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The Cauchy Problem for a Symmetric System of Keyfitz-Kranzer Type with Linear Damping

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Abstract

We prove the existence of a weak solution for the Cauchy problem associated with a 2×2 symmetric system of Keyfitz-Kranzer type with linear damping.

Keywords: Symmetric system of Keyfitz-Kranzer type, linear damping, existence, weak solution, compensated compactness method

1 Introduction

The following system of partial differential equations

$$\begin{cases} u_t + (u\phi(r))_x = 0\\ v_t + (v\phi(r))_x = 0, \end{cases}$$
(1.1)

where $\phi(r)$ is a nonlinear symmetric function of u and v, is a 2×2 symmetric system of Keyfitz-Kranzer type.

A system of the form (1.1) was first introduced in [6] by B. Keyfitz and H. Kranzer as a model in elasticity theory. Also, this type of system appear in magnetohydrodynamics, chromatography and enhanced oil recovery [1, 4, 7]. Symmetric systems of Keyfitz-Kranzer type have been studied by many authors [1, 2, 4, 5, 6, 7, 8, 9].

In this paper we shall apply the vanishing viscosity method together with the Murat's lemma and the div-curl lemma, to study the Cauchy problem for the 2×2 symmetric system of Keyfitz-Kranzer type with linear damping

$$\begin{cases} u_t + (u\phi(r))_x = -au \\ v_t + (v\phi(r))_x = -bv, \end{cases}$$
 (1.2)

with bounded measurable initial data

$$(u(x,0), v(x,0)) = (u_0(x), v_0(x)), \tag{1.3}$$

where $\phi(r) \in C^2(\mathbb{R}^+)$ and $\phi(r)$ is strictly increasing or decreasing for positive r, a, b are constants such that $b \geq a > 0$ and

$$r = |u|^{\alpha} + |v|^{\alpha},\tag{1.4}$$

for any $\alpha > 1$ fixed.

From [3], we see that the system (1.2) has eigenvalues

$$\lambda_1 = \phi(r) + \alpha r \phi'(r), \quad \lambda_2 = \phi(r),$$

and Riemann invariants

$$z(u,v) = \frac{v}{u}, \quad w(u,v) = \phi(r). \tag{1.5}$$

2 Existence of Weak Solution

We consider the Cauchy problem for the system

$$\begin{cases} u_t^{\epsilon} + \left(u^{\epsilon}\phi(r^{\epsilon})\right)_x = -au^{\epsilon} + \epsilon u_{xx}^{\epsilon} \\ v_t^{\epsilon} + \left(v^{\epsilon}\phi(r^{\epsilon})\right)_x = -bv^{\epsilon} + \epsilon v_{xx}^{\epsilon} \end{cases}, \tag{2.1}$$

with initial data (1.3).

Lemma 2.1. For any $\epsilon > 0$ and any T > 0, we have the a-priori bounds for the Cauchy problem (2.1)-(1.3)

$$|u^{\epsilon}(x,t)| \le M(T), \quad |v^{\epsilon}(x,t)| \le M(T), \quad (x,t) \in \mathbb{R} \times [0,T],$$
 (2.2)

for a positive constant M(T) independent of ϵ .

Proof. We multiply the first and second equations of system (2.1) respectively by $\alpha |u|^{\alpha-2}u$ and $\alpha |v|^{\alpha-2}v$, adding the results, we obtain

$$r_t + \lambda_1 r_x + a\alpha |u|^{\alpha} + b\alpha |v|^{\alpha} = \epsilon r_{xx} - \epsilon \alpha (\alpha - 1) \Big(|u|^{\alpha - 2} u_x^2 + |v|^{\alpha - 2} v_x^2 \Big). \quad (2.3)$$

We have from (2.3) the inequality

$$r_t + \lambda_1 r_x + a\alpha r \le \epsilon r_{xx}. \tag{2.4}$$

Applying the maximum principle to (2.4) we get the estimate $r^{\epsilon} \leq N(T)$, where N(T) is a positive constant, being independent of ϵ . Estimate from which we obtain the a-priori bounds in (2.2).

A consequence of the previous lemma is the following.

Corollary 2.2. For $\epsilon > 0$ and T > 0 the viscosity solution $(u^{\epsilon}(x,t), v^{\epsilon}(x,t))$ for the Cauchy problem (2.1)-(1.3) exists on $\mathbb{R} \times [0,T]$.

Lemma 2.3. If $u_0(x) \ge c_1$ for a positive constant c_1 , then

$$u^{\epsilon}(x,t) \ge c(t,\epsilon,c_1) > 0, \tag{2.5}$$

where $c(t, \epsilon, c_1)$ could tend to 0 as $t \to +\infty$ or $\epsilon \to 0$.

