

# **Bistability between a Chaotic Bursting State and a Depolarized Steady State in a Mathematical Model of Snail RPa1 Neurons**

**Takaaki Shirahata**

Kagawa School of Pharmaceutical Sciences  
Tokushima Bunri University  
8-53 Hamano-cho, Takamatsu-shi  
Kagawa 760-8542, Japan

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## **Abstract**

This study presents a novel form of bistability in a mathematical model of snail RPa1 neurons. This model was previously reported as a system of nonlinear ordinary differential equations simulating the membrane potential oscillations of snail RPa1 neurons. Here, we perform a numerical simulation revealing that whether the model shows a chaotic bursting state or a depolarized steady state is dependent on the initial condition. Notably, a certain transient current pulse can change the dynamical state of the model from a chaotic bursting state to a depolarized steady state, or vice versa. Taken together, these results indicate that the snail RPa1 neuron model can show bistability between chaotic bursting and depolarized steady states.

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**Keywords:** bistability, mathematical model of snail RPa1 neurons, chaotic bursting state, depolarized steady state

## 1 Introduction

The mathematical model of snail RPa1 neurons is noteworthy because it shows rich dynamical behaviors. Specifically, this model shows two types of chaotic activity: a chaotic spiking state and a chaotic bursting state [1]. This model was originally developed as a system of nonlinear ordinary differential equations (ODEs) based on the Hodgkin–Huxley concept to simulate the dynamics of the membrane potential of snail RPa1 neurons [1]. In addition to showing chaotic activity, a previous study demonstrated that the model shows a certain type of bistability, namely, the coexistence of a chaotic spiking state and a periodic bursting state under a certain specific parameter condition [1]. Therefore, it is natural to ask whether the model also exhibits bistability that involves a chaotic bursting state and, if so, the type of bistability that the model shows. In other words, what kind of dynamical state can coexist with a chaotic bursting state? Answering this question will provide a deeper understanding of the relationship between bistability and chaotic activity in the snail RPa1 neuron model. However, this question remains unanswered. Previous studies report that a chaotic bursting state can change into different dynamical states due to variation in the various model parameters [2, 3]. For example, a chaotic bursting state can change into a depolarized steady state due to variation in the time constant of potassium conductance [2]. Given this previous finding, we hypothesized that the model may exhibit bistability between a chaotic bursting state and a depolarized steady state. To evaluate this hypothesis, the present study performs a numerical simulation of the snail RPa1 neuron model.

## 2 Mathematical Model

The snail RPa1 neuron model investigated in the present study is described by the following system of nonlinear ODEs (see [1] for further details):

$$\begin{aligned} \frac{dV}{dt} = & \frac{1}{0.02} \left( I_{app} - 0.13 \left( \frac{1}{1 + e^{-0.2(V+45)}} \right) (V - 40) - 0.18m_B h_B (V + 58) \right. \\ & - 0.02(V - 40) - 0.25(V + 70) \\ & - 400m^3 h(V - 40) - 10n^4(V + 70) \\ & \left. - m_{Ca}^2(V - 150) - 0.01 \left( \frac{1}{1 + e^{-0.06(V+45)}} \right) \left( \frac{1}{1 + e^{15000([Ca] - 0.00004)}} \right) (V - 150) \right) \end{aligned} \quad (1)$$

$$\frac{dm_B}{dt} = \frac{1}{0.05} \left( \frac{1}{1 + e^{0.4(V+34)}} - m_B \right) \quad (2)$$

$$\frac{dh_B}{dt} = \frac{1}{1.5} \left( \frac{1}{1 + e^{-0.55(V+43)}} - h_B \right) \quad (3)$$

$$\frac{dm}{dt} = \frac{1}{0.0005} \left( \frac{1}{1 + e^{-0.4(V+31)}} - m \right) \quad (4)$$

$$\frac{dh}{dt} = \frac{1}{0.01} \left( \frac{1}{1 + e^{0.25(V+45)}} - h \right) \quad (5)$$

