

Jordan Decomposition and its Uniqueness of Signed Lattice Measure

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Abstract

In this paper, we define a signed lattice measure on σ -Algebras, as well as give the definition of positive and negative lattice, positive and negative parts of ν and mutually singular lattice measures. Herein, we show that the Hann Decomposition theorem decomposes any space X into a positive lattice A and a negative lattice B such that $A \vee B = X$ and $A \wedge B = \phi$. Also we prove the Jordan Decomposition theorem and its uniqueness of signed lattice measure.

Keywords: Signed Lattice Measure, Hann-Decomposition, Measurable Space

1. INTRODUCTION

In this paper, we illustrate the Jordan Decomposition theorem and its uniqueness of signed lattice measure. We consider a set X be a nonempty arbitrary

set and L a lattice subsets of X . All lattices considered throughout the paper will contain ϕ and X .

We define a signed lattice measure on σ -Algebras, and we show that the lattice Hann Decomposition theorem decomposes any set X into a Positive lattice A and a negative lattice B such that $A \vee B = X$ and $A \wedge B = \phi$. We also prove the Jordan Decomposition Theorem and its uniqueness.

In section 2, according to [4][5] the definition of a signed lattice measure on σ -Algebras, as well as give the definition of positive and negative lattice. Also we define positive and negative parts of ν and mutually singular lattice measures. Further, we show that any countable union of positive lattice is a positive lattice as well as if E is a measurable lattice such that $0 < \nu(E) < \infty$, then there is a positive lattice $A \leq E$ with $\nu(A) > 0$ [2]. Also we show the Hann Decomposition theorem for signed lattice measure [3]. Finally illustrate the Jordan Decomposition theorem and its uniqueness for a signed lattice measure.

2. PRELIMINARIES

In this section, we shall briefly review the well-known facts about lattice theory due to Birkhoff [1], propose an extension lattice, and investigate its properties. (L, \wedge, \vee) is called a lattice if it is enclosed under operations \wedge and \vee and satisfies, for any elements x, y, z , in L :

(L1) the commutative law: $x \wedge y = y \wedge x$ and $x \vee y = y \vee x$.

(L2) the associative law: $x \wedge (y \wedge z) = (x \wedge y) \wedge z$ and $x \vee (y \vee z) = (x \vee y) \vee z$.

(L3) the absorption law: $x \vee (y \wedge x) = x$ and $x \wedge (y \vee x) = x$.

Hereafter, the lattice (L, \wedge, \vee) will often be written as L for simplicity.

A mapping h from a lattice L to another lattice L^1 is called a lattice-homomorphism, if it satisfies

$h(x \wedge y) = h(x) \wedge h(y)$ and $h(x \vee y) = h(x) \vee h(y)$, $\forall x, y \in L$.

If h is a bijection, that is, h is one-to-one and onto, it is called a lattice isomorphism, and in this case, L^1 is said to be lattice-isomorphic to L .

A lattice (L, \wedge, \vee) is called distributive if, for any x, y, z , in L .

(L4) the distributive law holds:

$x \vee (y \wedge z) = (x \vee y) \wedge (x \vee z)$ and $x \wedge (y \vee z) = (x \wedge y) \vee (x \wedge z)$.

A lattice L is called complete if, for any subset A of L , L contains the supremum $\vee A$ and the infimum $\wedge A$. If L is complete, then L itself includes the maximum and minimum elements which are often denoted by 1 and 0 or I and O respectively [1].

A distributive lattice is called a Boolean lattice if for any element x in L , there exists a unique complement x^c such that

$x \vee x^c = 1$ (L5) the law of excluded middle

$$x \wedge x^c = 0 \quad (\text{L6) the law of non-contradiction}$$

Let L be a lattice and $\epsilon: L \rightarrow L$ be an operator. Then ϵ is called a lattice complement in L if the following conditions are satisfied.

$$(\text{L5) and (L6);} \quad \forall x \in L, x \vee x^c = 1 \text{ and } x \wedge x^c = 0,$$

$$(\text{L7) the law of contrapositive;} \quad \forall x, y \in L, x \leq y \text{ implies } x^c \geq y^c,$$

$$(\text{L8) the law of double negation;} \quad \forall x \in L, (x^c)^c = x.$$

Throughout this paper, we consider lattices as complete lattices which obey (L1) - (L8) except for (L6) the law of non-contradiction.

2.1 Definition. Unless otherwise stated, X is the entire set and L is a lattice of any subsets of X . If a lattice L satisfies the following conditions, then it is called a lattice σ -Algebra;

$$(1) \quad \forall h \in L, h^c \in L$$

$$(2) \quad \text{if } h_n \in L \text{ for } n = 1, 2, 3, \dots, \text{ then } \bigvee_{n=1}^{\infty} h_n \in L.$$

We denote $\sigma(L)$, as the lattice σ -Algebra generated by L .