Proof. We set $\nu = -\ln u$ and deduce from the first equation of the system (2.1) that

$$\nu_t - \epsilon \nu_{xx} \le \frac{\left(\phi(r)\right)^2}{\epsilon} + \phi(r)_x + a.$$

Then using the previous inequality, we obtain

$$\nu(x,t) \le \nu_0(x) * \frac{1}{\sqrt{4\epsilon\pi t}} e^{-\frac{x^2}{4\epsilon t}} + \int_0^t \left(\frac{1}{\epsilon} (\phi(r))^2 + \phi(r)_x + a\right) *_x \frac{1}{\sqrt{4\epsilon\pi (t-s)}} e^{-\frac{x^2}{4\epsilon (t-s)}} ds ,$$

where $\nu_0(x) = -\ln u_0^{\epsilon}(x)$. Hence

$$\nu(x,t) \le -\ln c_1 + \frac{N_1}{\epsilon}t + N_2\sqrt{\frac{t}{\epsilon}}.$$

It follows that

$$u(x,t) \ge c_1 \exp{-\left(\frac{N_1}{\epsilon}t + N_2\sqrt{\frac{t}{\epsilon}}\right)} \ge c(t,\epsilon,c_1) > 0.$$

Lemma 2.4. Let z be the Riemann invariant given in (1.5). If in addition to the assumption of lemma 2.3, $z_0(x) = z(x,0) \in L^{\infty}(\mathbb{R})$ and $z'_0(x) \in L^1(\mathbb{R})$, then $\left(\frac{v^{\epsilon}}{u^{\epsilon}}\right)(x,t) \in L^{\infty}(\mathbb{R} \times [0,T])$, $\left(\frac{v^{\epsilon}}{u^{\epsilon}}\right)_x(\cdot,t) \in L^1(\mathbb{R})$. Moreover

$$TV\left(\left(\frac{v^{\epsilon}}{u^{\epsilon}}\right)(\cdot,t)\right) = \int_{-\infty}^{+\infty} \left| \left(\frac{v^{\epsilon}}{u^{\epsilon}}\right)_{x}(x,t) \right| dx \le \int_{-\infty}^{+\infty} \left| \left(\frac{v_{0}}{u_{0}}\right)'(x) \right| dx = TV\left(\left(\frac{v_{0}}{u_{0}}\right)(x)\right), \tag{2.6}$$

where TV is the total variation.

Proof. Multiplying the first equation of (2.1) by $-\frac{v}{u^2}$ and the second equation by $\frac{1}{u}$ and summing them up, one obtains

$$\left(\frac{v}{u}\right)_t + \phi(r)\left(\frac{v}{u}\right)_x + (b-a)\frac{v}{u} = \epsilon\left(\frac{v}{u}\right)_{xx} + 2\epsilon\frac{u_x}{u}\left(\frac{v}{u}\right)_x,\tag{2.7}$$

Appliying the maximum principle to (2.7), we thus find that $\left(\frac{v^{\epsilon}}{u^{\epsilon}}\right)(x,t) \in L^{\infty}(\mathbb{R} \times [0,T])$. Now we differentiate (2.7) with respect to x and then we do $\theta = \left(\frac{v}{u}\right)_x$ to get

$$\theta_t + (\phi(r)\theta)_x + (b-a)\theta = \epsilon\theta_{xx} + (2\epsilon u^{-1}u_x\theta)_x,$$

multiplying this equation by the sequence of smooth functions $g'(\theta, \alpha)$, where α is a parameter, we obtain

$$g(\theta,\alpha)_t + (\phi(r)g(\theta,\alpha))_x + \phi(r)_x (g'(\theta,\alpha)\theta - g(\theta,\alpha)) + (b-a)g'(\theta,\alpha)\theta_x = \epsilon g(\theta,\alpha)_{xx} - \epsilon g''(\theta,\alpha)\theta_x^2 + (2\epsilon u^{-1}u_x g(\theta,\alpha))_x + (2\epsilon u^{-1}u_x)_x (g'(\theta,\alpha)\theta - g(\theta,\alpha)).$$
(2.8)

We choose $g(\theta, \alpha)$ such that $g''(\theta, \alpha) \geq 0$, $g'(\theta, \alpha) \rightarrow \text{sign}\theta$ and $g(\theta, \alpha) \rightarrow |\theta|$ as $\alpha \rightarrow 0$, we have from (2.8)

$$|\theta|_t + (\phi(r)|\theta|)_x \le \epsilon |\theta|_{xx} + (2\epsilon u^{-1}u_x|\theta|)_x. \tag{2.9}$$

Integrating (2.9) in $\mathbb{R} \times [0, t]$, we obtain (2.6).