$$\frac{dn}{dt} = \frac{1}{0.015} \left( \frac{1}{1 + e^{-0.18(V+25)}} - n \right) \quad (6)$$

$$\frac{dm_{Ca}}{dt} = \frac{1}{0.01} \left( \frac{1}{1 + e^{-0.2V}} - m_{Ca} \right) \quad (7)$$

$$\frac{d[Ca]}{dt} = 0.002 \left( -\frac{m_{Ca}^2 (V - 150)}{2F \left( \frac{4}{3} \pi 0.1^3 \right)} - 50[Ca] \right) \quad (8)$$

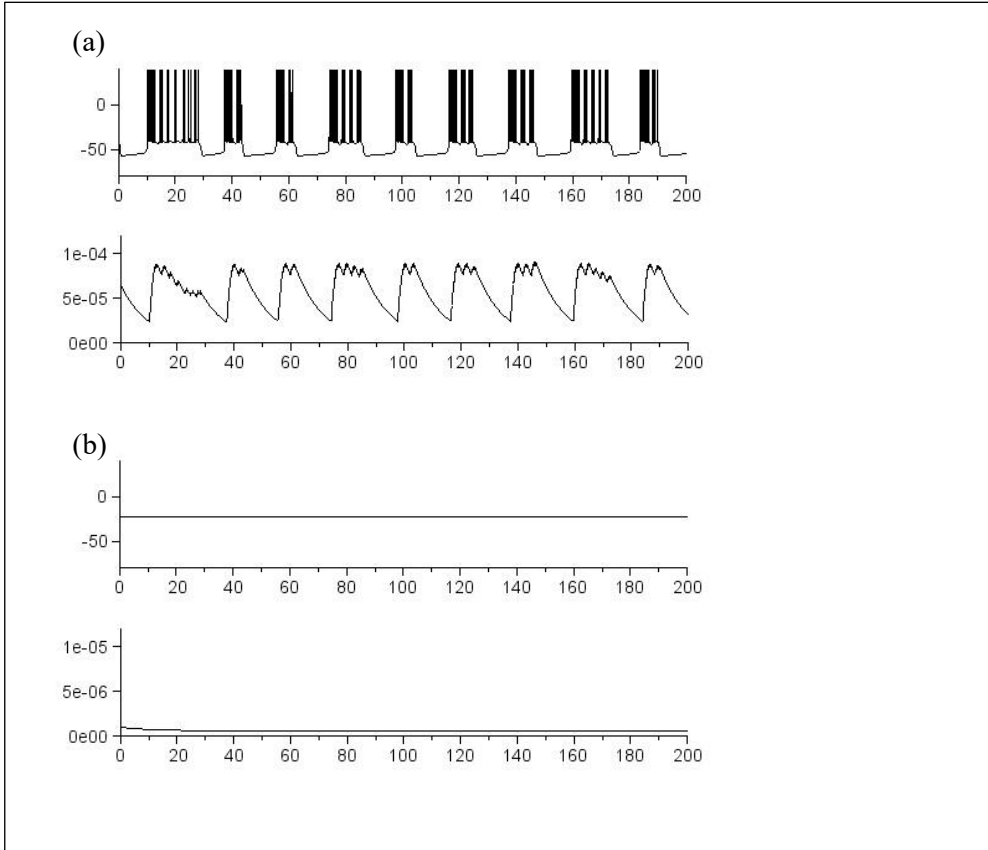
, in which  $V$  (mV) (the membrane potential of snail RPa1 neurons),  $m_B$ ,  $h_B$ ,  $m$ ,  $h$ ,  $n$ , and  $m_{Ca}$  (various gating variables), and  $[Ca]$  (mM) (the concentration of intracellular calcium) are state variables;  $t$  (s) is time;  $I_{app}$  (nA) (a transient current pulse) is a system parameter; and  $F$  is the Faraday constant. Equations (1)–(8) are numerically solved using the free and open source software Scilab (<http://www.scilab.org/>).

