

About a Solution of $\frac{\partial^3 x^\nu}{\partial x^\mu \partial x^\mu \partial x^\mu} = M : \{M = 0, M > 0\}$ in Tensor Satisfying Binary Law

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Abstract

I handle the case that all coordinate systems satisfies Binary Law in this article. If all coordinate systems satisfies Binary Law, I have already reported that $\frac{\partial^3 x^\mu}{\partial x^\nu \partial x^\nu \partial x^\nu} = \frac{\partial x_\mu}{\partial x^\nu} = \left(\frac{\partial x_\nu}{\partial x_\nu} - x_\mu \frac{1}{2} \left(\frac{\partial g^{\nu\mu}}{\partial x^\nu} \right) \right) = M$ is established. I report it about the solution of the equation of $\frac{\partial^3 x^\nu}{\partial x^\mu \partial x^\mu \partial x^\mu} = M : \{M = 0, M > 0\}$ in this article. I report that these provided solutions can decide metric $g^{\mu\nu}$ of the space. Metric $g^{\mu\nu}$ of the space here is quantity of the filed. And I report that these solutions provided more can decide power F_μ . Power F_μ here is quantity of the filed.

Keywords: Tensor, Covariant Derivative

1 Introduction

If all coordinate systems satisfies Binary Law, I have already reported that $\frac{\partial^3 x^\mu}{\partial x^\nu \partial x^\nu \partial x^\nu} = \frac{\partial x_\mu}{\partial x^\nu} = \left(\frac{\partial x_\nu}{\partial x_\nu} - x_\mu \frac{1}{2} \left(\frac{\partial g^{\nu\mu}}{\partial x^\nu} \right) \right) = M$ is established.[2] I report it about the solution of the equation of $\frac{\partial^3 x^\nu}{\partial x^\mu \partial x^\mu \partial x^\mu} = M : \{M = 0, M > 0\}$ in this article. I report that two kinds of $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = -Mx^\nu$ and $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = \frac{M}{x^\nu}$ can rewrite equation $\frac{\partial^3 x^\nu}{\partial x^\mu \partial x^\mu \partial x^\mu} = M$ each. If $M = 0$ is established, I get $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = 0$ from $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = -Mx^\nu$, $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = \frac{M}{x^\nu}$. On the other hand, two kinds of $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = -Mx^\nu$, $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = \frac{M}{x^\nu}$ can exist if $M > 0$ is established.

2 Definition

Definition 1 $\bar{x}^\mu \neq x^\mu, \bar{x}^\nu \neq x^\nu, \bar{x}^\mu = x^\nu, \bar{x}^\nu = x^\mu$ is established.[1]

I named $\bar{x}^\mu \neq x^\mu, \bar{x}^\nu \neq x^\nu, \bar{x}^\mu = x^\nu, \bar{x}^\nu = x^\mu$ "Binary Law".[1]

Definition 2 If $\bar{x}^\mu \neq x^\mu, \bar{x}^\nu \neq x^\nu, \bar{x}^\mu = x^\nu, \bar{x}^\nu = x^\mu$ is established, $x^\mu = x_\nu$ is established.[1]

Definition 3 If $\bar{x}^\mu \neq x^\mu, \bar{x}^\nu \neq x^\nu, \bar{x}^\mu = x^\nu, \bar{x}^\nu = x^\mu$ is established, $x^\nu = x_\mu$ is established.[1]

Definition 4 If $\bar{x}^\mu \neq x^\mu, \bar{x}^\nu \neq x^\nu, \bar{x}^\mu = x^\nu, \bar{x}^\nu = x^\mu$ is established, $x_\nu = -x_\mu$ is established.[1]

Definition 5 If $\bar{x}^\mu \neq x^\mu, \bar{x}^\nu \neq x^\nu, \bar{x}^\mu = x^\nu, \bar{x}^\nu = x^\mu$ is established, $x^\nu = -x^\mu$ is established.[1]

Definition 6 If $\bar{x}^\mu \neq x^\mu, \bar{x}^\nu \neq x^\nu, \bar{x}^\mu = x^\nu, \bar{x}^\nu = x^\mu$ is established, $\frac{\partial M}{\partial x^\mu} = 0$ is established. "M" expresses $\frac{\partial^3 x^\nu}{\partial x^\mu \partial x^\mu \partial x^\mu} = M$. [2]

Definition 7 If $\bar{x}^\mu \neq x^\mu, \bar{x}^\nu \neq x^\nu, \bar{x}^\mu = x^\nu, \bar{x}^\nu = x^\mu$ is established, $\frac{\partial^3 x^\mu}{\partial x^\nu \partial x^\nu \partial x^\nu} = \frac{\partial x_\mu}{\partial x^\nu} = \frac{\partial x_\nu}{\partial x^\mu} - x_\mu \frac{1}{2} \left(\frac{\partial g^{\nu\mu}}{\partial x^\nu} \right) = M$ is established. [2]

Definition 8 $\delta = \sqrt{(x_2 - x_1)^2 + (y_2 - y_1)^2}$ is established for distance δ between two points of $P_1(x_1, y_1)$ and $P_2(x_2, y_2)$. [3]

Definition 9 Central (x_0, y_0) , the equation of the circle of radius δ are expressed in $(x - x_0)^2 + (y - y_0)^2 = \delta^2$. [3]

Definition 10 $E = mc^2$ is established. [4]

"E" expresses Energy, "m" expresses Mass, and "c" expresses Speed of light.

Definition 11 $W(A \rightarrow B) = \int_A^B \vec{F} \cdot d\vec{r}$ is established. [5]

"W" expresses Work, \vec{F} expresses External force vector, and \vec{r} expresses Displacement vector.

Definition 12 $dx_i = \frac{\partial x'^j}{\partial x^i} dx'_j$ is established.

Definition 13 $ds^2 = (e^{\vec{i}} \cdot e^{\vec{j}}) dx'_i dx'_j = g^{ij} dx'_i dx'_j$ is established. [6]

Definition 14 $y(x) = C[2] + xC[1]$ is established as a solution of the equations of $\frac{d^2 y}{dx dx} = 0$. [7]

y is function $y=f(x)$ which assumes x an independent variable.

Definition 15 $y(x) = C[2]\cos\left(\frac{\sqrt{M}x}{\sqrt{2}}\right) + C[1]\sin\left(\frac{\sqrt{M}x}{\sqrt{2}}\right)$ is established as a solution of the equations of $\frac{d^2 y}{dx dx} = -\frac{M^*y}{2} : \{M > 0\}$. [7]

y is function $y=f(x)$ which assumes x an independent variable.

Hypothesis 1 $M \propto m$, $M = \epsilon m$ is established.

M expresses $\frac{\partial^3 x^\nu}{\partial x^\mu \partial x^\mu \partial x^\mu} = M$, ϵ expresses Proportional constant, and m expresses Mass.

3 About relations with solution of $\frac{\partial^3 x^\nu}{\partial x^\mu \partial x^\mu \partial x^\mu} = M : \{M = 0, M > 0\}$ and metric $g^{\mu\nu}$ of the space

Proposition 1 When all coordinate systems satisfies Binary Law, $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = -Mx^\nu$, $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = \frac{M}{x^\nu}$ can rewrite $\frac{\partial^3 x^\nu}{\partial x^\mu \partial x^\mu \partial x^\mu} = M$.