2.2 Definition. The ordered pair $(X, \sigma(L))$ is said to be lattice measurable space

2.3 Definition. If $m: \sigma(L) \rightarrow \mathbb{R} \cup \{\infty\}$ satisfies the following properties, then m is called a lattice measure on the lattice σ -Algebra $\sigma(L)$.

$$(1) \quad m(\phi) = m(0) = 0.$$

$$(2) \quad \forall h, g \in \sigma(L), \text{ s.t. } m(h), m(g) \geq 0; h \leq g \Rightarrow m(h) \leq m(g).$$

$$(3) \quad \forall h, g \in \sigma(L): m(h \vee g) + m(h \wedge g) = m(h) + m(g).$$

$$(4) \quad \text{If } h_n \in \sigma(L), n \in \mathbb{N} \text{ such that } h_1 \leq h_2 \leq \dots \leq h_n \leq \dots, \text{ then } m\left(\bigvee_{n=1}^{\infty} h_n\right) = \lim m(h_n).$$

Let m_1 and m_2 be lattice measures define on the same lattice σ -Algebra $\sigma(L)$.

If one of them is finite, the set function $m(E) = m_1(E) - m_2(E)$, $E \in \sigma(L)$ is well defined and countably additive on $\sigma(L)$. However, it is necessarily nonnegative; it is called a signed lattice measure.

2.4 Definition. By a signed lattice measure on the lattice measurable space $(X, \sigma(L))$ we mean $v: \sigma(L) \rightarrow \mathbb{R} \cup \{\infty\}$ or $\mathbb{R} \cup \{-\infty\}$, satisfying the following properties :

$$(1) \quad v(\phi) = v(0) = 0.$$

$$(2) \quad \forall h, g \in \sigma(L) \text{ s.t. } v(h), v(g) \geq 0; h \leq g \Rightarrow v(h) \leq v(g).$$

$$\forall h, g \in \sigma(L) \text{ s.t. } v(h), v(g) \leq 0; h \leq g \Rightarrow v(g) \leq v(h).$$

$$(3) \quad \forall h, g \in \sigma(L): v(h \vee g) + v(h \wedge g) = v(h) + v(g).$$

(4) if $h_n \in \sigma(L)$, $n \in \mathbb{N}$ such that $h_1 \leq h_2 \leq \dots \leq h_n \leq \dots$, then $v(\bigvee_{n=1}^{\infty} h_n) = \lim v(h_n)$.

This is meant in the sense that if the left-hand side is finite, the limit on the right-hand side is convergent, and if the left-hand side is $\pm\infty$, then the limit on the right-hand side diverges accordingly.

2.5 Definition. A set A is said to be lattice measurable set or measurable lattice, if A belongs to $\sigma(L)$.

2.6 Definition. A is a positive lattice if for any lattice measurable set E in A , $v(E) \geq 0$; Similarly, B is a negative lattice if for any lattice measurable set E in B , $v(E) \leq 0$.

2.7 Definition. Let $(X, \sigma(L))$ be a lattice measurable space and v be a signed lattice measure defined on $\sigma(L)$ then there exists a positive set A and a negative set B such that $A \vee B = X$ and $A \wedge B = \phi$. We denote lattice Hann decomposition by $\{A, B\}$

2.8 Definition. (Positive and Negative parts of v):

Let $(X, \sigma(L))$ be a lattice measurable space and v be a signed lattice measure defined on $\sigma(L)$. If we define v^+ and v^- on for all $E \in \sigma(L)$ such that $v^+(E) = v(E \wedge A)$ and $v^-(E) = -v(E \wedge B)$ with $v = v^+ - v^-$. Then v^+ and v^- are called the positive and the negative parts of a signed measure lattice respectively. Where $\{A, B\}$ be a Hann decomposition on X .

2.9 Definition. (Mutually singular lattice measures):

Let $(X, \sigma(L))$ be a lattice measurable space and v be a signed lattice measure defined on $\sigma(L)$. If we define two measures v_1 and v_2 are said to be mutually singular lattice measures with each other to be denoted by $v_1 \perp v_2$. If there exists two measurable lattices A and B such that $A \vee B = X$ and $A \wedge B = \phi$ and $v_1(A) = v_2(B) = 0$.