### 3 Numerical Results

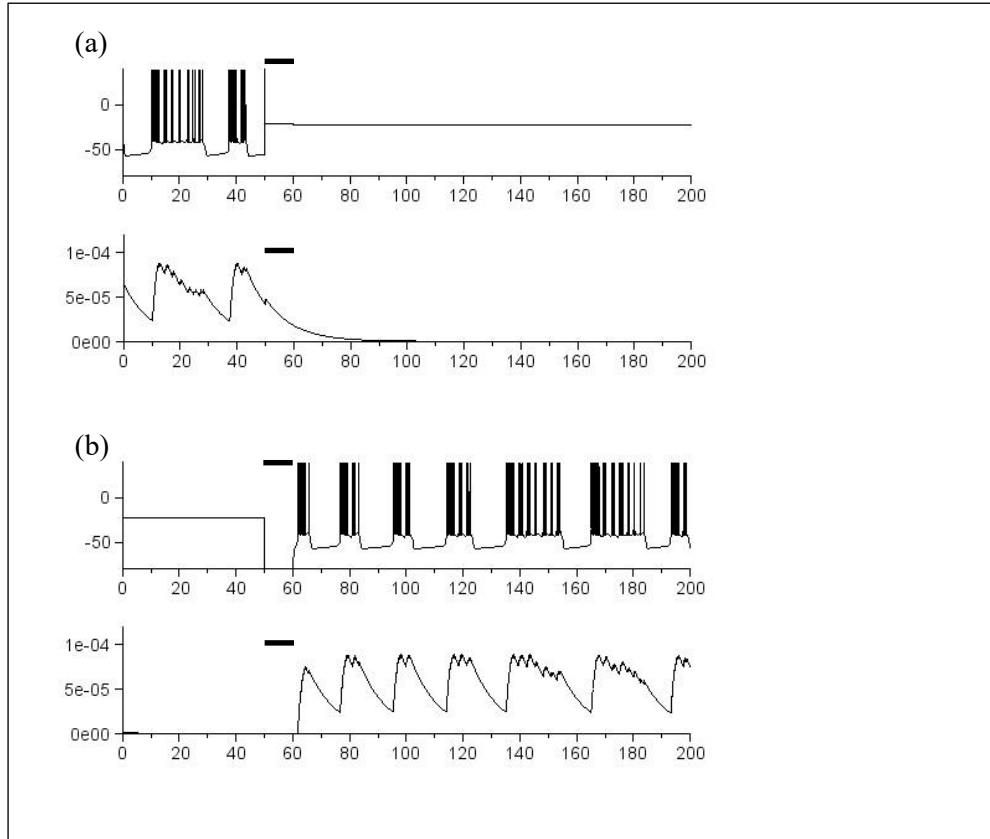
First, we perform a numerical simulation of the snail RPa1 neuron model varying the initial conditions in which  $I_{app}$  is set to 0. As  $V$  and  $[Ca]$  are experimentally interesting state variables [1], we illustrate the time courses of  $V$  and  $[Ca]$  in Figure 1. When the initial condition is set that  $V = -42$  mV,  $m_B = 0.95$ ,  $h_B = 0.77$ ,  $m = 0.14$ ,  $h = 0.1$ ,  $n = 0.048$ ,  $m_{Ca} = 0.0002$ , and  $[Ca] = 6.5 \times 10^{-5}$  mM, the model shows a chaotic bursting state: bursting with various burst duration lengths is observed (Figure 1a, top panel). In contrast, when the initial condition is set that  $V = -22$  mV,  $m_B = 0.01$ ,  $h_B = 0.999$ ,  $m = 0.98$ ,  $h = 0.003$ ,  $n = 0.62$ ,  $m_{Ca} = 0.01$ , and  $[Ca] = 1.0 \times 10^{-6}$  mM, the model shows a depolarized steady state: no oscillatory activity is observed (Figure 1b top).