Proof: μ, ν -reverses Definition7 and gets

$$\frac{\partial^3 x^\nu}{\partial x^\mu \partial x^\mu \partial x^\mu} = M. \tag{1}$$

I get

$$\begin{aligned} \frac{\partial^3 \dot{x}^1}{\partial x^1 \partial x^1 \partial x^1} = M, \quad \frac{\partial^3 \dot{x}^1}{\partial x^2 \partial x^2 \partial x^2} = M, \\ \frac{\partial^3 \dot{x}^2}{\partial x^1 \partial x^1 \partial x^1} = M, \quad \frac{\partial^3 \dot{x}^2}{\partial x^2 \partial x^2 \partial x^2} = M \end{aligned} \tag{2}$$

from (1) if I assume a dimensional number 2. Furthermore, I get

$$\begin{aligned} \int \frac{\partial^3 \dot{x}^1}{\partial x^1 \partial x^1 \partial x^1} \partial x^1 = M \int \partial x^1, \quad \int \frac{\partial^3 \dot{x}^1}{\partial x^2 \partial x^2 \partial x^2} \partial x^2 = M \int \partial x^2, \\ \int \frac{\partial^3 \dot{x}^2}{\partial x^1 \partial x^1 \partial x^1} \partial x^1 = M \int \partial x^1, \quad \int \frac{\partial^3 \dot{x}^2}{\partial x^2 \partial x^2 \partial x^2} \partial x^2 = M \int \partial x^2 \end{aligned} \tag{3}$$

in consideration of Definition6 for (2). I get

$$\begin{aligned} \frac{\partial^2 \dot{x}^1}{\partial x^1 \partial x^1} = Mx^1, \quad \frac{\partial^2 \dot{x}^1}{\partial x^2 \partial x^2} = Mx^2, \\ \frac{\partial^2 \dot{x}^2}{\partial x^1 \partial x^1} = Mx^1, \quad \frac{\partial^2 \dot{x}^2}{\partial x^2 \partial x^2} = Mx^2 \end{aligned} \tag{4}$$

from (3). I get

$$\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = Mx^\mu \quad (5)$$

from (4). Two next

$$\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = -Mx^\nu, \quad (6)$$

$$\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = Mx_\nu = \frac{M}{x^\nu} \quad (7)$$

can rewrite (5) each using Definition2, Definition5. I get (7) as $x_\nu = \frac{1}{x^\nu}$ here.

– End Proof –

Proposition 2 *When all coordinate systems satisfies Binary Law, $x^1 = x^2, \dot{x}^1 = \dot{x}^2$ is established for $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = -Mx^\nu, \frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = \frac{M}{x^\nu}$ if the number of the dimensions is 2.*

Proof: I get

$$\begin{aligned} \frac{\partial^2 \dot{x}^1}{\partial x^1 \partial x^1} &= -M\dot{x}^1, \quad \frac{\partial^2 \dot{x}^1}{\partial x^2 \partial x^2} = -M\dot{x}^1, \\ \frac{\partial^2 \dot{x}^2}{\partial x^1 \partial x^1} &= -M\dot{x}^2, \quad \frac{\partial^2 \dot{x}^2}{\partial x^2 \partial x^2} = -M\dot{x}^2, \end{aligned} \quad (8)$$

$$\begin{aligned} \frac{\partial^2 \dot{x}^1}{\partial x^1 \partial x^1} &= \frac{M}{\dot{x}^1}, \quad \frac{\partial^2 \dot{x}^1}{\partial x^2 \partial x^2} = \frac{M}{\dot{x}^1}, \\ \frac{\partial^2 \dot{x}^2}{\partial x^1 \partial x^1} &= \frac{M}{\dot{x}^2}, \quad \frac{\partial^2 \dot{x}^2}{\partial x^2 \partial x^2} = \frac{M}{\dot{x}^2} \end{aligned} \quad (9)$$

from (6),(7) if I assume a dimensional number 2. I get

$$\frac{\partial^2 \dot{x}^1}{\partial x^1 \partial x^1} = \frac{\partial^2 \dot{x}^1}{\partial x^2 \partial x^2}, \quad \frac{\partial^2 \dot{x}^2}{\partial x^1 \partial x^1} = \frac{\partial^2 \dot{x}^2}{\partial x^2 \partial x^2} \quad (10)$$

from (8),(9). I get

$$\frac{\partial^2 \dot{x}^1}{\partial x^1 \partial x^1} = \frac{\partial^2 \dot{x}^1}{\partial x^1 \partial x^1} \quad (false), \quad \frac{\partial^2 \dot{x}^2}{\partial x^1 \partial x^1} = \frac{\partial^2 \dot{x}^2}{\partial x^1 \partial x^1} \quad (false) \quad (11)$$

from (10) if I assume establishment of $x^1 = x^2$ (*false*). Because (11) isn't established,

$$x^1 = x^2 \quad (12)$$

is established. I get

$$\frac{d^2 \dot{x}^1}{dx^1 dx^1} = -M \dot{x}^1, \tag{13}$$

$$\frac{d^2 \dot{x}^2}{dx^1 dx^1} = -M \dot{x}^2, \tag{14}$$

$$\frac{d^2 \dot{x}^1}{dx^1 dx^1} = \frac{M}{\dot{x}^1}, \tag{15}$$

$$\frac{d^2 \dot{x}^2}{dx^1 dx^1} = \frac{M}{\dot{x}^2} \tag{16}$$

in consideration of (12) for (8),(9). I get

$$\frac{d^2 \dot{x}^2}{dx^1 dx^1} = -M \dot{x}^2 \text{ (false)}, \frac{d^2 \dot{x}^2}{dx^1 dx^1} = \frac{M}{\dot{x}^2} \text{ (false)} \tag{17}$$

from (13),(15) if I assume establishment of $\dot{x}^1 = \dot{x}^2$ (false). I get

$$\frac{d^2 \dot{x}^2}{dx^1 dx^1} = \frac{d^2 \dot{x}^2}{dx^1 dx^1} \text{ (false)} \tag{18}$$

from (14),(16),(17). Because (18) isn't established,

$$\dot{x}^1 = \dot{x}^2 \tag{19}$$

is established. I get

$$\frac{d^2 \dot{x}^1}{dx^1 dx^1} = -M \dot{x}^1, \tag{20}$$

$$\frac{d^2 \dot{x}^1}{dx^1 dx^1} = \frac{M}{\dot{x}^1} \tag{21}$$

in consideration of (19) for (13),(14),(15),(16). I get

$$\delta = \sqrt{(x^1)^2 + (x^2)^2} \tag{22}$$

as $x_1 = 0, y_1 = 0, x_2 \rightarrow x^1, y_2 \rightarrow x^2$ for Definition8. I get

$$x^1 = \frac{\delta}{\sqrt{2}} \tag{23}$$

in consideration of (12) for (22). I rewrite (20),(21) in consideration of (23) and get

$$\frac{d^2 \dot{x}^1}{d\delta d\delta} = -\frac{M\dot{x}^1}{2}, \quad (24)$$

$$\frac{d^2 \dot{x}^1}{d\delta d\delta} = \frac{M}{2\dot{x}^1}. \quad (25)$$

– End Proof –

Proposition 3 *When all coordinate systems satisfies Binary Law, and $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = -Mx^\nu$, $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = \frac{M}{x^\nu}$ is established more, $g^{11} = g^{12} = g^{21} = g^{22} = 2\frac{d\dot{x}^1}{dx^1} \frac{d\dot{x}^1}{dx^1}$: $g^{\mu\nu} = g^{11}, g^{12}, g^{21}, g^{22}$ is established if a dimensional number is 2.*