3. JORDAN-DECOMPOSTION AND ITS UNIQUENESS OF SIGNED LATTICE MEASURE

3.1 Lemma: Every measurable sublattice of a positive lattice is a positive lattice and a countable union of positive lattice is a positive lattice

Proof : Part (i): Let A be a positive lattice and P be a sublattice of A Let E be a

measurable sub lattice of P implies E be a measurable sublattice of A, By definition of a positive lattice $v(E) \geq 0$. Therefore P is a positive lattice.

Part (ii): To prove that countable union of positive lattice is a positive lattice.

Let $\langle A_i \rangle$ be a countable collection of positive lattices such that $A_1 \leq A_2 \leq \dots \leq A_n \leq \dots$

by the definition4, $v(\bigvee_{i=1}^{\infty} A_i) = \lim v(A_i)$. Put $A = \bigvee_{i=1}^{\infty} A_i$

Now each A_i is a positive lattice implies each A_i is a measurable lattice by the definition3,

$A = \bigvee_{i=1}^{\infty} A_i$ is measurable. Therefore A is measurable lattice.

To prove that A is positive.

Let E be a measurable sublattice of A that is $E \leq A$ implies $E \leq \bigvee_{i=1}^{\infty} A_i$

Define $E_n = E \wedge A_n \wedge A_{n-1}^c \wedge \dots$ implies $E_n \leq E$ and $E_n \leq A_n$ for all $n= 1, 2, 3, \dots$

implies E_n is measurable and E_n is positive lattice therefore $v(E_n) \geq 0$ for all $n=$

$1, 2, 3, \dots$ also in particular $E_n \leq E_{n+1}$ for all n. Evidently $E = \bigvee_{n=1}^{\infty} E_n$.

It implies $v(E) = v(\bigvee_{n=1}^{\infty} E_n) = \lim v(E_n) \geq 0$ implies $v(E) \geq 0$. Therefore A is a positive lattice.

3.2 Lemma: Let E be a measurable lattice such that $0 < v(E) < \infty$ then there is a positive lattice $A \leq E$ with $v(A) > 0$.

Proof: If E is a positive lattice we take $A = E$. Therefore $v(A) > 0$ and A is a positive lattice.

It is given that $0 < v(E) < \infty$ Evidently $v(E) > 0$ and finite. Suppose E contains a sublattice of negative measures, let n_1 be the smallest positive integer such that

there exists a measurable lattice $E_1 < E$ with $v(E_1) < \frac{-1}{n_1}$. If $E \wedge E_1^c$ is not already

a positive lattice, let n_2 be the least positive integer such that there exists

measurable lattice $E_2 < E \wedge E_1^c$ such that $v(E_2) < \frac{-1}{n_2}$. Proceeding inductively, if

$E \wedge (\bigvee_{j=1}^{k-1} E_j)^c$ is not already a positive lattice let n_k be the smallest positive integer

$> n_{k-1}$ such that there exists a measurable lattice $E_k < E \wedge (\bigwedge_{j=1}^{k-1} E_j^c)$ such that

$v(E_k) < \frac{-1}{n_k}$. Continuing like this let $A = E \wedge (\bigwedge_{k=1}^{\infty} E_k^c)$

implies $A = E - (\bigvee_{k=1}^{\infty} E_k)$ implies $E = A \vee (\bigvee_{k=1}^{\infty} E_k)$ implies $v(E) = v(A) + v(\bigvee_{k=1}^{\infty} E_k)$

implies $v(E) = v(A) + \lim v(E_k)$ ----- (1)

As $v(E)$ is finite it implies $\lim v(E_k)$ is finite and negative from (1) $v(A) = v(E) > 0$

Therefore $v(A) > 0$. To prove A is a positive lattice, $A = E \wedge (\bigwedge_{k=1}^{\infty} E_k^c)$

Since E is measurable it implies A is measurable. Let P be a measurable sublattice of A .

$P \leq A$ implies $P \leq E - (\bigvee_{k=1}^{\infty} E_k)$ implies $P \leq E$ and $P \leq (\bigvee_{k=1}^{\infty} E_k)$

implies $P \leq E$ and $P \not\leq E_k$ for any k . Therefore P is not a sublattice of a negative measure that is A does not have any sublattice of negative measures. Therefore A is a positive lattice. Therefore E has a positive lattice A such that $v(A) > 0$.

3.3 THEOREM (Lattice Hann-Decomposition):

Let $(X, \sigma(L))$ be a lattice measurable space and v be a signed lattice measure defined on $\sigma(L)$ then there exists a positive set A and a negative set B such that $A \vee B = X$ and $A \wedge B = \phi$.

Proof: Without loss of generality $+\infty$ is a maximum value omitted by v

Define $\lambda = \sup \{v(A_i) / A_i \text{ is a positive lattice}\}$ since $v(\phi) = 0$ implies $\lambda \geq 0$

Let $\langle A_i \rangle$ be a sequence of positive lattices such that $\lambda = \lim v(A_i)$ Put $A = \bigvee_{i=1}^{\infty} A_i$.