We establish the results related to compactness in H_{loc}^{-1} that allow us to apply the div-curl lemma.

Lemma 2.5. We assume the same conditions given in the Lemma 2.1. Then

$$r_t^{\epsilon} + \left(\int^{r^{\epsilon}} (\phi(s) + \alpha s \phi'(s)) \, ds \right)_x \tag{2.10}$$

is compact in $H^{-1}_{loc}(\mathbb{R} \times \mathbb{R}^+)$.

Proof. We rewrite (2.3) as

$$r_t + \left(\int_{-\infty}^{\infty} (\phi(s) + \alpha s \phi'(s)) ds \right)_x = \epsilon r_{xx} - \epsilon \alpha (\alpha - 1) \left(|u|^{\alpha - 2} u_x^2 + |v|^{\alpha - 2} v_x^2 \right) - \alpha \left(a|u|^{\alpha} + b|v|^{\alpha} \right). \quad (2.11)$$

Noting that the last term in the right-hand side of (2.11) is bounded in $L_{loc}^1(\mathbb{R} \times \mathbb{R}^+)$ and using the same type of argument given in Lemma 5 of [3], we obtain the conclusion of the lemma.

Lemma 2.6. Under the assumptions of Lemma 2.1, it follows that

$$\left(\int^{r^{\epsilon}} (\phi(s) + \alpha s \phi'(s)) ds\right)_{t} + \left(\int^{r^{\epsilon}} (\phi(s) + \alpha s \phi'(s))^{2} ds\right)_{x}$$
(2.12)

is compact in $H^{-1}_{loc}(\mathbb{R} \times \mathbb{R}^+)$.

Proof. Multiplying the equation (2.3) by $\phi(r) + \alpha r \phi'(r)$, we obtain

$$\left(\int^{r^{\epsilon}} (\phi(s) + \alpha s \phi'(s)) ds\right)_{t} + \left(\int^{r^{\epsilon}} (\phi(s) + \alpha s \phi'(s))^{2} ds\right)_{x} = \epsilon r_{xx} (\phi(r) + \alpha r \phi'(r)) - \epsilon \alpha (\alpha - 1) \left(|u|^{\alpha - 2} u_{x}^{2} + |v|^{\alpha - 2} v_{x}^{2}\right) (\phi(r) + \alpha r \phi'(r)) - \alpha \left(a|u|^{\alpha} + b|v|^{\alpha}\right) (\phi(r) + \alpha r \phi'(r)), \quad (2.13)$$

as the last term is bounded in $L^1_{loc}(\mathbb{R} \times \mathbb{R}^+)$, we conclude that (2.12) is compact in $H^{-1}_{loc}(\mathbb{R} \times \mathbb{R}^+)$ by following the proof given in Lemma 6 of [3].

Lemma 2.7. Suppose the conditions of Lemma 2.4 holds. Then

$$\left(\left(u^{\epsilon}\right)^{\alpha}\right)_{t} + \left(\frac{\left(u^{\epsilon}\right)^{\alpha}}{r^{\epsilon}} \int^{r^{\epsilon}} \left(\phi(s) + \alpha s \phi'(s)\right) ds\right)_{t} \tag{2.14}$$

is compact in $H^{-1}_{loc}(\mathbb{R} \times \mathbb{R}^+)$.

Proof. The first equation of the system (2.1) can be written as

$$u_t + u_x \left(\phi(|u|^{\alpha} \varphi) + \alpha |u|^{\alpha} \varphi \phi'(|u|^{\alpha} \varphi) \right) = \epsilon u_{xx} - |u|^{\alpha} u \varphi_x \phi'(|u|^{\alpha} \varphi) - au, \quad (2.15)$$

where the auxiliary function $\varphi(x,t)$ is defined by

$$\varphi = 1 + \left| \frac{v}{u} \right|^{\alpha},\tag{2.16}$$

by Lemma 2.4 $\varphi(\cdot,t)_x$ is bounded in $L^1(\mathbb{R})$. Multiplying the above equation by $\alpha |u|^{\alpha-2}u$, we obtain

$$(u^{\alpha})_t + \left(\frac{u^{\alpha}}{r} \int_{-r}^{r} (\phi(s) + \alpha s \phi'(s)) ds \right)_x = \epsilon (u^{\alpha})_{xx} - \epsilon \alpha (\alpha - 1) u^{\alpha - 2} u_x^2$$

$$- \alpha (u^{\alpha})^2 \varphi_x \phi'(u^{\alpha} \varphi) - a\alpha |u|^{\alpha}, \quad (2.17)$$

the term $a\alpha|u|^{\alpha}$ is bounded in $L^1_{loc}(\mathbb{R}\times\mathbb{R}^+)$. We skip the rest of the proof since the result is derived by following exactly the same proof of Lemma 7 in [3].