Proof: If a dimensional number is 2,

$$d\vec{r} = dx_1 \vec{e}^1 + dx_2 \vec{e}^2, \quad (26)$$

$$dx_1 = \frac{\partial x'^1}{\partial x^1} dx'_1 + \frac{\partial x'^2}{\partial x^1} dx'_2, \quad dx_2 = \frac{\partial x'^1}{\partial x^2} dx'_1 + \frac{\partial x'^2}{\partial x^2} dx'_2 \quad (27)$$

is established. I considered Definition12 here. I get

$$\begin{aligned} ds^2 &= d\vec{r} \cdot d\vec{r} = (\vec{e}^1 \cdot \vec{e}^1) dx_1 dx_1 + (\vec{e}^1 \cdot \vec{e}^2) dx_1 dx_2 \\ &\quad + (\vec{e}^2 \cdot \vec{e}^1) dx_2 dx_1 + (\vec{e}^2 \cdot \vec{e}^2) dx_2 dx_2 \\ &= dx_1 dx_1 + dx_2 dx_2 \end{aligned} \quad (28)$$

from (26). I selected it as $(\vec{e}^1 \cdot \vec{e}^1) = (\vec{e}^2 \cdot \vec{e}^2) = 1, (\vec{e}^1 \cdot \vec{e}^2) = (\vec{e}^2 \cdot \vec{e}^1) = 0$ here. I get

$$\begin{aligned} dx_1 dx_1 &= \frac{\partial x'^1}{\partial x^1} \frac{\partial x'^1}{\partial x^1} dx'_1 dx'_1 + \frac{\partial x'^2}{\partial x^1} \frac{\partial x'^1}{\partial x^1} dx'_2 dx'_1 \\ &\quad + \frac{\partial x'^1}{\partial x^1} \frac{\partial x'^2}{\partial x^1} dx'_1 dx'_2 + \frac{\partial x'^2}{\partial x^1} \frac{\partial x'^2}{\partial x^1} dx'_2 dx'_2, \\ dx_2 dx_2 &= \frac{\partial x'^1}{\partial x^2} \frac{\partial x'^1}{\partial x^2} dx'_1 dx'_1 + \frac{\partial x'^2}{\partial x^2} \frac{\partial x'^1}{\partial x^2} dx'_2 dx'_1 \\ &\quad + \frac{\partial x'^1}{\partial x^2} \frac{\partial x'^2}{\partial x^2} dx'_1 dx'_2 + \frac{\partial x'^2}{\partial x^2} \frac{\partial x'^2}{\partial x^2} dx'_2 dx'_2 \end{aligned} \quad (29)$$

from (27). I get

$$\begin{aligned} ds^2 &= \left(\frac{\partial x'^1}{\partial x^1} \frac{\partial x'^1}{\partial x^1} + \frac{\partial x'^1}{\partial x^2} \frac{\partial x'^1}{\partial x^2} \right) dx'_1 dx'_1 + \left(\frac{\partial x'^2}{\partial x^1} \frac{\partial x'^1}{\partial x^1} + \frac{\partial x'^2}{\partial x^2} \frac{\partial x'^1}{\partial x^2} \right) dx'_2 dx'_1 \\ &+ \left(\frac{\partial x'^1}{\partial x^1} \frac{\partial x'^2}{\partial x^1} + \frac{\partial x'^1}{\partial x^2} \frac{\partial x'^2}{\partial x^2} \right) dx'_1 dx'_2 + \left(\frac{\partial x'^2}{\partial x^1} \frac{\partial x'^2}{\partial x^1} + \frac{\partial x'^2}{\partial x^2} \frac{\partial x'^2}{\partial x^2} \right) dx'_2 dx'_2 \\ &= g^{11} dx'_1 dx'_1 + g^{21} dx'_2 dx'_1 + g^{12} dx'_1 dx'_2 + g^{22} dx'_2 dx'_2 \end{aligned} \quad (30)$$

from (28),(29),Definition13. I get

$$\begin{aligned} g^{11} &= \frac{\partial \dot{x}^1}{\partial x^1} \frac{\partial \dot{x}^1}{\partial x^1} + \frac{\partial \dot{x}^1}{\partial x^2} \frac{\partial \dot{x}^1}{\partial x^2}, g^{21} = \frac{\partial \dot{x}^2}{\partial x^1} \frac{\partial \dot{x}^1}{\partial x^1} + \frac{\partial \dot{x}^2}{\partial x^2} \frac{\partial \dot{x}^1}{\partial x^2}, \\ g^{12} &= \frac{\partial \dot{x}^1}{\partial x^1} \frac{\partial \dot{x}^2}{\partial x^1} + \frac{\partial \dot{x}^1}{\partial x^2} \frac{\partial \dot{x}^2}{\partial x^2}, g^{22} = \frac{\partial \dot{x}^2}{\partial x^1} \frac{\partial \dot{x}^2}{\partial x^1} + \frac{\partial \dot{x}^2}{\partial x^2} \frac{\partial \dot{x}^2}{\partial x^2} \end{aligned} \quad (31)$$

as $x'^1 \rightarrow \dot{x}^1, x'^2 \rightarrow \dot{x}^2$ for (30). I get

$$\begin{aligned} g^{11} &= 2 \frac{d\dot{x}^1}{dx^1} \frac{d\dot{x}^1}{dx^1}, g^{21} = 2 \frac{d\dot{x}^2}{dx^1} \frac{d\dot{x}^1}{dx^1}, \\ g^{12} &= 2 \frac{d\dot{x}^1}{dx^1} \frac{d\dot{x}^2}{dx^1}, g^{22} = 2 \frac{d\dot{x}^2}{dx^1} \frac{d\dot{x}^2}{dx^1}, \\ g^{11} &= g^{12} = g^{21} = g^{22} = 2 \frac{d\dot{x}^1}{dx^1} \frac{d\dot{x}^1}{dx^1} \end{aligned} \quad (32)$$

from (12),(19) in consideration of Proposition2 for (31). I rewrite (32) in consideration of (23) and get

$$g^{11} = g^{12} = g^{21} = g^{22} = 4 \frac{d\dot{x}^1}{d\delta} \frac{d\dot{x}^1}{d\delta}. \quad (33)$$

– End Proof –

Proposition 4 *When all coordinate systems satisfies Binary Law, and $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = -Mx^\nu, \frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = \frac{M}{x^\nu}$ is established more, $M = M' - \frac{4}{\dot{x}^1} \frac{d^2 \dot{x}^1}{dx^1 dx^1}, F'_1 = -\frac{8\sqrt{2}}{(\dot{x}^1)^2} \frac{d\dot{x}^1}{d\delta} \frac{d^2 \dot{x}^1}{d\delta d\delta} + \frac{8\sqrt{2}}{\dot{x}^1} \frac{d^3 \dot{x}^1}{d\delta d\delta d\delta}$ is established if a dimensional number is 2.*

