

The Random Phase Approximation for Binary Hard-Core Mixture

Arkadiy B. Finkel'shtein

Ural Federal University, Mira st. 19, 620002, Ekaterinburg, Russia

Copyright © 2014 Arkadiy B. Finkel'shtein. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.

Abstract

Expressions of the Ashcroft-Langreth partial structure factors for an arbitrary two-component hard-core fluid in the random phase approximation are presented.

Keywords: Hard-core mixture, random phase approximation, partial structure factor

In general case, the random phase approximation (RPA) for classical binary fluid is formulated as follows:

$$c_{ij}^{\text{RPA}}(r) = c_{0ij}(r) - \beta \varphi_{1ij}(r) \quad , \quad (1)$$

where $c_{ij}(r)$ is the partial direct correlation function, $\varphi_{ij}(r)$ - partial pair interatomic potential, $\beta = (k_B T)^{-1}$, k_B - Boltzmann constant, T - temperature, symbols “0” and “1” are attributes of a reference system and perturbation, respectively, $i, j=1, 2$.

For any hard-core (HC) $\varphi_{ij}(r)$,

$$\varphi_{ij}^{\text{HC}}(r) = \begin{cases} \infty, & r < \sigma_{ij} \\ \phi_{ij}(r), & r \geq \sigma_{ij} \end{cases} \quad , \quad (2)$$

where σ_{ij} is the partial HC diameter,

$$c_{ij}^{\text{RPA-HC}}(r) = c_{ij}^{\text{HS}}(r) - \beta \phi_{ij}(r) \quad . \quad (3)$$

Here, “HS” denotes the hard-sphere model, which for a mixture is described by the following $\varphi_{ij}(r)$:

$$\phi_{ij}^{\text{HS}}(r) = \begin{cases} \infty, & r < \sigma_{ij} \\ 0, & r \geq \sigma_{ij} \end{cases}, \quad (4)$$

In the wave space Eq. (3) is being rewritten as

$$c_{ij}^{\text{RPA-HC}}(q) = c_{ij}^{\text{HS}}(q) - \beta \phi_{ij}(q). \quad (5)$$

The Lebowitz exact solution [1] obtained for the additive HS binary mixture in the Percus-Yevick approximation [2] can be used to calculate $c_{ij}^{\text{HS}}(r)$ and $c_{ij}^{\text{HS}}(q)$.

The Ashcroft-Langreth (AL) partial structure factors [3] are expressed as follows ($i \neq j$):

$$S_{ii}(q) = \frac{1 - c_j \rho c_{jj}(q)}{\prod_{k=1,2} (1 - c_k \rho c_{kk}(q)) - c_i c_j \rho^2 c_{ij}^2(q)}, \quad (6)$$

$$S_{ij}(q) = \frac{\sqrt{c_i c_j} \rho c_{ij}(q)}{\prod_{k=1,2} (1 - c_k \rho c_{kk}(q)) - c_i c_j \rho^2 c_{ij}^2(q)}, \quad (7)$$

where c_i is the concentration of the i -th component, ρ - mean atomic density of a mixture. Combining Eqs. (6) and (7) with Eq. (5) one can obtain after some transformations the expressions of AL partial structure factors within the RPA for an arbitrary $\phi_{ij}(r)$ in a general form:

$$S_{ii}^{\text{RPA-HC}}(q) = \frac{1 - c_j \rho c_{jj}^{\text{HS}}(q) + c_j \rho \beta \phi_{jj}(q)}{Z(q)}, \quad (8)$$

$$S_{ij}^{\text{RPA-HC}}(q) = \frac{\sqrt{c_i c_j} \rho c_{ij}^{\text{HS}}(q) - \sqrt{c_i c_j} \rho \beta \phi_{ij}(q)}{Z(q)}, \quad (9)$$

$$Z(q) = \prod_{k=1,2} (1 - c_k \rho c_{kk}^{\text{HS}}(q)) - c_i c_j \rho^2 c_{ij}^{\text{HS}2}(q) + \prod_{k=1,2} (1 + c_k \rho \beta \phi_{kk}(q)) - c_i c_j \rho^2 \beta^2 \phi_{ij}^2(q) - \\ - 1 - c_i c_j \rho^2 \beta \left(\sum_{i,j} c_{ii}^{\text{HS}}(q) \phi_{jj}(q) - 2 c_{ij}^{\text{HS}}(q) \phi_{ij}(q) \right). \quad (10)$$

Expressions obtained can be used to describe the square-well, triangular-well, hard-core-Yukawa, charged-hard-sphere and other model binary mixtures.

References

- [1] J.L. Lebowitz, Exact solution of generalized Percus-Yevick equation for a mixture of hard spheres, Phys. Rev. A, 133 (1964), 895-899.
- [2] J.K. Percus, G.Y. Yevick, Analysis of classical statistical mechanics by means of collective coordinates, Phys. Rev., 110 (1958), 1-13.
- [3] N.W. Ashcroft, D.C. Langreth, Structure of binary liquid mixtures. II., Phys. Rev., 156 (1967), 685-692.

Received: February 28, 2014