Lemma 2.8. Let the assumptions in Lemma 2.4 hold. Then

$$u_t^{\epsilon} + \left(u^{\epsilon}\phi(r^{\epsilon})\right)_x \tag{2.18}$$

is compact in $H^{-1}_{loc}(\mathbb{R} \times \mathbb{R}^+)$.

We skip the proof of this lemma since it is similar to the one exposed to establish Lemma 8 in [3].

Corollary 2.9. Suppose the conditions of Lemma 2.4. Then we have

$$u_t^{\epsilon} + \left(u^{\epsilon}\phi(r^{\epsilon}) + \frac{v^{\epsilon}}{u^{\epsilon}}\right)_r \tag{2.19}$$

is compact in $H_{loc}^{-1}(\mathbb{R} \times \mathbb{R}^+)$.

The proof of the following result is an easy adaptation of the proof given in [3], Lemma 10.

Lemma 2.10. Assuming the hypotheses as in lemma 2.4, then

$$v_t^{\epsilon} + \left(v^{\epsilon}\phi(r^{\epsilon})\right)_r \tag{2.20}$$

is compact in $H^{-1}_{loc}(\mathbb{R} \times \mathbb{R}^+)$.

Corollary 2.11. Suppose the conditions of Lemma 2.4. Then we have

$$v_t^{\epsilon} + \left(v^{\epsilon}\phi(r^{\epsilon}) + \left(\frac{v^{\epsilon}}{u^{\epsilon}}\right)^2\right)_x \tag{2.21}$$

is compact in $H^{-1}_{loc}(\mathbb{R} \times \mathbb{R}^+)$.

Lemma 2.12. When the assumptions of Lemma 2.1 are satisfied and

$$meas\{r: (2n+1)\phi'(r) + 2nr\phi''(r) = 0\} = 0,$$
(2.22)

then there exists a subsequence of $\{r^{\epsilon}(x,t)\}$ which converges pointwisely.

Proof. We use the div-curl lemma, which can be applied to the functions (2.10) and (2.12) (for more details see [3], Lemma 12).

Lemma 2.13. Assume the hypotheses of the lemmas 2.4 and 2.12, then there is a subsequence of $\{u^{\epsilon}(x,t)\}$ which converges pointwisely.

Proof. The proof is the same to that of Lemma 13, [3].

Lemma 2.14. Under the assumptions of lemma 2.13, there is a subsequence of $\{v^{\epsilon}\}$ such that it converges pointwise.

Proof. The proof is the same as for Lemma 14 in [3].

Theorem 2.15. Suppose that $\left(\frac{v_0}{u_0}\right)(x) \in L^{\infty}(\mathbb{R}), \left(\frac{v_0}{u_0}\right)'(x) \in L^1(\mathbb{R}), \phi(r)$ is strictly increasing or decreasing for positive $r, \phi(r) \in C^2(\mathbb{R}^+)$ and meas $\{r : (2n+1)\phi'(r)+2nr\phi''(r)\}=0$. Then there exist a subsequence of $(u^{\epsilon},v^{\epsilon})$ which converges pointwisely and the limit is a weak solution of the Cauchy problem (1.2)-(1.3).

Proof. We consider the sequence of viscosity solutions $(u^{\epsilon}, v^{\epsilon})$ of the system (2.1). Let us consider $\varphi, \psi \in C_0^{\infty}(\mathbb{R} \times [0, \infty))$. By multiplying the first equation of the system (2.1) by φ , the second by ψ , adding the resulting equations and integrating by parts in $\mathbb{R} \times [0, \infty)$, we obtain that u^{ϵ} and v^{ϵ} satisfy the weak formulation of the Cauchy problem (2.1)-(1.3),

$$\int_{\mathbb{R}} \int_{0}^{+\infty} \left(u^{\epsilon} \varphi_{t} + u^{\epsilon} \phi(r^{\epsilon}) \varphi_{x} - a u^{\epsilon} \varphi \right) dt dx + \int_{\mathbb{R}} u_{0} \varphi(x, 0) dx
+ \int_{\mathbb{R}} \int_{0}^{+\infty} \left(v^{\epsilon} \psi_{t} + v^{\epsilon} \phi(r^{\epsilon}) \psi_{x} - b v^{\epsilon} \psi \right) dt dx + \int_{\mathbb{R}} v_{0} \psi(x, 0) dx =
- \epsilon \int_{\mathbb{R}} \int_{0}^{+\infty} \left(u^{\epsilon} \varphi_{xx} + v^{\epsilon} \psi_{xx} \right) dt dx .$$
(2.23)