Proof: μ, ν -reverses Definition7 and gets

$$M = \frac{\partial x_\mu}{\partial x_\mu} - x_\nu \frac{1}{2} \left(\frac{\partial g^{\mu\nu}}{\partial x^\mu} \right). \quad (34)$$

If the second term of the right side of (34) is zero, I get

$$M = \frac{\partial x_\mu}{\partial x_\mu}. \quad (35)$$

"M" in (34),(35) is the same each. Therefore, I express (35) in

$$M' = \frac{\partial x_\mu}{\partial x_\mu} \quad (36)$$

to distinguish these. I get

$$M = M' - x_\nu \frac{1}{2} \left(\frac{\partial g^{\mu\nu}}{\partial x^\mu} \right) \quad (37)$$

from (34),(36). I get

$$\begin{aligned} M = M' - \dot{x}_1 \frac{1}{2} \left(\frac{\partial g^{11}}{\partial x^1} \right) - \dot{x}_1 \frac{1}{2} \left(\frac{\partial g^{12}}{\partial x^2} \right) \\ - \dot{x}_2 \frac{1}{2} \left(\frac{\partial g^{21}}{\partial x^1} \right) - \dot{x}_2 \frac{1}{2} \left(\frac{\partial g^{22}}{\partial x^2} \right) \end{aligned} \quad (38)$$

from (37) if I assume a dimensional number 2. I get

$$\begin{aligned} M = M' - \frac{1}{\dot{x}^1} \frac{1}{2} \left(\frac{\partial g^{11}}{\partial x^1} \right) - \frac{1}{\dot{x}^1} \frac{1}{2} \left(\frac{\partial g^{12}}{\partial x^2} \right) \\ - \frac{1}{\dot{x}^2} \frac{1}{2} \left(\frac{\partial g^{21}}{\partial x^1} \right) - \frac{1}{\dot{x}^2} \frac{1}{2} \left(\frac{\partial g^{22}}{\partial x^2} \right) \end{aligned} \quad (39)$$

as $\dot{x}_1 = \frac{1}{\dot{x}^1}, \dot{x}_2 = \frac{1}{\dot{x}^2}$ for (38). I get

$$\begin{aligned} M &= M' - \frac{2}{\dot{x}^1} \left(\frac{dg^{11}}{dx^1} \right) \\ &= M' - \frac{4}{\dot{x}^1} \left(\frac{d \frac{d\dot{x}^1}{dx^1} \frac{d\dot{x}^1}{dx^1}}{dx^1} \right) = M' - \frac{8}{\dot{x}^1} \frac{d^2 \dot{x}^1}{dx^1 dx^1} \frac{d\dot{x}^1}{dx^1}, \\ M' &= M + \frac{8}{\dot{x}^1} \frac{d^2 \dot{x}^1}{dx^1 dx^1} \frac{d\dot{x}^1}{dx^1} \end{aligned} \quad (40)$$

from (12),(19),(32) in consideration of Proposition2,Proposition3 for (39). I rewrite (40) in consideration of (23) and get

$$M' = M + \frac{16\sqrt{2}}{\dot{x}^1} \frac{d^2 \dot{x}^1}{d\delta d\delta} \frac{d\dot{x}^1}{d\delta}. \quad (41)$$

I rewrite Definition11 and can get

$$\begin{aligned} W &= \int \vec{F} \cdot d\vec{r} = \int F_\mu \partial r^\mu = \int F^\mu \partial r_\mu, \\ E &= \int F_\mu \partial x^\mu = \int F^\mu \partial x_\mu. \end{aligned} \quad (42)$$

I get (42) as $W \rightarrow E, r^\mu \rightarrow x^\mu, r_\mu \rightarrow x_\mu$ here. I get

$$F_\mu = \frac{\partial E}{\partial x^\mu}, \quad (43)$$

$$F^\mu = \frac{\partial E}{\partial x_\mu} \quad (44)$$

from (42). I get

$$F_\mu = \frac{\partial \frac{M}{\epsilon} c^2}{\partial x^\mu} = \frac{\partial M}{\partial x^\mu} \quad (45)$$

as $c = 1, \epsilon = 1$ from (43), Definition10, Hypothesis1.

$$F_\mu = \frac{\partial M}{\partial x^\mu} = 0 \quad (46)$$

is established in consideration of Definition6 in (45). I get

$$F_1 = \frac{\partial M}{\partial x^1} = 0, F_2 = \frac{\partial M}{\partial x^2} = 0 \quad (47)$$

from (46) if I assume a dimensional number 2. I get

$$F_1 = F_2 = \frac{dM}{dx^1} = 0 \quad (48)$$

from (12) in consideration of Proposition2 for (47). I get

$$\begin{aligned} \frac{dM'}{dx^1} &= \frac{dM}{dx^1} + \frac{d \frac{8}{\dot{x}^1} \frac{d^2 \dot{x}^1}{dx^1 dx^1} \frac{d\dot{x}^1}{dx^1}}{dx^1} = \frac{d \frac{8}{\dot{x}^1} \frac{d^2 \dot{x}^1}{dx^1 dx^1} \frac{d\dot{x}^1}{dx^1}}{dx^1} = F'_1, \\ F'_1 &= -\frac{8}{(\dot{x}^1)^2} \frac{d\dot{x}^1}{dx^1} \frac{d^2 \dot{x}^1}{dx^1 dx^1} \frac{d\dot{x}^1}{dx^1} + \frac{8}{\dot{x}^1} \frac{d^3 \dot{x}^1}{dx^1 dx^1 dx^1} \frac{d\dot{x}^1}{dx^1} + \frac{8}{\dot{x}^1} \frac{d^2 \dot{x}^1}{dx^1 dx^1} \frac{d^2 \dot{x}^1}{dx^1 dx^1} \end{aligned} \quad (49)$$

from (40),(48). I rewrite (49) in consideration of (23) and get

$$F'_1 = -\frac{32}{(\dot{x}^1)^2} \frac{d\dot{x}^1}{d\delta} \frac{d^2 \dot{x}^1}{d\delta d\delta} \frac{d\dot{x}^1}{d\delta} + \frac{32}{\dot{x}^1} \frac{d^3 \dot{x}^1}{d\delta d\delta d\delta} \frac{d\dot{x}^1}{d\delta} + \frac{32}{\dot{x}^1} \frac{d^2 \dot{x}^1}{d\delta d\delta} \frac{d^2 \dot{x}^1}{d\delta d\delta}. \quad (50)$$

– End Proof –

Proposition 5 *When all coordinate systems satisfies Binary Law, $g^{11} = 2\sqrt{2}C_\delta$, $F'_1 = 0$ is established for $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = -Mx^\nu$, $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = \frac{M}{x^\nu}$: $M = 0$ if the number of the dimensions is 2.*

Proof: When $M = 0$ is established, I get

$$\frac{d^2 \dot{x}^1}{d\delta d\delta} = 0 \quad (51)$$

from (24),(25). I get

$$\dot{x}^1 = f(\delta) = C_\delta[2] + \delta C_\delta[1] \quad (52)$$

as a solution of the equations of (51) in consideration of Definition14. C_δ expresses a constant term for variable δ here. I get

$$\dot{x}^1 = f(\delta) = \delta C_\delta[1] = \delta C_\delta \quad (53)$$

as $C_\delta[2] = 0$ for (52). I get

$$\frac{d\dot{x}^1}{d\delta} = C_\delta, \frac{d^2 \dot{x}^1}{d\delta d\delta} = 0, \frac{d^3 \dot{x}^1}{d\delta d\delta d\delta} = 0 \quad (54)$$

from (53). I get

$$g^{11} = g^{12} = g^{21} = g^{22} = 4(C_\delta)^2 \quad (55)$$

from (33),(54). I get

$$F'_1 = 0 \quad (56)$$

from (50),(53),(54).