By Lemma 1 A is a positive lattice. Therefore $v(A) \leq \lambda$ ----- (1)

Also each $A_i \leq A$ implies $A - A_i \leq A$ implies $v(A - A_i) \geq 0$ Evidently $A = A_i \vee (A - A_i)$

implies $v(A) = v(A_i) + v(A - A_i)$ implies $v(A) \geq v(A_i)$ for all i , implies $v(A) \geq \lambda$ ----- (2)

From (1) and (2) $v(A) = \lambda$. Put $B = A^c$, to prove that B is a negative lattice.

Let us suppose that B has a positive measure lattice E (say). Now E is a positive lattice and A is a positive lattice implies $A \vee E$ is a positive lattice implies $v(A \vee E) \leq \lambda$ implies $v(A) + v(E) \leq \lambda$ implies $\lambda + v(E) \leq \lambda$ implies $v(E) \leq 0$. This is a contradiction. Therefore B has no sublattice of positive measures. This implies B is a negative lattice.

Hence by above, we conclude that "Let X be an entire set. Then by above theorem we can find a positive lattice A and a negative lattice $B (=A^c)$. By the lattice measurability $A \vee B = X$ and $A \wedge B = \phi$. These characteristics provide the

following $X = A \cup B$ and $\phi = A \cap B$.

3.4 THEOREM (Jordan–Decomposition on Signed Lattice Measure):

Let $(X, \sigma(L))$ be a lattice measurable space and ν be a signed lattice measure defined on $\sigma(L)$ then there exists a unique pair of mutually singular lattice measures ν^+ and ν^- with $\nu = \nu^+ - \nu^-$.

Proof: Let $\{A, B\}$ be a Hann decomposition on X then there exists a positive lattice A and a negative lattice B such that $A \vee B = X$ and $A \wedge B = \phi$

Now for all $E \in \sigma(L)$. Define $\nu^+(E) = \nu(E \wedge A)$ (1)

$\nu^-(E) = -\nu(E \wedge B)$ (2)

from(1) $\nu^+(B) = \nu(B \wedge A) = \phi$

from(2) $\nu^-(A) = -\nu(A \wedge B) = \phi$

therefore ν^+ and ν^- are mutually singular lattice measures. Also $E=(E \wedge A) \vee (E \wedge B)$
 $\nu(E) = \nu(E \wedge A) + \nu(E \wedge B)$, $\nu(E) = \nu^+(E) - \nu^-(E)$ therefore $\nu = \nu^+ - \nu^-$ (3)

UNIQUENESS :

Let $\{A, B\}$ and $\{A^1, B^1\}$ be two lattice Hann decompositions on X

$$A \vee A^1 = (A \wedge A^1) \vee (A \Delta A^1) \dots\dots\dots (4)$$

$$\text{where } (A \Delta A^1) = (A - A^1) \vee (A^1 - A) \dots\dots\dots (5)$$

$$\nu(A \vee A^1) = \nu(A \wedge A^1) + \nu(A \Delta A^1) \dots\dots\dots (6)$$

$$\text{where } \nu(A \Delta A^1) = \nu(A - A^1) + \nu(A^1 - A) \dots\dots\dots (7)$$

Clearly $A - A^1 \leq A \wedge B^1$ Now $A - A^1 \leq A$ and $A - A^1 \leq B^1$

It implies $\nu(A - A^1) \geq 0$ and $\nu(A - A^1) \leq 0$ (since A is a positive lattice and B^1 is a negative lattice) therefore $\nu(A - A^1) = 0$ similarly $\nu(A^1 - A) = 0$ therefore from (7) $\nu(A \Delta A^1) = 0$

$$\text{from (6) } \nu(A \vee A^1) = \nu(A \wedge A^1) \dots\dots\dots (8)$$

Also evidently for all $E \in \sigma(L)$

$$\nu(E \wedge A \wedge A^1) \leq \nu(E \wedge A) \leq \nu(E \wedge A \wedge A^1) \dots\dots\dots (9)$$

$$\nu(E \wedge A \wedge A^1) \leq \nu(E \wedge A^1) \leq \nu(E \wedge A \wedge A^1) \dots\dots\dots (10)$$

From (9) and (10) we have $\nu(E \wedge A) = \nu(E \wedge A^1)$.

Therefore ν^+ is unique. From (3) $\nu = \nu^+ - \nu^-$ implies ν^- is unique. Therefore there is only one such pair of mutually singular lattice measures.

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