

Completeness of Lorentz Metric on 3-Dimensional Heisenberg Group

Essin Turhan

Firat University, Department of Mathematics
23119, Elazığ, Turkey
eturhan@firat.edu.tr

Abstract

We consider three-dimensional Heisenberg group which is given with possible three situations for Lorentz metrics and obtain completeness of Lorentz metrics on 3-dimensional Heisenberg group.

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Introduction:

It is shown that we can define three metrics on three dimensional Heisenberg group in [7]. Then in [6], Rahmani showed that one of the three metrics is flat and he obtain some geodesics using the Euler-Lagrange equations. We know from [5] that a Lie group with left invariant metrics flat if and only if the associated Lie algebra splits as an orthogonal direct sum $\mathfrak{b} \oplus \mathfrak{u}$ where \mathfrak{b} is a commutative subalgebra, \mathfrak{u} is a commutative ideal.

It is known that every left invariant pseudo-Riemannian metric on 2-step Lie group is complete. And then in [2] it is studied that completeness for the left invariant Lorentz metrics on some Lie groups which are a special class of solvable Lie groups. Guediri showed that for some groups every left invariant Lorentz metrics are non-complete.

In this paper, we study geodesically completeness of Lorentz metric on 3-dimensional Heisenberg group.

Preliminaries:

Definition: A curve $\gamma : I \longrightarrow H_3$ is inextendible if there exist no curve $\gamma' : I' \longrightarrow H_3$ such that $I' \subset I$ and $\gamma(p) = \gamma'(p)$ for all $p \in I$.

Definition: The curve $\gamma : I \longrightarrow H_3$ is an incomplete geodesic if γ is an inextendible geodesic and $I \neq \mathbb{R}$. We say that 3-dimensional Heisenberg group is geodesically incomplete if it contains an incomplete geodesic. And H_3 is geodesically complete if it isn't geodesically incomplete.

Lemma: Let H_3 be a 3-dimensional Heisenberg group with a left invariant Lorentz metric induced by a nondegenerate inner product \langle, \rangle on the \mathfrak{h}_3 . For any C^1 -curve in H_3 , $\gamma : I \longrightarrow H_3$ associate the curve $L_{\gamma(t)}^{-1} \gamma'(t)$ in \mathfrak{h}_3 .

$$t \longrightarrow \gamma(t)$$

Then the curve of \mathfrak{h}_3 associated to geodesics of H_3 are solutions of the equation

$$X' = ad_X^* X. \tag{1}$$

The Heisenberg group H_{2k+1} admits a flat left invariant Lorentz metric if and only if $k = 1$ It is proved that

$$g_3 = dx^2 + (xdy + dz)^2 - ((1 - x)dy - dz)^2$$

is flat by N. Rahmani.

Theorem: The Heisenberg group H_3 with left invariant metric is flat if and only if the associated Lie algebra \mathfrak{h}_3 splits an orthogonal direct sum $\mathfrak{h}_3 = \mathfrak{b} \oplus \mathfrak{u}$ where \mathfrak{b} is a commutative subalgebra, \mathfrak{u} is a commutative ideal and where the linear transformation $ad(b)$ is skew-adjoint for every $b \in \mathfrak{b}$.

Let g be a Lorentz metric one of the left invariant Lorentz metrics g_1, g_2, g_3 . If $g(\cdot, \cdot)$ is nondegenerate, $e_0 \notin \mathfrak{u}$ such that

$$g(e_0, \mathfrak{u}) = 0, \quad \mathfrak{h}_3 = \mathbb{R}e_0 \oplus \mathfrak{u}.$$

We can obtain

$$ad_{e_0}^* e_0 = 0 \quad \text{and} \quad \forall y \in \mathfrak{u} \quad ad_{e_0}^* y \in \mathfrak{u}.$$

So, (1) equation is

$$\begin{cases} \dot{x}_1 = -\frac{\langle [e_0, x], x \rangle}{\langle e_0, e_0 \rangle} = -\frac{\langle Sx, x \rangle}{\langle e_0, e_0 \rangle} \\ \dot{x} = x_1 (ad_{e_0}^* x). \end{cases}$$

Here $L_{\gamma(t)}^{-1} \gamma'(t) = x_1 e_0 + x$, $x \in \mathfrak{u}$ and $S = \frac{1}{2} (ad e_0 + ad^* e_0)$.

It is found out that if (G, \langle, \rangle) is a flat Lorentz 2-step nilpotent Lie group, the restriction of \langle, \rangle to the of \mathfrak{b} is degenerate by [3].

Completeness of Lorentz Metric:

Theorem: Let (H_3, g_1) is given with the Lorentz metric

$$g_1 = -dx^2 + dy^2 + (xdy + dz)^2.$$

The left invariant Lorentz metric g_1 is geodesically complete.