By Lemmas 2.12, 2.13 and 2.14, we can find a subsequence of $(u^{\epsilon}, v^{\epsilon})$ (no relabeled), which converges pointwise, a. e. $(x,t) \in \mathbb{R} \times [0,T]$, to (u,v) and it is such that $r^{\epsilon} \to |u|^{\alpha} + |v|^{\alpha}$, a. e. $(x,t) \in \mathbb{R} \times [0,T]$. Since ϕ is continuous, $\phi(r^{\epsilon}) \to \phi(|u|^{\alpha} + |v|^{\alpha})$, a. e. $(x,t) \in \mathbb{R} \times [0,T]$. Note that by (2.2),

$$\left| \epsilon \int_{\mathbb{R}} \int_{0}^{+\infty} u^{\epsilon} \varphi_{xx} \, dt \, dx \right| \leq \epsilon ||\varphi_{xx}||_{L^{\infty}} \int_{\operatorname{supp}(\varphi)} |u^{\epsilon}| \, dt \, dx$$
$$\leq \epsilon N(T) \Big(meas \big(\operatorname{supp}(\varphi) \big) \Big),$$

thus we obtain

$$\lim_{\epsilon \to 0} \epsilon \int_{\mathbb{R}} \int_{0}^{+\infty} u^{\epsilon} \varphi_{xx} \, dt \, dx = 0. \tag{2.24}$$

Using the above argument we also have

$$\lim_{\epsilon \to 0} \epsilon \int_{\mathbb{R}} \int_{0}^{+\infty} v^{\epsilon} \psi_{xx} \, dt \, dx = 0. \tag{2.25}$$

We want to pass to the limit the weak formulation (2.23) to complete the proof. From (2.24) and (2.25), it follows immediately that the integral on the right-hand side of (2.23) converges to 0 as $\epsilon \to 0$. Due to the convergence almost

everywhere, we can apply the Lebesgue dominated convergence theorem to (2.23) to obtain that (u, v) is a weak solution of the Cauchy problem (1.2)-(1.3).

References

- [1] G.-Q. Chen, Hyperbolic systems of conservation laws with a symmetry, *Commun. in Partial Differential Equations*, **16** (1991), 1461 1487. https://doi.org/10.1080/03605309108820806
- [2] H. Freisthuhler, On the Cauchy problem for a class of hyperbolic systems of conservation laws, *J. Diff. Eqs.*, **112** (1994), 170 178. https://doi.org/10.1006/jdeq.1994.1099
- [3] J. C. Hernández, Existence of weak entropy solution for a symmetric system of Keyfitz-Kranzer type, *Revista Colombiana de Matemáticas*, **47** (2013), 13 28.
- [4] F. James, Y.-J. Peng and B. Perthame, Kinetic formulation for chromatography and some other hyperbolic systems, *J. Math. Pure Appl.*, **74** (1995), 367 385.
- [5] A. J. Kearsley, A. M. Reiff, Existence of weak solutions to a class of nonstrictly hyperbolic conservation laws with non-interacting waves, *Pacific J. of Math.*, **205** (2002), 153 - 170. https://doi.org/10.2140/pjm.2002.205.153
- [6] B. L. Keyfitz, H. C. Kranzer, A system of non-strictly hyperbolic conservation laws arising in elasticity theory, Arch. Rat. Mech. Anal., 72 (1980), 219 241. https://doi.org/10.1007/bf00281590
- [7] T.-P. Liu, J.-H.Wang, On a nonstrictly hyperbolic system of conservation laws, *Journal of Differential Equations*, **57** (1985), 1 14. https://doi.org/10.1016/0022-0396(85)90068-3
- [8] E. Yu. Panov, On the theory of generalized entropy solutions of the Cauchy problem for a class of non-strictly hyperbolic systems of conservation laws, *Sbornik: Mathematics*, **191** (2000), 121 150. https://doi.org/10.1070/sm2000v191n01abeh000450
- [9] E. Yu. Panov, On infinite-dimensional Keyfitz-Kranzer systems of conservation laws, *Differential Equations*, **45** (2009), 274 278. https://doi.org/10.1134/S0012266109020165

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