– End Proof –

Proposition 6 *When all coordinate systems satisfies Binary Law,*

$g^{11} = 2M \left(C_\delta \cos \left(\frac{\sqrt{M}\delta}{\sqrt{2}} \right) \right)^2$, $F'_1 = 8M^2 C_\delta \sin \left(\frac{\sqrt{M}\delta}{\sqrt{2}} \right)$, $g^{11} = 4 \frac{dh(\delta)}{d\delta} \frac{dh(\delta)}{d\delta}$, $F'_1 = -\frac{32}{(h(\delta))^2} \frac{dh(\delta)}{d\delta} \frac{d^2 h(\delta)}{d\delta d\delta} \frac{dh(\delta)}{d\delta} + \frac{32}{h(\delta)} \frac{d^3 h(\delta)}{d\delta d\delta d\delta} \frac{dh(\delta)}{d\delta} + \frac{32}{h(\delta)} \frac{d^2 h(\delta)}{d\delta d\delta} \frac{d^2 h(\delta)}{d\delta d\delta}$ is established for $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = -Mx^\nu$, $\frac{\partial^2 x^\nu}{\partial x^\mu \partial x^\mu} = \frac{M}{x^\nu}$: $M > 0$ if the number of the dimensions is 2.

Proof: When $M > 0$ is established, I get

$$\dot{x}^1 = g(\delta) = C_\delta[2] \cos\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right) + C_\delta[1] \sin\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right) \quad (57)$$

as a solution of the equations of (24) in consideration of Definition15. I get

$$\dot{x}^1 = g(\delta) = C_\delta[1] \sin\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right) = C_\delta \sin\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right) \quad (58)$$

as $C_\delta[2] = 0$ for (57). I get

$$\begin{aligned} \frac{d\dot{x}^1}{d\delta} &= \frac{\sqrt{M}}{\sqrt{2}} C_\delta \cos\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right), \quad \frac{d^2\dot{x}^1}{d\delta d\delta} = -\frac{M}{2} C_\delta \sin\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right), \\ \frac{d^3\dot{x}^1}{d\delta d\delta d\delta} &= -\frac{M}{2} \frac{\sqrt{M}}{\sqrt{2}} C_\delta \cos\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right) \end{aligned} \quad (59)$$

from (58). I get

$$g^{11} = g^{12} = g^{21} = g^{22} = 2M \left(C_\delta \cos\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right) \right)^2 \quad (60)$$

from (33),(59). I get

$$\begin{aligned} F'_1 &= 8M^2 \cot\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right) C_\delta \cos\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right) \\ &\quad - 8M^2 \cot\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right) C_\delta \cos\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right) \\ &\quad + 8M^2 C_\delta \sin\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right) \\ &= 8M^2 C_\delta \sin\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right) \end{aligned} \quad (61)$$

from (50),(58),(59). I show details form about (61) to Fig.1. I get

$$0 = F'_1 = 8M^2 C_\delta \sin\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right) = 8M^2 C_\delta \sin(\pi n) \quad (62)$$

from (61).

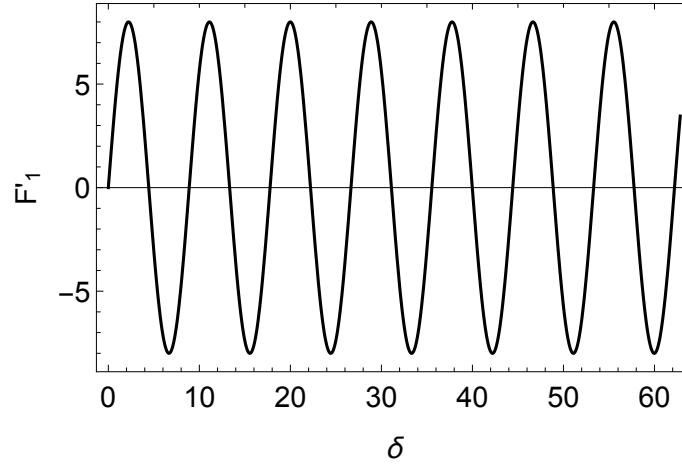


Figure 1: $(F_1' - \delta) : F_1' = 8M^2C_\delta \sin\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right)$ for $\frac{d^2\dot{x}^1}{dx^1 dx^1} = -M\dot{x}^1 : \{M = 1, C_\delta = 1\}$. [8]

When (62) is established,

$$\delta = \sqrt{2}\pi n \quad (63)$$

is established. I assumed it $M = 1$ here. δ when $F_1' = 0$ is satisfied is

$$\delta = 0, \sqrt{2}\pi, 2\sqrt{2}\pi, 3\sqrt{2}\pi, 4\sqrt{2}\pi, 5\sqrt{2}\pi, \dots \quad (64)$$

from (63). I show details form about Central $(0, 0)$, the equation of the circle of radius δ in consideration of Definition9 to Fig.2. I assumed it δ according to (64) here. $F_1' \neq 0$ is established any place other than the domain of δ according to (64). I get

$$0 = F_1' = 8M^2C_\delta \sin\left(\frac{\sqrt{M}\delta}{\sqrt{2}}\right) = 8M^2C_\delta \sin(2\pi) \quad (65)$$

from (61). When (65) is established,

$$\delta = \frac{2\pi\sqrt{2}}{\sqrt{M}} \quad (66)$$

is established. δ of (66) accords in wavelength λ of (61).

