

The Characteristic Polynomial of Some Perturbed Tridiagonal k -Toeplitz Matrices¹

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Abstract. We generalize some recent results on the spectra of tridiagonal matrices, providing explicit expressions for the characteristic polynomial of some perturbed tridiagonal k -Toeplitz matrices. The calculation of the eigenvalues (and associated eigenvectors) follows straightforward.

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1. PRELIMINARIES

An $n \times n$ matrix $A = [a_{ij}]$ is said tridiagonal if $a_{ij} = 0$, whenever $|i - j| > 1$, i.e.,

$$(1.1.1) \quad A = \begin{pmatrix} a_1 & c_1 & & & \\ b_1 & \ddots & \ddots & & \\ & \ddots & \ddots & c_{n-1} & \\ & & b_{n-1} & a_n & \end{pmatrix},$$

with non-mentioned entries equal to zero. Since our aim will be the study of the characteristic polynomial of some tridiagonal matrices, the cases when $b_1 \cdots b_{n-1} = 0$ or $c_1 \cdots c_{n-1} = 0$ (i.e., when A is reducible) can be reduced to small order cases. In fact, throughout we will assume that $b_\ell c_\ell > 0$, for $\ell = 1, \dots, n - 1$.

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relation:

$$x \begin{pmatrix} P_0(x) \\ P_1(x) \\ \vdots \\ P_n(x) \end{pmatrix} = J_{n+1} \begin{pmatrix} P_0(x) \\ P_1(x) \\ \vdots \\ P_n(x) \end{pmatrix} + \alpha_n P_{n+1}(x) \begin{pmatrix} 0 \\ 0 \\ \vdots \\ 1 \end{pmatrix},$$

where J_{n+1} is a tridiagonal matrix of order $n + 1$, defined by

$$J_{n+1} := \begin{pmatrix} \beta_0 & \alpha_0 & & & \\ \gamma_1 & \beta_1 & \alpha_1 & & \\ & \gamma_2 & \ddots & \ddots & \\ & & \ddots & \ddots & \alpha_{n-1} \\ & & & \gamma_n & \beta_n \end{pmatrix}, \quad \text{for } n = 0, 1, 2, \dots$$

Hence each zero of the polynomial P_n , say λ_{nj} (for $j = 1, \dots, n$), is an eigenvalue of the corresponding tridiagonal matrix J_n of order n , and an associated eigenvector is

$$(2.2.2) \quad (P_0(\lambda_{nj}), P_1(\lambda_{nj}), \dots, P_{n-1}(\lambda_{nj}))^t.$$

The (monic) characteristic polynomial of J_n is precisely P_n , with $\alpha_n = 1$, i.e.,

$$P_n(x) = \det(xI_n - J_n), \quad n = 1, 2, \dots,$$

where I_n denotes the identity matrix of order n .

It is important to point out that, since we are interested on the eigenvalues of tridiagonal matrices, the eigenvalues of J_n are the same of the matrix

$$\begin{pmatrix} \beta_0 & 1 & & & \\ \gamma_1 \alpha_0 & \beta_1 & 1 & & \\ & \gamma_2 \alpha_1 & \ddots & \ddots & \\ & & \ddots & \ddots & 1 \\ & & & \gamma_{n-1} \alpha_{n-2} & \beta_{n-1} \end{pmatrix}$$

or

$$\begin{pmatrix} \beta_0 & \sqrt{\gamma_1 \alpha_0} & & & \\ \sqrt{\gamma_1 \alpha_0} & \beta_1 & \sqrt{\gamma_2 \alpha_1} & & \\ & \sqrt{\gamma_2 \alpha_1} & \ddots & \ddots & \\ & & \ddots & \ddots & \sqrt{\gamma_{n-1} \alpha_{n-2}} \\ & & & \sqrt{\gamma_{n-1} \alpha_{n-2}} & \beta_{n-1} \end{pmatrix}.$$

Since all these matrices have the same characteristic polynomial, we will simply consider

$$\tilde{J}_n = \begin{pmatrix} \beta_0 & 1 & & & \\ \gamma_1 & \beta_1 & 1 & & \\ & \gamma_2 & \ddots & \ddots & \\ & & \ddots & \ddots & 1 \\ & & & \gamma_{n-1} & \beta_{n-1} \end{pmatrix},$$

with positive γ_i 's, according to our original assumption. It follows from this fact the uselessness of the separation in three "different" families of tridiagonal matrices made in [10].

