1}{2}l_1\right)^{1/3}}{2 \times 3^{2/3}} + \frac{il_1^{1/3}}{2 \times 2^{1/3}3^{1/6}})^{1/2}, & \quad (23) \\
\phi_0 = \frac{1}{72}i(12i + 3 \times 2^{2/3}3^{5/6}l_2^{1/3} + 3 \times 2^{1/3}3^{1/6}l_2^{2/3} - 3i2^{2/3} (3l_2)^{1/3} \\
+i2^{1/3} (3l_2)^{2/3} - 3 \times 2^{2/3}3^{5/6}l_1^{1/3} - 3 \times 2^{1/3}3^{1/6}l_1^{2/3} - 3i2^{2/3} (3l_1)^{1/3} \\
+i2^{1/3} (3l_1)^{2/3}), \kappa^2 = -\frac{2}{3}l_3^{1/3} \mu^2 - \frac{2il_3^{1/3} \mu^2}{\sqrt{3}} - \left(\frac{2}{3}\right)^{2/3} l_1^{1/3} \mu^2 + \frac{i2^{2/3}l_1^{1/3} \mu^2}{3^{1/6}})
\end{aligned}$$

	A_1	C_1	F
1	1/2	-1/2	$\coth(\xi) \pm \cosh(\xi), \tanh(\xi) \pm \operatorname{sech}(\xi)$
2	1/2	1/2	$\sec(\xi) \pm i \tan(\xi)$
3	-1/2	-1/2	$\csc(\xi) \pm i \cot(\xi)$
4	1	-1	$\tanh(\xi), \coth(\xi)$
5	1	1	$\tan(\xi)$
6	-1	-1	$\cot(\xi)$

Table 1: Solutions for eq. (11), [6].

$$f_6 \rightarrow (a_0 = 0, a_1 = (1 - \frac{1}{6}l_3^{1/3} - \frac{il_3^{1/3}}{2\sqrt{3}} - \frac{(\frac{1}{2}l_1)^{1/3}}{2 \times 3^{2/3}} + \frac{il_1^{1/3}}{2 \times 2^{1/3}3^{1/6}})^{1/2}), \quad (24)$$

$$\phi_0 = \frac{1}{72}i(12i + 3 \times 2^{2/3}3^{5/6}l_2^{1/3} + 3 \times 2^{1/3}3^{1/6}l_2^{2/3} - 3i2^{2/3}(3l_2)^{1/3}$$

$$+ i2^{1/3}(3l_2)^{2/3} - 3 \times 2^{2/3}3^{5/6}l_1^{1/3} - 3 \times 2^{1/3}3^{1/6}l_1^{2/3} - 3i2^{2/3}(3l_1)^{1/3}$$

$$+ i2^{1/3}(3l_1)^{2/3}), \kappa^2 = -\frac{2}{3}l_3^{1/3}\mu^2 - \frac{2il_3^{1/3}\mu^2}{\sqrt{3}} - \left(\frac{2}{3}\right)^{2/3}l_1^{1/3}\mu^2 + \frac{i2^{2/3}l_1^{1/3}\mu^2}{3^{1/6}}$$

Then, we get 6 families of solutions

3 Solitary wave method 2, Riccati equation

We use the method in [6], in order to find solutions for eq. (11). So:

$$\Phi = \sum_{i=1}^n a_i F^i \quad (25)$$

where F solves, table (1), the Riccati equation, i.e.

$$F' = (C_1 F^2 + A_1), \quad F'' = 2C_1 F(C_1 F^2 + A_1) \quad (26)$$

here A_1 and C_1 are constants, table (1). Replacing in eq. (11), and balancing nonlinear terms, we get $n = 1$. Then, eq. (25) is, $v = (a_0 + a_1 F)$. So, the first derivative is:

$$v' = (a_1 F') = a_1(C_1 F^2 + A_1) \quad (27)$$

$$v'' = (a_1 F'') = 2a_1 C_1 F(C_1 F^2 + A_1)$$

Replacing in eq. (11)

$$4(1 - (a_0 + a_1 F)^2)(1 - (a_0 + a_1 F)^2 + 2\phi_0(a_0 + a_1 F)^2)2a_1 C_1 F(C_1 F^2 + A_1) + 8(a_0 + a_1 F)(1 - (a_0 + a_1 F)^2 + 2\phi_0(a_0 + a_1 F)^2)(a_1(C_1 F^2 + A_1))^2 - \kappa^2 2(a_0 + a_1 F)(1 - (a_0 + a_1 F)^2)^2 = 0 \quad (28)$$

Then, we get a set of equations, order by order in F^i . And doing some algebra, we get:

$$g_1 \rightarrow \left\{ a_0 = 0, a_1 = -\frac{i\sqrt{C_1}}{\sqrt{A_1}}, \kappa = 0 \right\}, g_2 \rightarrow \left\{ a_0 = 0, a_1 = \frac{i\sqrt{C_1}}{\sqrt{A_1}}, \kappa = 0 \right\} \quad (29)$$

Then, we get 12 families of solutions, g_i , using Ricatti method.

4 Conclusions

We solved the modified Poisson-Boltzmann equation using the Tanh and Riccati methods. We find 6 families of solutions using tanh and 12 Riccati methods. As a future work, we can extend the method to investigate spherical and cylindrical geometries. The solutions are:

$$\Psi(x, y) = 2\frac{k_B T}{e} \tanh^{-1}(a_1 \tanh(\mu(x + y))) \quad (30)$$

$$\Psi(x, y) = 2\frac{k_B T}{e} \tanh^{-1}(a_1 F(x + y)) \quad (31)$$

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