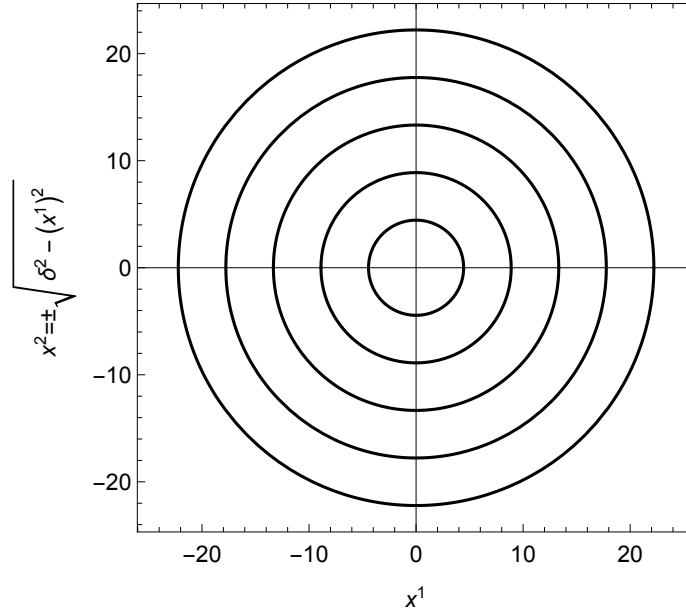


Figure 2: $(x^2 - x^1) : x^2 = \pm\sqrt{\delta^2 - (x^1)^2}$ details form about Central $(0,0)$, the equation of the circle of radius $\delta = \sqrt{2}\pi, 2\sqrt{2}\pi, 3\sqrt{2}\pi, 4\sqrt{2}\pi, 5\sqrt{2}\pi$. [8]

Therefore, I get

$$\lambda = \frac{2\pi\sqrt{2}}{\sqrt{M}} \tag{67}$$

from (66).

When $M > 0$ is established, I get

$$\dot{x}^1 = h(\delta) \tag{68}$$

as a solution of the equations of (25). I show details form by the numerical computation about (68) to Fig.3. I get

$$g^{11} = g^{12} = g^{21} = g^{22} = 4 \frac{dh(\delta)}{d\delta} \frac{dh(\delta)}{d\delta} \tag{69}$$

from (33),(68). I show details form by the numerical computation about (69) to Fig.4.

I get

$$M' = M + \frac{16\sqrt{2}}{h(\delta)} \frac{d^2h(\delta)}{d\delta d\delta} \frac{dh(\delta)}{d\delta} \tag{70}$$

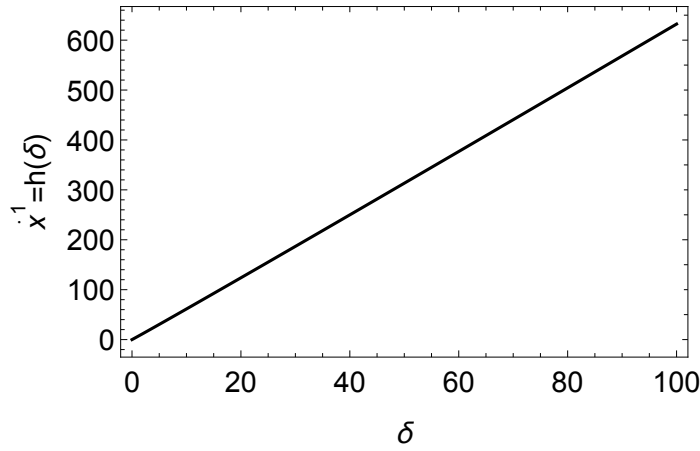


Figure 3: $(\dot{x}^1 - \delta) : \dot{x}^1 = h(\delta)$ numerical computation plot for $\frac{d^2 \dot{x}^1}{dx^1 dx^1} = \frac{M}{\dot{x}^1} : \{M = 1\}$. [8]

from (41),(68). I get

$$F'_1 = -\frac{32}{(h(\delta))^2} \frac{dh(\delta)}{d\delta} \frac{d^2 h(\delta)}{d\delta d\delta} \frac{dh(\delta)}{d\delta} + \frac{32}{h(\delta)} \frac{d^3 h(\delta)}{d\delta d\delta d\delta} \frac{dh(\delta)}{d\delta} + \frac{32}{h(\delta)} \frac{d^2 h(\delta)}{d\delta d\delta} \frac{d^2 h(\delta)}{d\delta d\delta} \tag{71}$$

from (50),(68). I show details form by the numerical computation about (70),(71) to Fig.5, Fig.6.

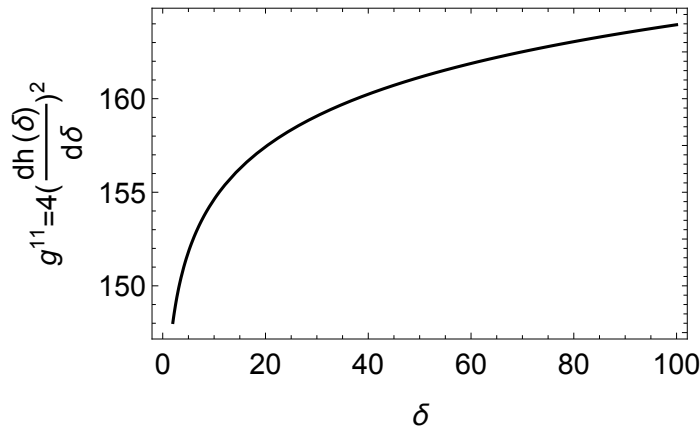


Figure 4: $(g^{11} - \delta) : g^{11} = 4 \left(\frac{dh(\delta)}{d\delta}\right)^2$ numerical computation plot for $\frac{d^2 \dot{x}^1}{dx^1 dx^1} = \frac{M}{\dot{x}^1} : \{M = 1\}$. [8]

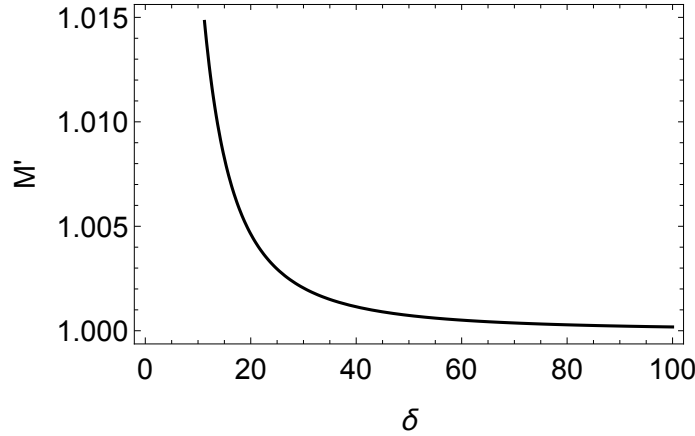


Figure 5: $(M' - \delta) : M' = M + \frac{16\sqrt{2}}{h(\delta)} \frac{d^2 h(\delta)}{d\delta d\delta} \frac{dh(\delta)}{d\delta}$ numerical computation plot for $\frac{d^2 \dot{x}^1}{dx^1 dx^1} = \frac{M}{\dot{x}^1} : \{M = 1\}$. [8]