3. CHEBYSHEV POLYNOMIALS OF SECOND KIND

The Chebyshev polynomials of second kind, denoted by $\{U_n\}_{n \geq 0}$, are a very useful example of a (monic) OPS. These polynomials satisfy the three-term recurrence relations

$$(3.3.1) \quad U_{n+1}(x) = xU_n(x) - U_{n-1}(x), \quad \text{for all } n = 1, 2, \dots$$

with initial conditions $U_0(x) = 1$ and $U_1(x) = 2x$. Since each U_n satisfies

$$(3.3.2) \quad U_n(x) = \frac{\sin(n+1)\theta}{\sin \theta}, \quad \text{with } x = \cos \theta \quad (0 \leq \theta < \pi),$$

for all $n = 0, 1, 2, \dots$, we can deduce the orthogonality relations

$$\int_{-1}^1 U_n(x)U_m(x) \sqrt{1-x^2} dx = \frac{\pi}{2} \delta_{n,m}$$

(cf. [2], e.g.). It is also well known the explicit formula for Chebyshev polynomials of second kind

$$U_n(x) = \sum_{k=0}^{\lfloor n/2 \rfloor} (-1)^k \frac{(n-k)!}{k!(n-2k)!} (2x)^{n-2k}.$$

Consider now the $n \times n$ real tridiagonal Toeplitz matrix T_n defined by

$$(3.3.3) \quad T_n = \begin{pmatrix} a & 1 & & & \\ b & a & \ddots & & \\ & \ddots & \ddots & 1 & \\ & & b & a & \end{pmatrix} \in \mathbb{R}^{n \times n},$$

with $b > 0$. The characteristic polynomial of T_n is $P_n(\lambda) = (\sqrt{b})^n U_n\left(\frac{\lambda-a}{2\sqrt{b}}\right)$ and it follows immediately that the eigenvalues of T_n are

$$\lambda_\ell = a + 2\sqrt{b} \cos\left(\frac{\ell\pi}{n+1}\right),$$

for $\ell = 1, 2, \dots, n$, since they are the zeros of P_n , with associated eigenvectors

$$\left(\sin\left(\frac{\ell\pi}{n+1}\right), \sqrt{b} \sin\left(\frac{2\ell\pi}{n+1}\right), \dots, (\sqrt{b})^{n-1} \sin\left(\frac{n\ell\pi}{n+1}\right) \right)^t,$$

from (2.2.2).

Notice that if $b < 0$, then $P_n(\lambda) = (i\sqrt{-b})^n U_n\left(\frac{\lambda-b}{2i\sqrt{-a}}\right)$.

4. EIGENVALUES OF TRIDIAGONAL 2-TOEPLITZ MATRIX

In 1966, Rózsa held a seminar at University of Hamburg on tridiagonal k -Toeplitz matrices, called at that time as "periodic continuants", motivated mainly by some problems of lattice dynamics, of ladder networks and of structural analysis. In this occasion, Elsner and Redheffer had the idea to use Chebyshev polynomials of the second kind to derive formulas to the characteristic polynomial and the eigenvectors of that matrices in some particular cases. They published some new results one year later in [3] and independently, but slightly later, Rózsa in [16] presented similar propositions using some tools from theory of matrices. Later, these results were generalized to periodic block-tridiagonal matrices by Rózsa and Romani [17].

Recently the study of general tridiagonal k -Toeplitz matrices was considered again by the author and Petronilho in [6], after the characteristic polynomial of a tridiagonal 2-Toeplitz matrix had been considered by Gover [7], Marcellán and Petronilho [14] and Fonseca and Petronilho [5]. Based on this last paper, recently Álvarez-Nodarse, Petronilho and Quintero [1] studied the spectra of tridiagonal 2 and 3-Toeplitz matrices motivated by some Quantum Physics problems already mentioned.

Considering the tridiagonal 2-Toeplitz matrix

$$(4.4.1) \quad T_N^{(2)} = \begin{pmatrix} a_1 & 1 & & & \\ b_1 & a_2 & 1 & & \\ & b_2 & a_1 & 1 & \\ & & b_1 & \ddots & \ddots \\ & & & \ddots & \ddots \end{pmatrix}_{N \times N},$$

let us define the polynomials

$$\pi_2(x) = (x - a_1)(x - a_2)$$

and

$$P_k^*(x) = \left(\sqrt{b_1 b_2}\right)^k U_k \left(\frac{x - b_1 - b_2}{2\sqrt{b_1 b_2}}\right).$$

It is well known (cf. [5, 7, 14]) that the eigenvalues of the matrix (4.4.1) are the solutions of the polynomial Q_N defined by

$$Q_{2k+1}(x) = (x - a_1)P_k^*(\pi_2(x))$$

and

$$Q_{2k}(x) = P_k^*(\pi_2(x)) + b_2 P_{k-1}^*(\pi_2(x)).$$

In the case of $N = 2n + 1$, eigenvalues of $T_N^{(2)}$ are a_1 and

$$\frac{a_1 + a_2}{2} \pm \sqrt{\frac{(a_1 - a_2)^2}{4} + b_1^2 + b_2^2 + 2\sqrt{b_1 b_2} \cos\left(\frac{k\pi}{n+1}\right)},$$

The procedure used in the section 5 can be used to determine the eigenvalues (and naturally the eigenvectors by (2.2.2) of any perturbed tridiagonal k -Toeplitz matrix, with a perturbation similar to the one occurred in the tridiagonal 2-Toeplitz matrix.

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