Next, we perform a numerical simulation varying  $I_{app}$ , and the time courses of  $V$  and  $[Ca]$  (Figure 2). When the simulation starts with same initial conditions as in Figure 1a, the model shows a chaotic bursting state when  $t$  is at 50 s or earlier (Figure 2a, top panel). Under such conditions, if a transient current pulse is injected into the model (i.e.,  $t$  is between 50 and 60 s,  $I_{app}$  is 10 nA; otherwise,  $I_{app}$  is 0 nA), the dynamical state of the model changes to a depolarized steady state during the transient current pulse (see the horizontal bar in Figure 2a, top panel).

Even if the current pulse is terminated (i.e.,  $t$  is 60 s and longer), the dynamical state of the model does not recover to a chaotic bursting state, i.e., the depolarized steady state continues (Figure 2a, top). Finally, we investigate the opposite direction (i.e., a depolarized steady state  $\rightarrow$  a chaotic bursting state) where the simulation starts with the initial condition shown in Figure 1(b) and a transient current pulse is injected (i.e.,  $t$  is 50 to 60 s,  $I_{app}$  is  $-40$  nA; otherwise,  $I_{app}$  is 0 nA) (Figure 2b). In this case, we observe that the dynamical state of the model changes from a depolarized steady state to a chaotic bursting state (Figure 2b, top).



**Figure 1.** Time courses of  $V$  and  $[Ca]$  of the snail RPa1 neuron model under different initial conditions. The initial conditions are as follows: in (a):  $V = -42$  mV,  $m_B = 0.95$ ,  $h_B = 0.77$ ,  $m = 0.14$ ,  $h = 0.1$ ,  $n = 0.048$ ,  $m_{Ca} = 0.0002$ , and  $[Ca] = 6.5 \times 10^{-5}$  mM; and in (b)  $V = -22$  mV,  $m_B = 0.01$ ,  $h_B = 0.999$ ,  $m = 0.98$ ,  $h = 0.003$ ,  $n = 0.62$ ,  $m_{Ca} = 0.01$ , and  $[Ca] = 1.0 \times 10^{-6}$  mM. In both cases,  $I_{app}$  is 0. In (a) and (b), the horizontal axis indicates  $t$  (s) in both the top and bottom panels, while the vertical axis indicates  $V$  (mV) in the top panel and  $[Ca]$  (mM) in the bottom panel.



**Figure 2.** The transition from a chaotic bursting state to a depolarized steady state in (a) and vice versa in (b) due to the transient current pulse. (a) A current pulse with an amplitude of 10 nA and a duration of 10 s is injected at 50 s. The initial conditions are the same as Figure 1(a). (b) A current pulse with an amplitude of -40 nA and a duration of 10 s is injected at 50 s. The initial conditions are the same as Figure 1(b). In (a) and (b), the horizontal axis indicates  $t$  (s) in both the top and bottom panels, while the vertical axis indicates  $V$  (mV) in the top panel and  $[Ca]$  (mM) in the bottom panel. Horizontal bars between 50 and 60 s denote the duration of the transient current pulse.

## 4 Conclusion

This study focuses on the bistability of a mathematical model of snail RPa1 neurons, verifying a novel type of bistability between a chaotic bursting state and a depolarized steady state. This type of bistability differs from the previously reported bistability between a chaotic spiking state and a periodic bursting state [1]. Notably, bistability has been reported in previous studies of mathematical

models of neurons different from the snail RPa1 neuron model (e.g., a leech interneuron model [4], a circadian pacemaker neuron model [5]). Again, the type of bistability reported in previous studies differs from that observed in the present study: the leech interneuron model shows bistability between a periodic bursting state and a hyperpolarized steady state [4], while the circadian pacemaker neuron model shows bistability of a periodic spiking state and a depolarized steady state [5]. In conclusion, the present study contributes to a deeper understanding of the relationship between bistability and chaotic activity in the snail RPa1 neuron model.

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