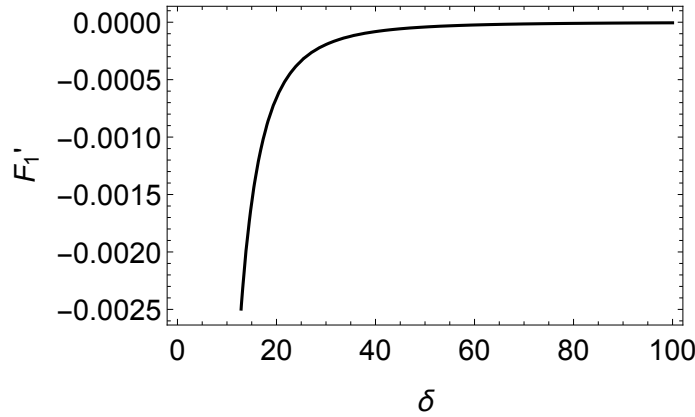


Figure 6: $(F_1' - \delta) : F_1' = -\frac{32}{(h(\delta))^2} \frac{dh(\delta)}{d\delta} \frac{d^2 h(\delta)}{d\delta d\delta} \frac{dh(\delta)}{d\delta} + \frac{32}{h(\delta)} \frac{d^3 h(\delta)}{d\delta d\delta d\delta} \frac{dh(\delta)}{d\delta} + \frac{32}{h(\delta)} \frac{d^2 h(\delta)}{d\delta d\delta} \frac{d^2 h(\delta)}{d\delta d\delta}$ numerical computation plot for $\frac{d^2 \dot{x}^1}{dx^1 dx^1} = \frac{M}{\dot{x}^1} : \{M = 1\}$. [8]

– End Proof –

4 About outbreak mechanism of the gravitation to occur between two matter

I show details form about Central $(0, 0)$, the equation of the circle of radius $\delta = 1$ in consideration of Definition9 to Fig.7. I get Fig.8 in consideration of

Fig.6 for Fig.7.

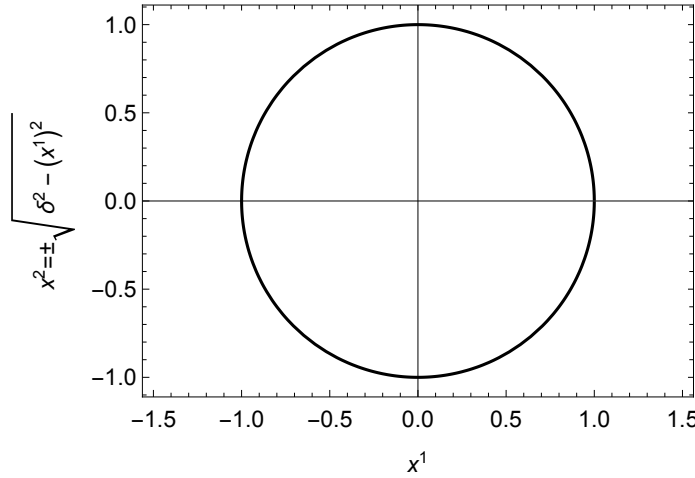


Figure 7: $(x^2 - x^1) : x^2 = \pm\sqrt{\delta^2 - (x^1)^2}$ details form about Central $(0,0)$, the equation of the circle of radius $\delta = 1$. [8]

The arrow of Fig.8 expresses Vector field $\vec{F}'_{\mu}(\delta) : (F'_1(\delta), F'_2(\delta)) = (F'_1(\delta), F'_1(\delta))$ in Fig.6.

All $|\vec{F}'_{\mu}(\delta)| = \sqrt{(F'_1(\delta))^2 + (F'_2(\delta))^2} = \sqrt{2}F'_1(\delta)$ of Vector field $\vec{F}'_{\mu}(\delta)$ showed it to equivalence. $F'_1 = F'_2$ is established from (48) here. Because I don't consider it in this article about the change for δ of $|\vec{F}'_{\mu}(\delta)|$, I express it as $|\vec{F}'_{\mu}(\delta)| = C_{\delta}$ in Fig.8. C_{δ} expresses invariable for variable δ .

Area $\alpha : (\delta > 0)$ is compressed equally if I switch from point $\alpha : (\delta = 0)$ to area of the circle α of radius $\delta > 0$ in Fig.8. However, the power to move area $\alpha : (\delta > 0)$ to a particular direction on two dimensions doesn't occur. In other words, area $\alpha : (\delta > 0)$ is only compressed equally.

If $|\vec{F}'_{\mu}(\delta)|$ of field of force $\vec{F}'_{\mu}(\delta) : \{1, 2\}$ is reduced in Fig.8 here, the power of the direction of field of force $\vec{F}'_{\mu}(\delta) : \{3, 4\}$ will appear for area $\alpha : (\delta > 0)$. In other words, a symmetric break of field of force $\vec{F}'_{\mu}(\delta)$ in Fig.8 causes power to move area $\alpha : (\delta > 0)$ to a particular direction. The power to appear grows big so that a symmetric break is big.

I think about the case that point α producing field of force $\vec{F}'_{\mu}(\delta)$ exists one more in this neighborhood. I express another point in β here. I get Fig.9 as $(-320 < x^1 < 320) \rightarrow (-250 < x^1 < 650), (-320 < x^2 < 320) \rightarrow (-250 < x^2 < 650)$ in Fig.8. I get Fig.10 as point $\alpha : (0, 0) \rightarrow \alpha : (400, 400), \alpha : (400, 400) \rightarrow \beta : (400, 400)$ in Fig.9. I superimpose Fig.10 on Fig.9 and get

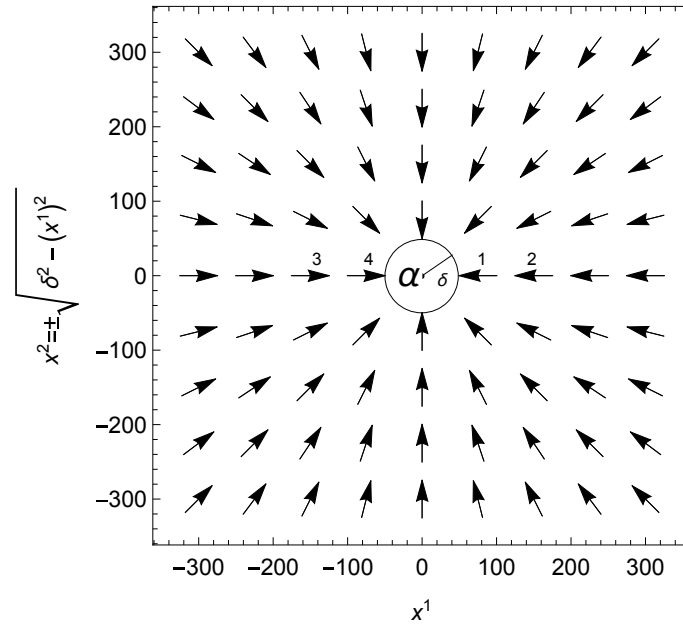


Figure 8: Vector field $\vec{F}'_\mu(\delta) : (F'_1(\delta), F'_2(\delta)) = (F'_1(\delta), F'_1(\delta))$ for point $\alpha : (\delta = 0)$. [8]

Fig.11. Then, I show that field of force $\vec{F}'_\mu(\delta)$ by point β breaks a symmetry of field of force $\vec{F}'_\mu(\delta)$ by point α . Furthermore, I show that field of force $\vec{F}'_\mu(\delta)$ by point α breaks a symmetry of field of force $\vec{F}'_\mu(\delta)$ by point β . The direction of field of force $\vec{F}'_\mu(\delta) : \{3, 4, 5\}$ of Fig.10 is reverse to field of force $\vec{F}'_\mu(\delta) : \{3, 4, 5\}$ of Fig.9. Therefore, $|\vec{F}'_\mu(\delta)|$ of field of force $\vec{F}'_\mu(\delta) : \{3, 4, 5\}$ is reduced in Fig.11. Thus, field of force $\vec{F}'_\mu(\delta) : \{1, 2\}$ appears for point α , and field of force $\vec{F}'_\mu(\delta) : \{6, 7\}$ appears for point β . There seems to be power to attract each other between point α and point β .