Proof: Let our basis be

$$e_1 = \frac{\partial}{\partial z}, \quad e_2 = \frac{\partial}{\partial y} - x \frac{\partial}{\partial z}, \quad e_3 = \frac{\partial}{\partial x}$$

and we have

$$g_1(e_1, e_1) = g_1(e_2, e_2) = -g_1(e_3, e_3) = 1.$$

We can find easily on the basis $\{e_2, e_3\}$ nontrivial adjoint maps as

$$ad_{e_2}^* e_1 = -e_3, \quad ad_{e_3}^* e_1 = -e_2.$$

The ideal \mathbf{u} is generated by e_1 and is non-null. So we may choose $e_2, e_3 \notin \mathbf{u}$ such that

$$g_1(e_2, \mathbf{u}) \Big|_{\mathbf{u}} = 0, \quad g_1(e_3, \mathbf{u}) \Big|_{\mathbf{u}} = 0 \quad \text{and} \quad \mathbf{h}_3 = \mathbb{R}e_0 \oplus \mathbf{u}.$$

With this choice we say $\{e_1, e_2, e_3\}$ is an orthogonal basis of \mathbf{h}_3 and we obtain

$$[e_2, e_3] = e_1, \quad [e_3, e_1] = 0, \quad [e_2, e_1] = 0.$$

Then, we get easily

$$ad^* e_1 = \begin{pmatrix} 0 & -1 \\ -1 & 0 \end{pmatrix}$$

If we take $L_{\gamma(t)}^{-1} \gamma'(t) = x_1 e_1 + x_2 e_2 + x_3 e_3$, our (1) equations are

$$\begin{aligned} \dot{x}_1 &= 0 \\ \dot{x}_2 &= -x_1 x_3 \\ \dot{x}_3 &= -x_1 x_2. \end{aligned}$$

So the metric g_1 is geodesically complete.

Theorem: Let (H_3, g_2) is given with the Lorentz metric

$$g_2 = dx^2 + dy^2 - (xdy + dz)^2.$$

The left invariant Lorentz metric g_2 is geodesically complete.

Proof: Let our basis be

$$e_1 = \frac{\partial}{\partial x}, \quad e_2 = x \frac{\partial}{\partial z} - \frac{\partial}{\partial y}, \quad e_3 = \frac{\partial}{\partial z}$$

and we have

$$g_2(e_1, e_1) = g_2(e_2, e_2) = -g_2(e_3, e_3) = 1.$$

We can find easily on the basis $\{e_1, e_2\}$ nontrivial adjoint maps as

$$ad_{e_2}^* e_3 = -e_1, \quad ad_{e_1}^* e_3 = e_2.$$

The ideal \mathbf{u} is generated by e_3 and is non-null. So we may choose $e_1, e_2 \notin \mathbf{u}$ such that

$$g_1(e_1, \mathbf{u})|_{\mathbf{u}} = 0, \quad g_1(e_2, \mathbf{u})|_{\mathbf{u}} = 0 \quad \text{and} \quad \mathbf{h}_3 = \mathbb{R}e_0 \oplus \mathbf{u}.$$

With this choice we say $\{e_1, e_2, e_3\}$ is an orthogonal basis of \mathbf{h}_3 and we obtain

$$[e_1, e_2] = e_3, \quad [e_3, e_1] = 0, \quad [e_2, e_3] = 0.$$

Then, we get easily

$$ad^* e_3 = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$$

If we take $L_{\gamma(t)}^{-1} \gamma'(t) = x_1 e_1 + x_2 e_2 + x_3 e_3$, our (1) equations are

$$\begin{aligned} \dot{x}_1 &= x_1 x_3 \\ \dot{x}_2 &= -x_2 x_3 \\ \dot{x}_3 &= 0. \end{aligned}$$

So the metric g_2 is complete.

Theorem: Let (H_3, g_3) is given with the Lorentz metric

$$g_3 = dx^2 + (xdy + dz)^2 - ((1-x)dy - dz)^2.$$

The left invariant Lorentz metric g_3 is geodesically complete.

Proof: It is proved that g_3 is flat by Rahmani. Since g_3 is flat, the ideal \mathbf{u} is degenerate. We can choose a basis

$$e_1 = \frac{\partial}{\partial x}, \quad e_2 = \frac{\partial}{\partial y} + (1-x) \frac{\partial}{\partial z}, \quad e_3 = \frac{\partial}{\partial y} - x \frac{\partial}{\partial z}$$

and we have

$$g_3(e_1, e_1) = g_3(e_2, e_2) = -g_3(e_3, e_3) = 1.$$

The ideal \mathbf{u} is generated by $e_2 - e_3$ and is null. So we can choose a new basis $\left\{x = \frac{1}{\sqrt{2}}(e_2 - e_3), y = \frac{1}{\sqrt{2}}(e_3 + e_2), z = e_1\right\}$. With this basis we can obtain

$$[y, z] = 2x = v$$

and

$$g_3(x, y) \Big|_{\mathbf{u}} = 1 \text{ and } g_3(z, z) \Big|_{\mathbf{u}} = 1.$$

We can get

$$ad_v^* y = z, \quad ad_z^* y = -x$$

and

$$ad^* y = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$$

If we take $L_{\gamma(t)^{-1}}^* \gamma'(t) = x_1 v + x_2 y + x_3 z$, our (1) equations are

$$\begin{aligned} \dot{x}_1 &= 0 \\ \dot{x}_2 &= x_1 x_3 \\ \dot{x}_3 &= -x_1 x_2. \end{aligned}$$

So the metric g_3 is complete.

Corollary: Let (H_3, g) be 3-dimensional Heisenberg group, g is left invariant Lorentz metric one of the g_1, g_2, g_3 on H_3 . g is geodesically complete on 3-dimensional Heisenberg group.

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