5 Discussion

Because $F'_1(\delta) = 0$ is established in the case of $M = 0$ in Proposition5, it is $\vec{F}'_\mu(\delta) : (F'_1(\delta), F'_1(\delta)) = (0, 0)$. In the case of $M > 0$, $F'_1(\delta) : Fig.1$ and $F'_1(\delta) : Fig.6$ are established together in Proposition6. Because size of the wavelength of $F'_1(\delta) : Fig.1$ is limited numerical value, we don't need to consider existence of $F'_1(\delta) : Fig.1$ on the scale to greatly exceed size of the wavelength.

About $\frac{\partial^3 x^\nu}{\partial x^\mu \partial x^\mu \partial x^\mu} = M$

When Binary Law is established, all coordinate systems is only two of x^μ, x^ν . [1]

I get $\frac{\partial^3 x^\mu}{\partial x^\nu \partial x^\nu \partial x^\nu} = M$ as μ, ν -inversion form of $\frac{\partial^3 x^\nu}{\partial x^\mu \partial x^\mu \partial x^\mu} = M$. Because

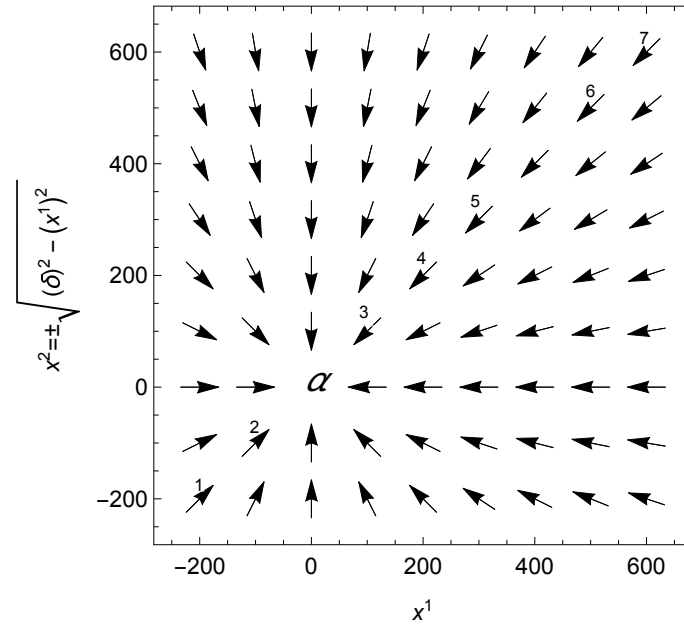


Figure 9: Vector field $\vec{F}'_\mu(\delta) : (F'_1(\delta), F'_2(\delta)) = (F'_1(\delta), F'_1(\delta))$ for point $\alpha : (\delta = 0)$. [8]

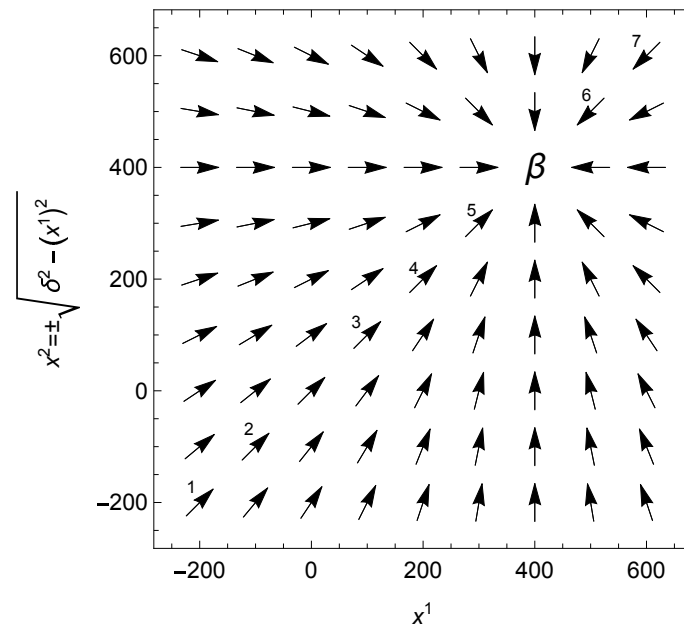


Figure 10: Vector field $\vec{F}'_\mu(\delta) : (F'_1(\delta), F'_2(\delta)) = (F'_1(\delta), F'_1(\delta))$ for point $\beta : (\delta = 0)$. [8]

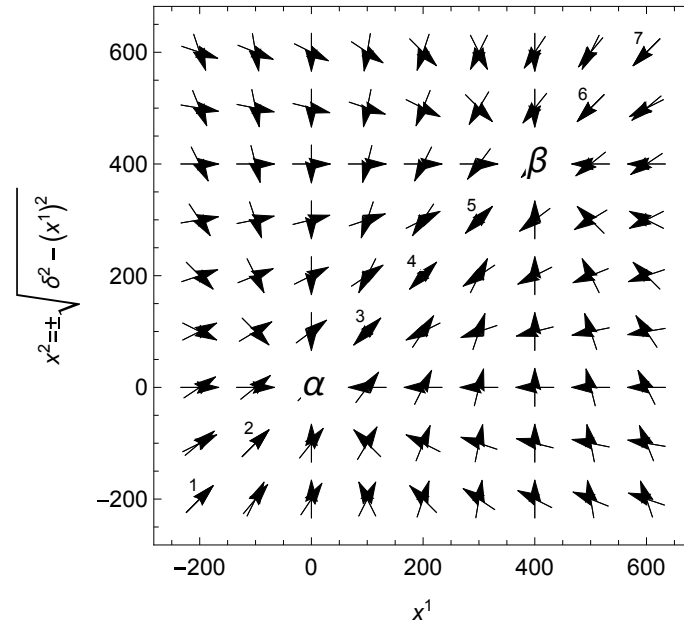


Figure 11: Vector field $\vec{F}'_\mu(\delta) : (F'_1(\delta), F'_2(\delta)) = (F'_1(\delta), F'_1(\delta))$ for point $\alpha : (\delta = 0), \beta : (\delta = 0)$. [8]

$\frac{\partial^3 x^\mu}{\partial x^\nu \partial x^\nu \partial x^\nu}, \frac{\partial^3 x^\nu}{\partial x^\mu \partial x^\mu \partial x^\mu}$ became the scalar each, it was found that all coordinate systems x^μ, x^ν obeyed these differential equations mathematically each. It wasn't known that all coordinate systems could have such a property in the past. This is a new discovery fact and is discovery in tensor satisfying Binary Law. M to show here is scalar. Therefore, all coordinate systems x^μ, x^ν can adopt three states based on $\{M = 0, M > 0, M < 0\}$ each.

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- [7] Calculation result obtained by the author using Mathematica 11.3J Home Edition.
- [8] Figure obtained by the author by Mathematica 11.3J Home Edition.

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